

# Neutrinos from Gamma-ray Burst

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4 April, 2011

## Outline

### Objectives and motives

Ultra high energy cosmic ray and Neutrinos:

Gamma-ray burst and fireball shock model

Dominance: pion decay or neutron beta decay?

For different parameters in GRBs:

Result:

### Summary and Conclusion



## Purpose of this research

Predict the neutrino flux...

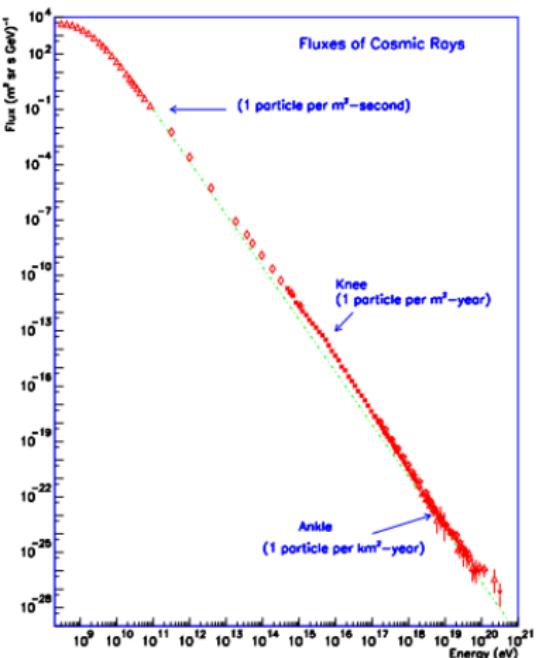
The flavor composition...

Searching for a pure neutron beam source...

- From typical gamma-ray bursts
- The magnetic field effects
- The dependence on measured photon spectrum



# Ultra high energy cosmic rays (UHECR)

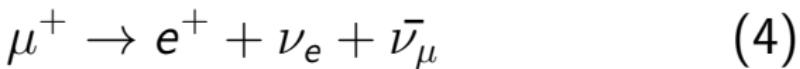
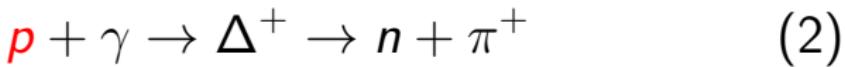
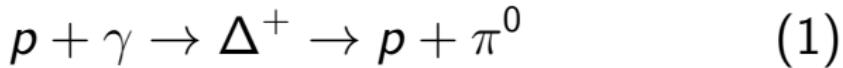


- Power law spectrum  $\frac{dl_{ob}}{d\mathcal{E}} \propto \mathcal{E}^{-2.7}$
- Energy up to  $\mathcal{E} > 10^{20}$  eV
- A cut-off at  $\mathcal{E} \approx 5 \times 10^{19}$  eV
- Fermi (shock) acceleration model:  
Initially:  $\frac{dl}{d\mathcal{E}} \propto \mathcal{E}^{-\alpha}$ ,  $\alpha = 2.0 - 2.4$

[Olinto (2005)]



