

### 3. Tensor analysis in special relativity

#### 3.1 Def of tensors

Def.: An abstract definition of a tensor is a <sup>linear</sup> map from the direct product

vector space  $V \otimes V \otimes \dots \otimes V \otimes \tilde{V} \otimes \dots \otimes \tilde{V}$  to real (or complex number);

$$T: \underbrace{V \otimes V \otimes \dots \otimes V}_{N} \otimes \underbrace{\tilde{V} \otimes \dots \otimes \tilde{V}}_{M} \rightarrow \mathbb{C} \quad (3.1.1)$$

This is called a tensor of type  $(N, M)$ .

Here  $V$  = Vector space of Contravariant vector

$\tilde{V}$  = " " " Covariant vector or 1-form. = Dual vector space of  $V$

Elements in  $\tilde{V}$  is called a dual vector.

- By def., a dual vector is a linear map that maps a vector to a scalar.

$$\tilde{v} \in \tilde{V} = V \rightarrow \mathbb{C} \quad (3.1.2)$$

$$\tilde{v} \mapsto \tilde{v}(v)$$

The mapping  $\tilde{v}(v)$  is bilinear. One can show that  $(\tilde{V})^* = V$ .

Useful to denote  $\tilde{v}(v)$  as  $(\tilde{v}, v)$  or  $\langle \tilde{v} | v \rangle$ , called inner product.

Def. The set of all linear functions on  $V$  form a vector space, called the dual vector space  $\tilde{V}$ .

NB. Linearity means  $f(\alpha \vec{v}_1 + \beta \vec{v}_2) = \alpha f(\vec{v}_1) + \beta f(\vec{v}_2)$ .

example: The "metric tensor" defined by

$$g(\vec{A}, \vec{B}) = \vec{A} \cdot \vec{B} = A^\mu g_{\mu\nu} B^\nu \quad \begin{matrix} \leftarrow \text{Components of metric tensor} \\ \uparrow \text{scalar product} \end{matrix} \quad (3.1.3)$$

is a tensor of type  $(0, 2)$ .

Def. The components of a tensor of type  $(M, N)$  <sup>in a frame  $\mathcal{O}$</sup>  are the values of the tensor when evaluated on the basis vectors  $\{\vec{e}_\alpha\}$  and the dual basis vectors  $\{\tilde{\omega}^\alpha\}$  of the frame  $\mathcal{O}$ .

N.B. Tensor is coord/ indep. <sub>frame</sub>, but its components are coord./frame dependent.

N.B. We will use the symbol  $\sim$  to denote a 1-form.

### 3.2 Some more details on 1-forms : $(1, 0)$ tensor

Consider  $\vec{A}$  vector,  $\tilde{p}$  1-form, and  $\{\vec{e}_\alpha\}$  the basis of a frame  $\mathcal{O}$ .

The components  $A^\alpha$  of  $\vec{A}$  wrt  $\mathcal{O}$  is defined by  $\vec{A} \triangleq A^\alpha \vec{e}_\alpha$ . (3.2.1)

The components  $\tilde{p}_\alpha$  of  $\tilde{p}$  wrt  $\mathcal{O}$  is defined by  $\tilde{p}(\vec{e}_\alpha) \triangleq \tilde{p}_\alpha$ .

Since the set of all 1-form is a vector space  $\tilde{V}$ , we can <sup>pick</sup> ~~choose~~ any set of four linearly independent 1-forms as a basis, denote it as  $\{\tilde{\omega}^\alpha\}$  s.t any  $\tilde{p}$  can be expanded as

However, suppose we have a basis  $\{\vec{e}_\alpha\}$  already for the frame  $\mathcal{O}$ .  $\tilde{p} = p_\alpha \tilde{\omega}^\alpha$

It is natural to consider a dual basis s.t.

$$\tilde{\omega}^\alpha(\vec{e}_\beta) = \delta_\beta^\alpha \quad (3.2.2)$$

This implies that for any,  $\tilde{p} = p_\alpha \tilde{\omega}^\alpha$

$$\vec{A} = A^\alpha \vec{e}_\alpha$$

$$\text{we have } \tilde{p}(\vec{A}) = p_\alpha A^\beta \tilde{\omega}^\alpha(\vec{e}_\beta) = p_\alpha A^\alpha \quad (3.2.3)$$

a contradiction!

Under a coord change,  $\vec{e}'_\beta = \vec{e}_\alpha \Lambda^\alpha_\beta$  (3.2.4)

$$\textcircled{O} \rightarrow \textcircled{O}'$$

This implies the transformation

$$\tilde{\omega}'^\alpha = \Lambda^\alpha_\beta \tilde{\omega}^\beta \quad (3.2.5)$$

for the dual basis.

### • Picture of 1-form

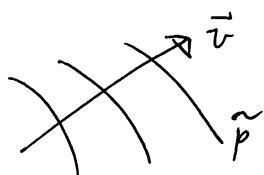
For vector, we usually represent it with an arrow, whose length corresponds to the magnitude of the vector.



For 1-form, we can represent it by a series of surfaces.



such that  $\tilde{p}(\vec{v})$  is the number of surfaces that the arrow  $\vec{v}$  crosses.



Note that a 1-form does not define a unique direction since it is not a vector.

