

Composite Operator

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In some cases, we need to consider Green's function of composite operator.

Consider the operator $\Omega(\frac{1}{2}\phi^2)$ in $\lambda\phi^4$ theory. Green's function with one insertion of Ω is

$$G_{\Omega}^{(n)}(x; x_1, x_2, \dots, x_n) = \langle 0 | T(\frac{1}{2}\phi^2(x) \phi(x_1) \phi(x_2) \dots \phi(x_n)) | 0 \rangle$$

In momentum space

$$(2\pi)^4 \delta^4(p+p_1+\dots+p_n) G_{\Omega}^{(n)}(p; p_1, \dots, p_n) = \int d^4x e^{-ipx} \int \prod_{i=1}^n d^4x_i e^{-ip_i x_i} G_{\Omega}^{(n)}(x; x_1, \dots, x_n)$$

In perturbation theory, we can use Wick's theorem to work out these Green's functions in terms of Feynman diagram

Example, to lowest order in λ

$$G_{\Omega}^{(2)}(x; x_1, x_2) = \frac{1}{2} \langle 0 | T(\phi^2(x) \phi(x_1) \phi(x_2)) | 0 \rangle = i\Delta(x-x_1)\Delta(x-x_2)$$

or in momentum space

$$G_{\Omega}^{(2)}(p; p_1, p_2) = i\Delta(p_1)i\Delta(p+p_1)$$

If we truncate the external propagators, we have

$$\Gamma_{\Omega}^{(2)}(p, p_1, -p_1-p) = 1$$

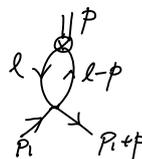
To first order in λ , we have

$$G_{\Omega}^{(2)}(x, x_1, x_2) = \int d^4y \langle 0 | T(\frac{1}{2}\phi^2(x) \phi(x_1) \phi(x_2) \frac{(i\lambda)}{4!}\phi^4(y)) | 0 \rangle = \int d^4y \left(\frac{i\lambda}{2}\right) [i\Delta(x-y)]^2 i\Delta(x_1-y) i\Delta(x_2-y)$$



The amputated 1PI momentum space Green's is

$$\Gamma_{\Omega}^{(2)}(p; p_1, -p-p_1) = \left(\frac{i\lambda}{2}\right) \int \frac{d^4k}{(2\pi)^4} \left(\frac{i}{k^2-\mu^2+i\epsilon}\right) \frac{i}{(k-p)^2-\mu^2+i\epsilon}$$



To calculate this type of Green's functions systematically, we can add a term $\chi(x)\Omega(x)$ to \mathcal{L}

$$\mathcal{L}[\chi] = \mathcal{L}[\phi] + \chi(x)\Omega(x)$$

where $\chi(x)$ is a c-number source function. We can construct the generating functional $W[\chi]$ in the presence of this external source. We obtain the connected Green's function by differentiating $\ln W[\chi]$ with respect to χ and then setting χ to zero.

Renormalization of Composite operators

Superficial degrees of divergence

$$D_{\Omega} = D + \delta_{\Omega} = D + (d_{\Omega} - 4)$$

where d_{Ω} is the canonical dimension of Ω

For the case of $\Omega(x) = \frac{1}{2}\phi^2(x)$, $d_{\phi^2} = 2$ and $D_{\phi^2} = 2 - n \Rightarrow$ only $\Gamma_{\phi^2}^{(2)}$ is divergent.

Taylor expansion

$$\Gamma_{\phi^2}^{(2)}(p; p_1) = \Gamma_{\phi^2}^{(2)}(0, 0) + \Gamma_{\phi^2}^{(2)}(p, p_1)$$

We can combine the counter term

$$-\frac{i}{2} \Gamma_{\phi^2}^{(2)}(0, 0) \chi(x) \phi^2(x)$$

with the original term to write

$$-\frac{i}{2} \chi \phi^2 = -\frac{i}{2} \int \frac{d^3x}{\phi^2} \chi \phi^2 = -\frac{i}{2} Z_\phi \chi \phi^2 \quad Z_\phi^2 = 1 + \Gamma_{\phi^2}(0,0)$$

In general, we need to insert counter term $\Delta\Omega$ into the original action

$$\mathcal{L} \rightarrow \mathcal{L} + \chi(\Omega + \Delta\Omega)$$

If $\Delta\Omega = c\Omega$, as in the case of $\Omega = \frac{1}{2}\phi^2$, we have

$$\mathcal{L}[\chi] = \mathcal{L}[\phi] + \chi Z_c \Omega = \mathcal{L}[\phi] + \chi \Omega_0$$

with $\Omega_0 = Z_c \Omega = (1+c)\Omega$

Such composite operators are said to be multiplicatively renormalizable. and Green's functions of unrenormalized operator Ω_0 is related to that of renormalized operator Ω by

$$G_{\Omega_0}^{(n)}(x; x_1, \dots, x_n) = \langle 0 | T(\Omega_0(x) \phi(x_1) \dots \phi(x_n)) | 0 \rangle \\ = Z_\Omega Z_\phi^{n/2} G_{\Omega}^{(n)}(x; x_1, \dots, x_n)$$

For more general cases, $\Delta\Omega \neq c\Omega$ and the renormalization of a composite operator may require counter term proportional to other composite operators

Example, 2 composite operators A and B

$$\mathcal{L}[\chi] = \mathcal{L}[\phi] + \chi_A(A + \Delta A) + \chi_B(B + \Delta B)$$

The counter terms ΔA and ΔB are linear combinations of A and B

$$\Delta A = C_{AA} A + C_{AB} B \\ \Delta B = C_{BA} A + C_{BB} B$$

We can write

$$\mathcal{L}[\chi] = \mathcal{L}[\phi] + (\chi_A, \chi_B) C \begin{pmatrix} A \\ B \end{pmatrix} \quad C = \begin{pmatrix} 1 + C_{AA} & C_{AB} \\ C_{BA} & 1 + C_{BB} \end{pmatrix}$$

Diagonalize C by bi-unitary transformation

$$U C V^t = \begin{pmatrix} Z_{A'} & 0 \\ 0 & Z_{B'} \end{pmatrix}$$

then

$$\mathcal{L}[\chi] = \mathcal{L}[\phi] + Z_{A'} \chi_{A'} A' + Z_{B'} \chi_{B'} B'$$

$$\begin{pmatrix} A' \\ B' \end{pmatrix} = V \begin{pmatrix} A \\ B \end{pmatrix}$$

$$(\chi_{A'}, \chi_{B'}) = (\chi_A, \chi_B) U$$

A', B' are multiplicatively renormalizable.