## Production of neutrinos

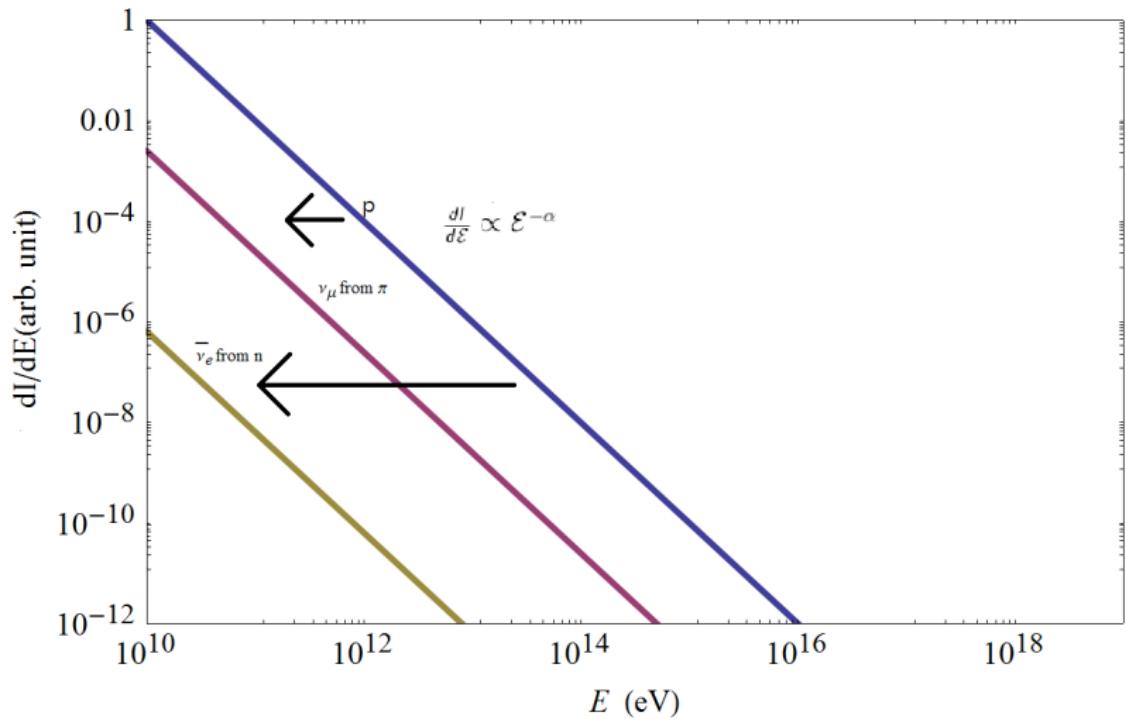


$$\mathcal{E}_{\nu_\mu} \approx 0.05 \mathcal{E}_p$$



$$\mathcal{E}_{\nu_e} \approx 0.0008 \mathcal{E}_p$$





# Possible source candidates of UHECR

- Gamma-ray Burst (GRB)
- Active Galactic Nucleus (AGN)

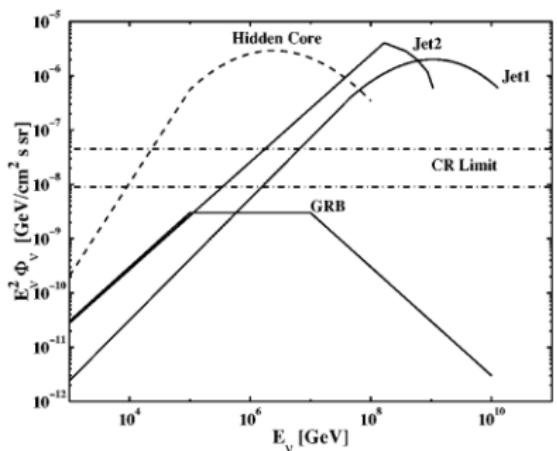


## Gamma-ray burst

- Gamma-ray burst (GRB): a possible astrophysical source of UHECR
- (Internal) shock:  $\Gamma \gtrsim 100$
- Protons(nuclei) are accelerated in the shock
- Photons from synchrotron emission by  $e^\pm$



## The cosmic ray limit



- Normalization from UHECR flux
- AGN-jets model: disfavored
- Gamma-ray burst: a more probable source

[Waxman and Bahcall (1999)]



# Relativistic Fireball Shock model

Gamma-ray burst:

- Massive objects collapse or merger
- Particles ejected relativistically
- (Internal) shock:  $\Gamma \gtrsim 100$
- Protons(nuclei) are accelerated in the shock
- Photon from synchrotron emission by  $e^\pm$



## Total energy of internal shock

To determine  $B'$ ....

- $L_\gamma : 10^{50} - 10^{54} \text{ erg/s}$
- $L_\gamma = \xi_e L_{int},$
- $\frac{2\pi r_i^2 c \Gamma^2 B'^2}{\mu_0} = \xi_B L_{int} = \frac{\xi_B L_\gamma}{\xi_e}$
- $\xi_e = \xi_B = 0.3$
- $r_i \sim 2\Gamma^2 c t_\nu$



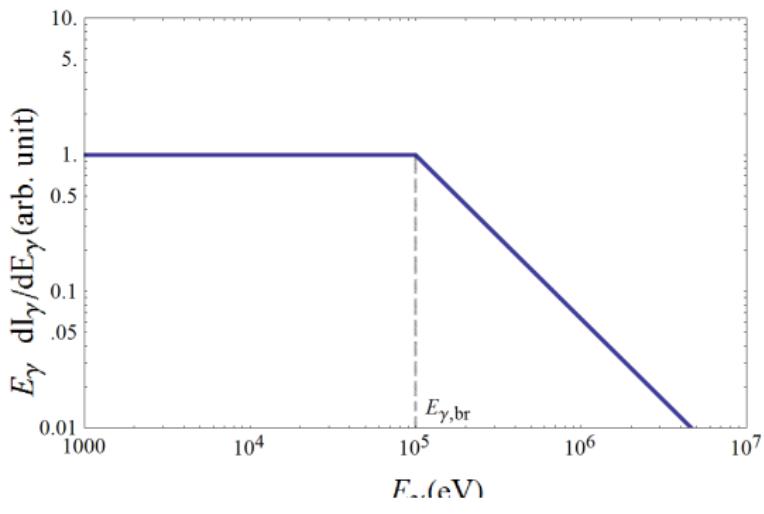
## The effects by magnetic field

- The proton (or nuclei) acceleration
- The synchrotron effects on charged particles
- The photon characteristic energy

$$\mathcal{E}_{\gamma, ch} = \frac{\Gamma \hbar \Gamma_e^2 e B'}{m_e}$$



# The photon spectrum of GRB



- $p + \gamma \rightarrow \Delta^+ \rightarrow n + \pi^+$
- $\mathcal{E}_\gamma \frac{dI_\gamma}{d\mathcal{E}_\gamma} \propto \mathcal{E}_\gamma^{-\alpha_1}$   
if  $\mathcal{E}_\gamma \leq \mathcal{E}_{\gamma,br}$
- $\mathcal{E}_\gamma \frac{dI_\gamma}{d\mathcal{E}_\gamma} \propto \mathcal{E}_\gamma^{-\alpha_2}$   
if  $\mathcal{E}_\gamma \geq \mathcal{E}_{\gamma,br}$
- $\alpha_1 \approx 0, \alpha_2 \approx 1 - 4$
- $\mathcal{E}_{\gamma,br} = \mathcal{E}_{\gamma,ch}$



## To determine the neutrino spectrum...

Take proton as the primary cosmic ray

1. Consider the magnetic field effect by comparing time scales:
2. Pions may lose energy before decay in magnetic field, Neutrons do not.
3. Comparing time scales

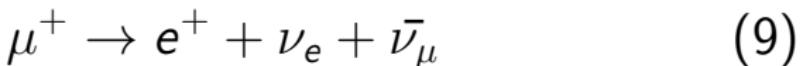
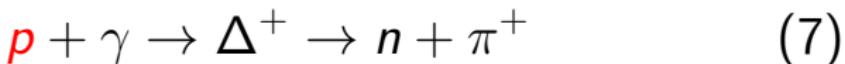
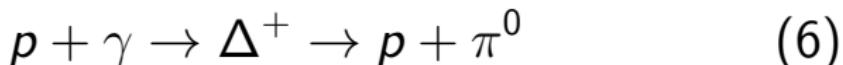
$$t_{acc,p} == t_{sync,p} \Rightarrow \mathcal{E}_{p,max}$$

$$t_{sync,\pi} == t_{decay,\pi} \Rightarrow \mathcal{E}_{\pi,sync}$$

4. Consider the photon spectrum indices



## Production of neutrinos



$$\mathcal{E}_{\nu_\mu} \approx 0.05 \mathcal{E}_p$$



$$\mathcal{E}_{\nu_e} \approx 0.0008 \mathcal{E}_p$$



## Time scales related to $B'$

Consider a charged particle (mass  $m$ , charge  $q$ ) in a uniform magnetic field  $B'$ (in GRB shock frame):

- 

$$t'_{acc} = \frac{\mathcal{E}'}{c^2 q B'}$$

- 

$$t'_{sync} = \frac{9\pi\epsilon_0 m^4 c^5}{q^4 B'^2 \mathcal{E}'}$$

- 

$$t'_{deacy} = \mathcal{E}' \tau / \mathcal{E}_0$$

$$dI_p/d\mathcal{E}_p \propto \mathcal{E}_p^{-2} \text{ if } \mathcal{E}_p < \mathcal{E}_{p,max}$$



## The proton spectrum

$$\frac{dI_p}{d\mathcal{E}_p}(\mathcal{E}_p) \propto \begin{cases} \mathcal{E}_p^{-2} & \text{if } \mathcal{E}_p \leq \mathcal{E}_{p,\max} \\ \exp(-\frac{\mathcal{E}_p}{\mathcal{E}_{p,\max}}) & \text{if } \mathcal{E}_p \geq \mathcal{E}_{p,\max} \end{cases}$$

- By Fermi acceleration:
- The maximal energy is derived by comparing time-scales
- $\mathcal{E}_{p,\max} \propto B'^{-1}$



# The pion source spectrum



$$\mathcal{E}_{\gamma,br} = \frac{\Gamma \hbar \Gamma_e^2 e B'}{m_e}$$

$$\mathcal{E}_{p,br} = \Gamma^2 \frac{m_\Delta^2 c^4 - m_p^2 c^4}{4 \mathcal{E}_{\gamma,br}}$$

$$\mathcal{E}_{\pi,br} \approx 0.2 \mathcal{E}_{p,br}$$



# The pion source spectrum

$$\phi_\pi(\mathcal{E}_\pi) \propto \begin{cases} \mathcal{E}_\pi^{\alpha_2} \mathcal{E}_\pi^{-2} & \text{if } \mathcal{E}_\pi \leq \mathcal{E}_{\pi,br} \\ \mathcal{E}_\pi^{\alpha_1} \mathcal{E}_\pi^{-2} & \text{if } \mathcal{E}_{\pi,br} \leq \mathcal{E}_\pi \leq \mathcal{E}_{\pi,max} \\ \exp(-\frac{\mathcal{E}_\pi}{\mathcal{E}_{\pi,max}}) & \text{if } \mathcal{E}_\pi \geq \mathcal{E}_{\pi,max} \end{cases}$$

Muon-neutrino from pion:

$$\frac{dl_{\nu_\mu}}{d\mathcal{E}_{\nu_\mu}}(\mathcal{E}_{\nu_\mu}) = -\frac{\partial}{\partial \mathcal{E}_{\nu_\mu}} \int_{4\mathcal{E}_{\nu_\mu}}^{\infty} d\mathcal{E}_i \phi_\pi(\mathcal{E}_i) P_\pi(\mathcal{E}_i, 4\mathcal{E}_{\nu_\mu})$$

where

$$P(\mathcal{E}_i, \mathcal{E}_f) = 1 - \exp \left( -\mathcal{E}_{\pi, sync}^2 \frac{(\mathcal{E}_f^{-2} - \mathcal{E}_i^{-2})}{2} \right)$$

$$\mathcal{E}_{\pi, sync}^2 \propto B'^{-2}$$



## The neutron source spectrum



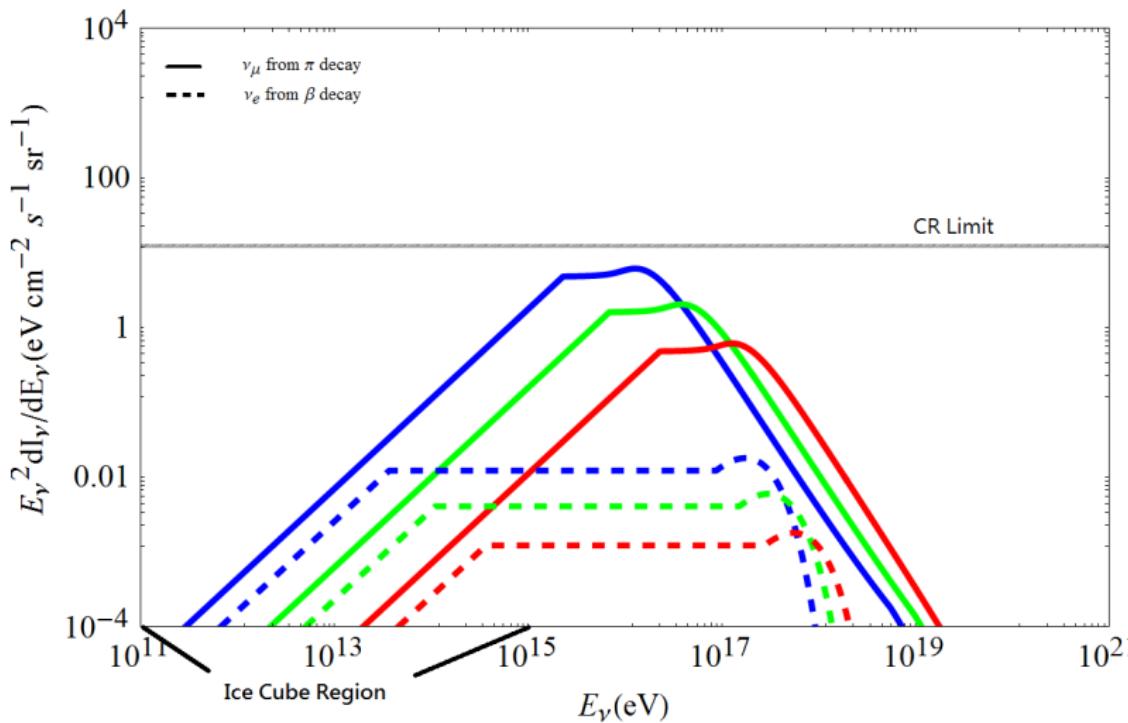
- Neutron not affected by  $B'$
- Constrained by  $\mathcal{E}_{p,max}$

$$\phi_{\nu_e} = \frac{dI_{\nu_e}}{d\mathcal{E}_{\nu_e}} \propto \begin{cases} \mathcal{E}_{\nu_e}^{\alpha_2} \mathcal{E}_{\nu_e}^{-2} & \text{if } \mathcal{E}_{\nu_e} \leq \mathcal{E}_{\nu_e,br} \\ \mathcal{E}_{\nu_e}^{\alpha_1} \mathcal{E}_{\nu_e}^{-2} & \text{if } \mathcal{E}_{\nu_e,br} \leq \mathcal{E}_{\nu_e} \leq \mathcal{E}_{\nu_e,max} \\ \exp\left(-\frac{\mathcal{E}_{\nu_e}}{\mathcal{E}_{\nu_e,max}}\right) & \text{if } \mathcal{E}_{\nu_e} \geq \mathcal{E}_{\nu_e,max} \end{cases}$$

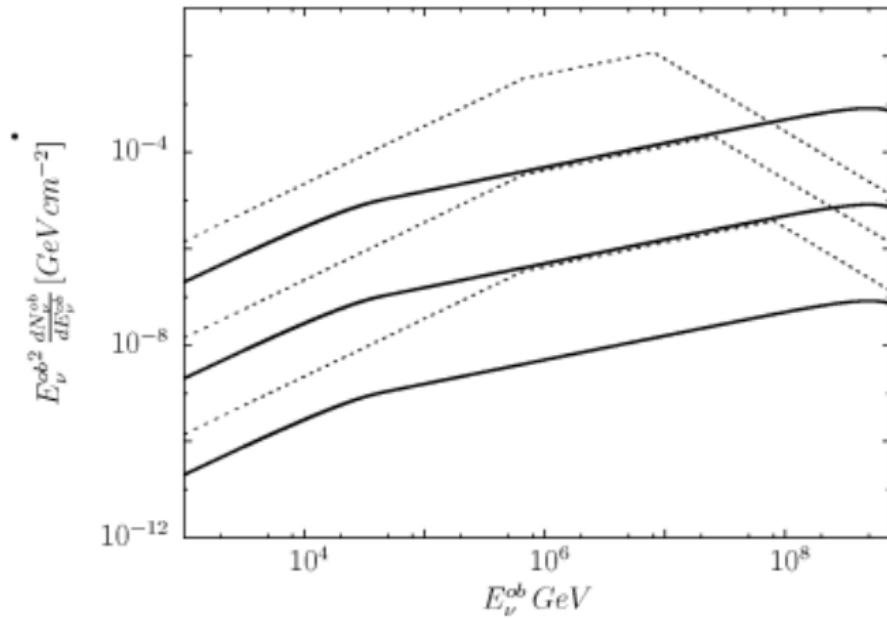


# Typical GRB

$\Gamma=300$ ,  $B=3, 10, 30$  Tesla,  $\alpha_2=1.2$



## Typical GRB: a comparison

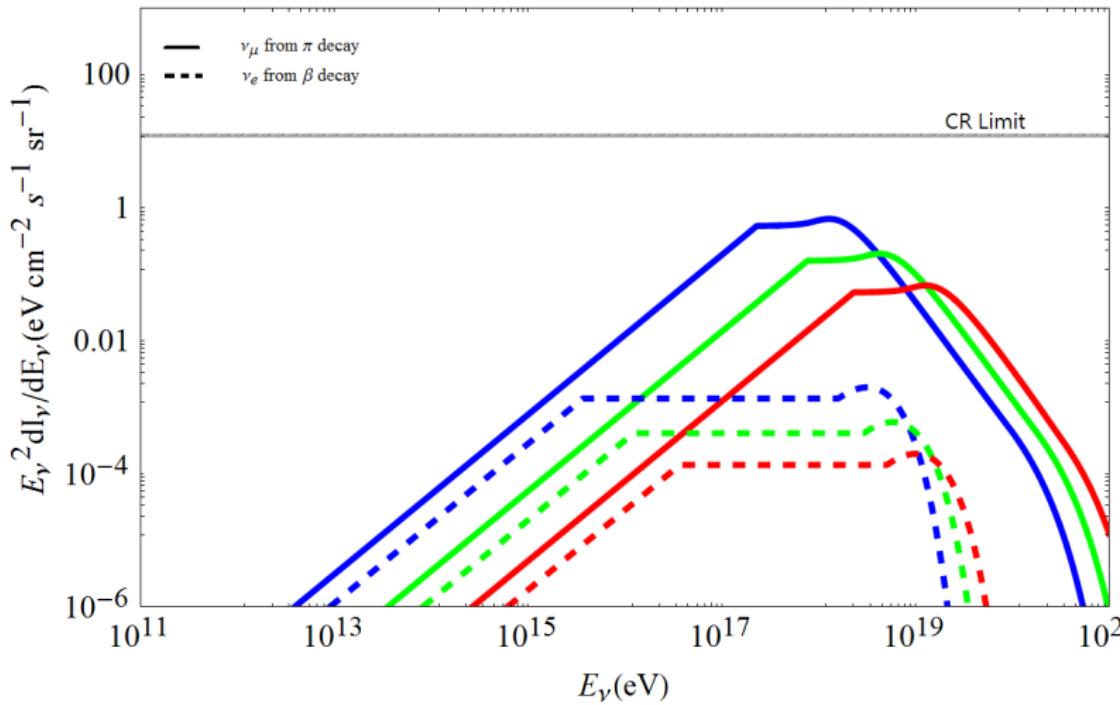


[Moharana and Gupta (2010)]



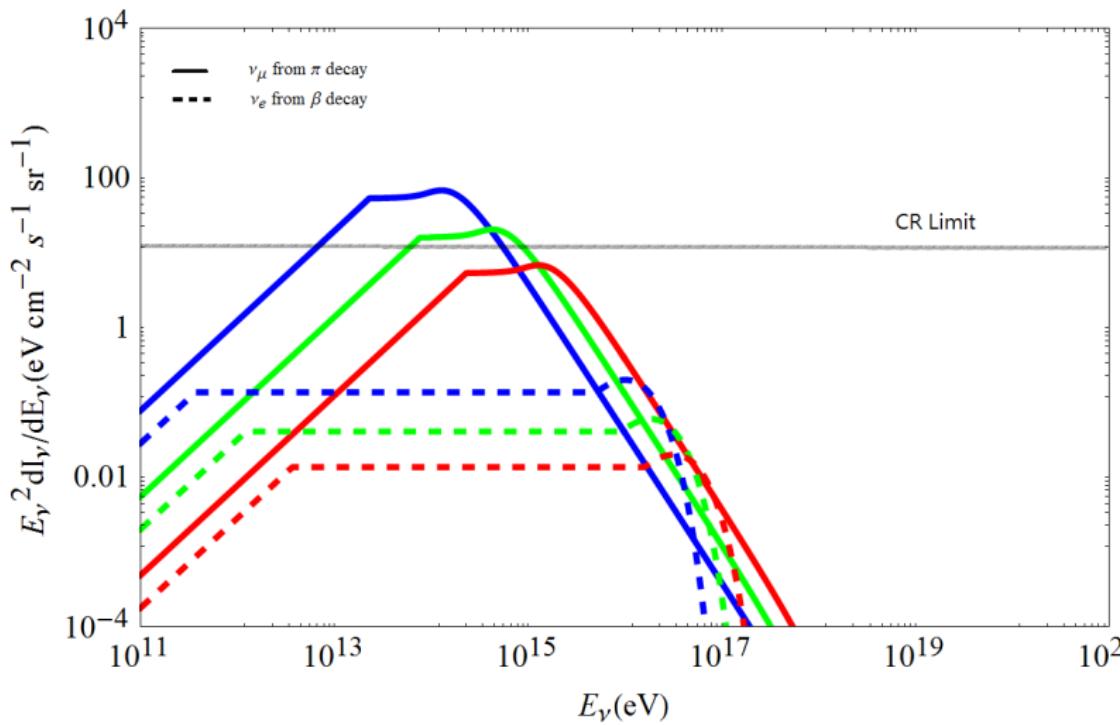
# Weak Magnetic field

$\Gamma=1000$ ,  $B=0.1, 0.3, 1$  Tesla,  $\alpha_2=1.2$



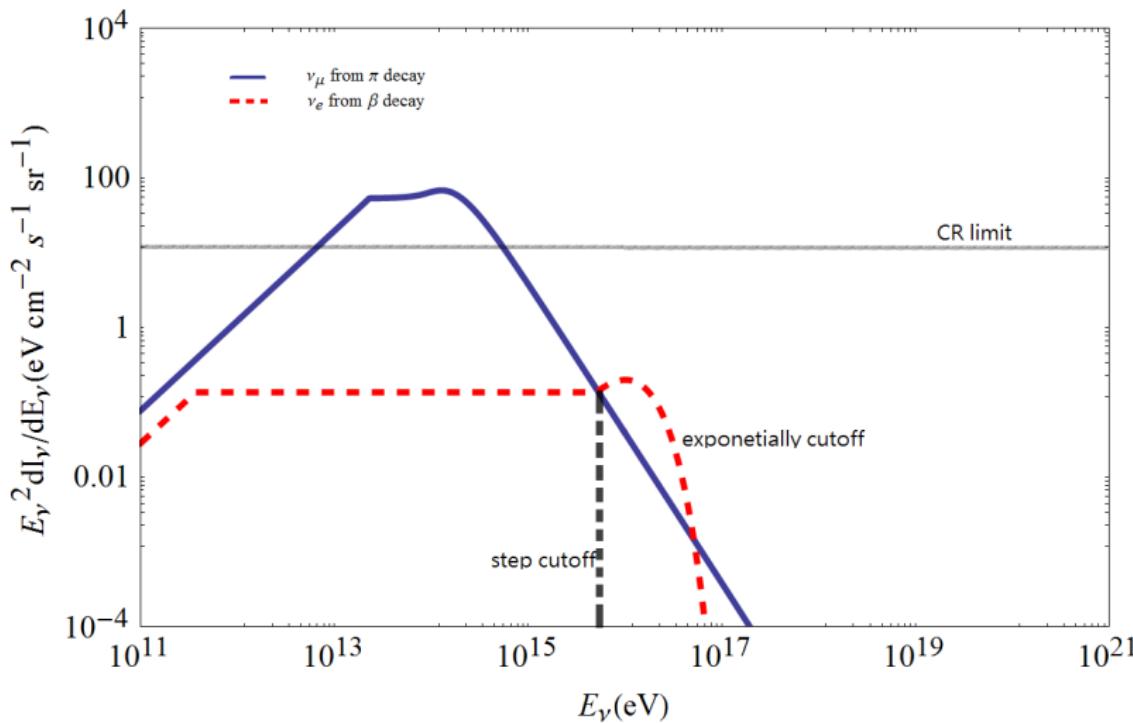
# Strong Magnetic field

$\Gamma=100$ ,  $B=100, 300, 1000$  Tesla,  $\alpha_2=1.2$



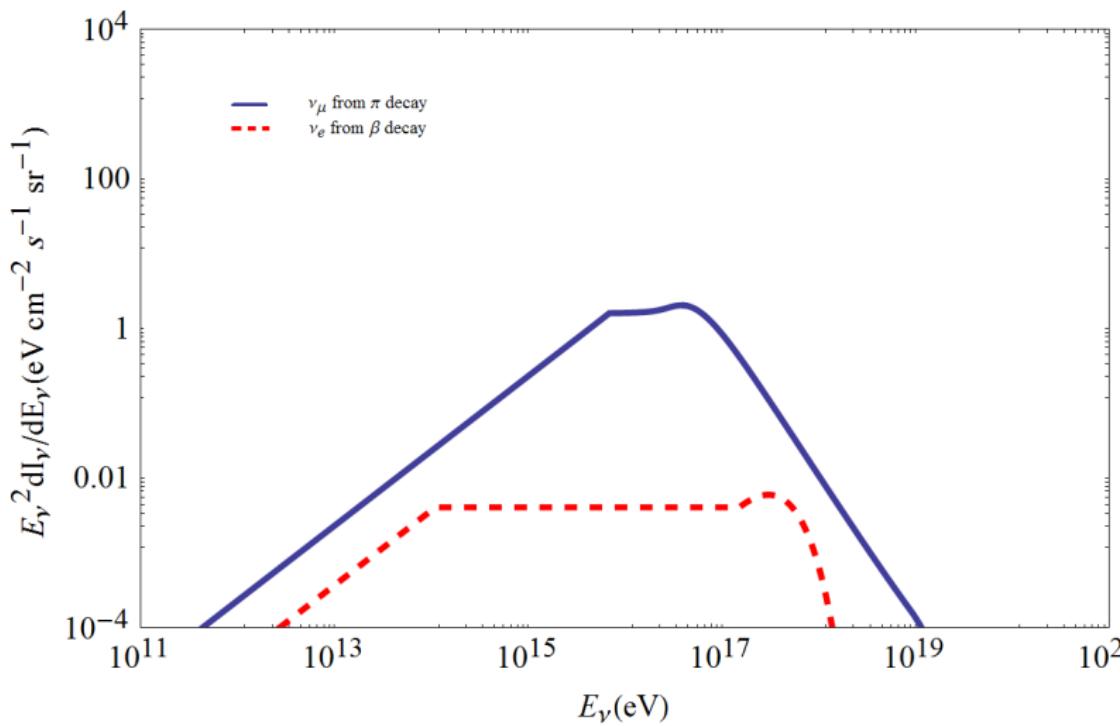
## Cut-off form

$$\Gamma=100, B=1000\text{Tesla}, \alpha_2=1.2$$



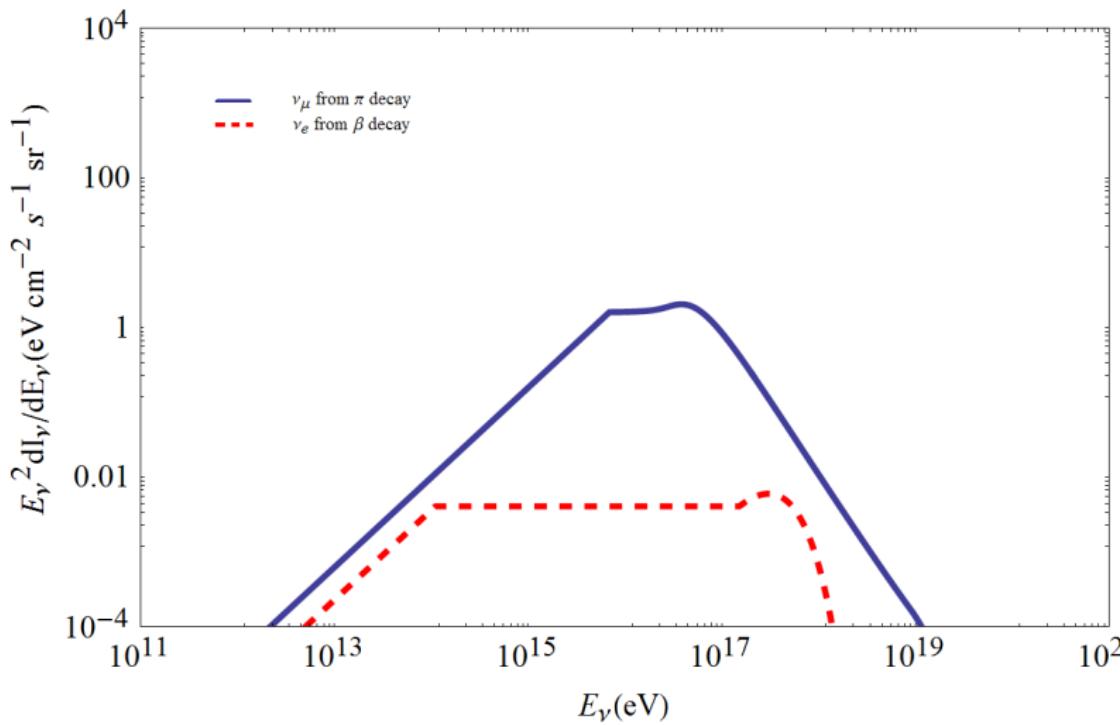
# Dependence on $\alpha_2$

$\Gamma=300$ ,  $B=10\text{Tesla}$ ,  $\alpha_2=1$ .



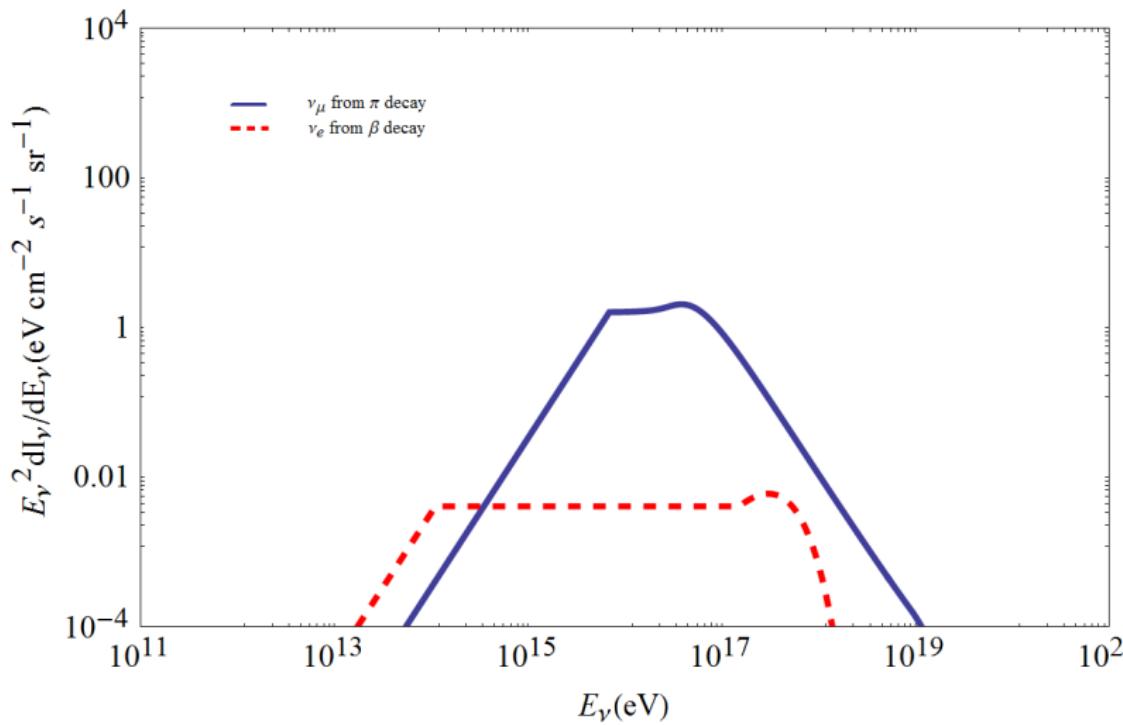
# Dependence on $\alpha_2$

$\Gamma=300$ ,  $B=10\text{Tesla}$ ,  $\alpha_2=1.2$



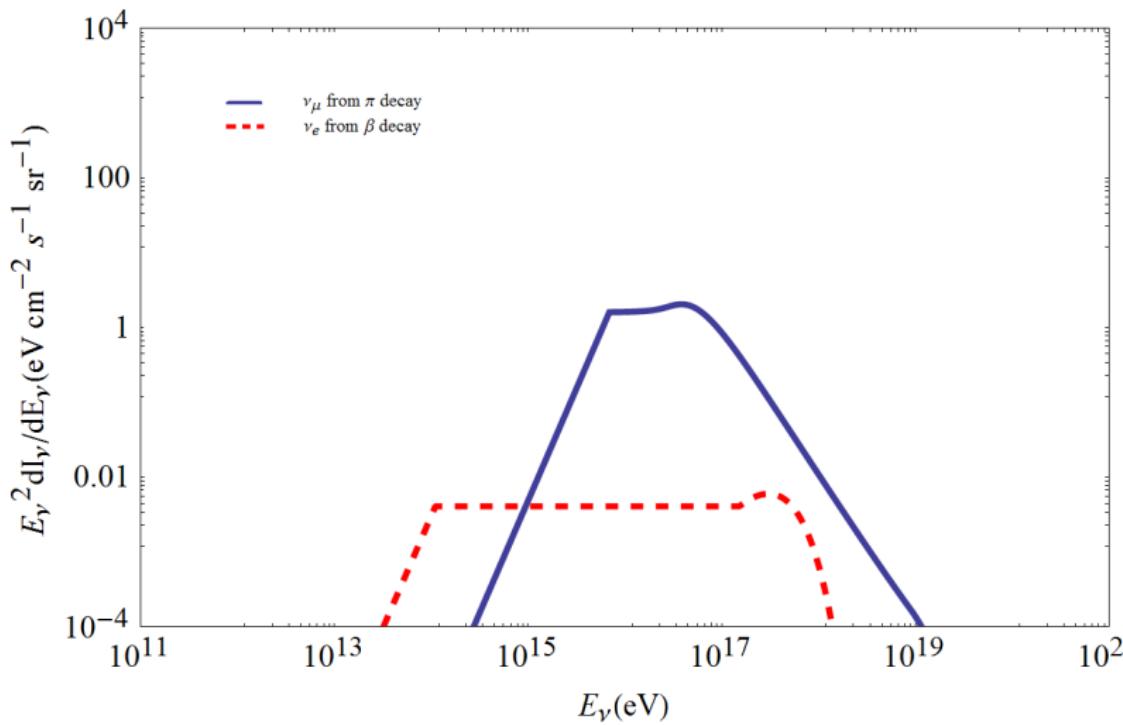
# Dependence on $\alpha_2$

$\Gamma=300$ ,  $B=10\text{Tesla}$ ,  $\alpha_2=2$ .



# Dependence on $\alpha_2$

$\Gamma=300$ ,  $B=10\text{Tesla}$ ,  $\alpha_2=3$ .



## Summary

- The expected flux:(Handout)
- Typical GRB: A different result with [Moharana and Gupta](Does neutron beam exist?)
- GRB with Strong magnetic field:
  1. Depends on the cut-off form
  2. Exceed the cosmic-ray limit
- Flavor dependence on  $\alpha_2$ ...



## Conclusion

### 1. The pure neutron beam

$$(\phi_{\nu_e} : \phi_{\nu_\mu} : \phi_{\nu_\tau} = 1 : 0 : 0) \dots$$

- Not from Typical GRB.
- Not likely from a GRB with Strong magnetic field too
- In the low energy region for some gamma-ray burst, depending on the photon-spectrum index  $\alpha_2$

### 2. Neutrino flavor composition before oscillation

from typical GRB:  $\phi_{\nu_e} : \phi_{\nu_\mu} : \phi_{\nu_\tau} \approx 1 : 1.5 : 0$  in IceCube region ( $10^{11}$  eV  $\leq \mathcal{E}_\nu \leq 10^{15}$  eV)



## Future works

1. The red-shift correction
2. Synchrotron effect on muons
3. Nuclei dominance of UHECR origin



Thank you for your attention