An example of a one-form is the gradient of a scalar function  $\phi$ , defined by

components  $(\tilde{d}\phi)_\alpha \stackrel{\Delta}{=} \partial_\alpha \phi = \frac{\partial \phi}{\partial x^\alpha}$

when evaluated on the tangent vector  $v^\alpha \stackrel{\Delta}{=} \frac{dx^\alpha}{d\tau}$  of a worldline  $x^\alpha(\tau)$ ,

we have  $\vec{v}(\tilde{d}\phi) = v^\alpha \partial_\alpha \phi = \frac{dx^\alpha}{d\tau} \frac{\partial \phi}{\partial x^\alpha} = \frac{d}{d\tau} \phi(x^\alpha(\tau))$

= rate of change of the function  $\phi$  along the curve  $x^\alpha(\tau)$ .

- It is sometimes useful to talk about a normal 1-form.

Given a surface, a normal vector is one which is orthogonal to the surface.  
 need a metric to define!

$$\vec{N} \cdot \vec{T} = 0 \quad \forall \vec{T} \text{ tangential to the surface.}$$

on the other hand, to define a normal 1-form, there is no need of a metric:

$$\tilde{\omega}(\vec{T}) = 0 \quad \forall \vec{T} \text{ tangential to the surface.}$$

Moreover we say a normal one-form is outward if

$$\tilde{\omega}(\vec{v}) > 0 \quad \text{for all vectors } \vec{v} \text{ which point outwards from the surface}$$

### 3.3 Some more details of the $(^0_2)$ tensors

We have seen that the metric tensor is a  $(^0_2)$  tensor.

Another example is to simply take product of two 1-forms:  $\tilde{p} \otimes \tilde{q}$ .

$$\tilde{p} \otimes \tilde{q} (\vec{A}, \vec{B}) \triangleq \tilde{p}(\vec{A}) \tilde{q}(\vec{B}) \quad (3.3.1)$$

(called Cartesian product  
outer)

Note that  $\tilde{p} \otimes \tilde{q} + \tilde{q} \otimes \tilde{p}$  are different  $(^0_2)$  tensors.

Thm. A general  $(^0_2)$  tensor can be written as a sum of outer products.

Pf. Let  $\tilde{p}$  be a  $(^0_2)$  tensor with components  $p_{\alpha\beta} \triangleq \tilde{p}(\vec{e}_\alpha, \vec{e}_\beta)$ .

On arbitrary vectors  $\vec{A}, \vec{B}$ , we have  $\tilde{p}(\vec{A}, \vec{B}) = p_{\alpha\beta} A^\alpha B^\beta$ .

The question is how is  $\tilde{p} + p_{\alpha\beta}$  related?

~~to see~~

In general we can assume that  $\tilde{p} = p_{\alpha\beta} \tilde{\omega}^{\alpha\beta}$  for some tensor  $\tilde{\omega}^{\alpha\beta}$ .

$$\text{then } \tilde{p}(\vec{A}, \vec{B}) = p_{\alpha\beta} \tilde{\omega}^{\alpha\beta}(\vec{A}, \vec{B}) \stackrel{!}{=} p_{\alpha\beta} A^\alpha B^\beta$$

$$p_{\alpha\beta} A^\gamma B^\delta \tilde{\omega}^{\alpha\beta}(\vec{e}_\gamma, \vec{e}_\delta)$$

This holds for arb  $A^\alpha, B^\beta$ , so need  $\tilde{\omega}^{\alpha\beta}(\vec{e}_\gamma, \vec{e}_\delta) = \delta_\gamma^\alpha \delta_\delta^\beta$

$\therefore \tilde{\omega}^{\alpha\beta}$  is a tensor whose value is just the product of two basis one-forms:

$$\tilde{\omega}^{\alpha\beta} = \tilde{\omega}^\alpha \otimes \tilde{\omega}^\beta$$

$\therefore$  the tensors  $\tilde{\omega}^\alpha \otimes \tilde{\omega}^\beta$  are a basis for all  $(^0_2)$  tensors and

$$\tilde{p} = p_{\alpha\beta} \tilde{\omega}^\alpha \otimes \tilde{\omega}^\beta$$

We call a  $(^0_2)$  tensor symmetric if  $\tilde{p}(\vec{A}, \vec{B}) = \tilde{p}(\vec{B}, \vec{A}) \quad \forall \vec{A}, \vec{B}$

$$\text{In terms of components, } p_{\alpha\beta} = p_{\beta\alpha} \quad (3.3.2)$$

Similarly,  $\tilde{p}$  is antisymmetric if  $\tilde{p}(\vec{A}, \vec{B}) = -\tilde{p}(\vec{B}, \vec{A}) \quad \forall \vec{A}, \vec{B}$

$$p_{\alpha\beta} = -p_{\beta\alpha} \quad (3.3.3)$$

In general, given a  $(^0_2)$  tensor, it may not be symmetric or antisymmetric,

but we can always define its symmetric and antisymmetric parts as

$$\left\{ \begin{array}{l} h_{\alpha\beta}^{(S)} \triangleq \frac{1}{2}(h_{\alpha\beta} + h_{\beta\alpha}) \\ h_{\alpha\beta}^{(AS)} \triangleq \frac{1}{2}(h_{\alpha\beta} - h_{\beta\alpha}) \end{array} \right. \quad (3.3.4)$$

$$\text{Note that } h_{\alpha\beta} = h_{\alpha\beta}^{(S)} + h_{\alpha\beta}^{(AS)}$$

The factor  $\frac{1}{2}$  in  $h_{\alpha\beta}^{(S)}$  is defined s.t. for symmetric tensor, it is equal to its symmetric part.

### 3.4 Identification of vectors & 1-forms using $(^0_2)$ or $(^2_0)$ tensors

39 36

- In general, vectors & 1-forms are dual to each other in the sense of linear vector spaces. In the presence of a  $(^0_2)$  tensor, one can map vector to 1-form or vice-versa directly.

Let  $g$  be a  $(^0_2)$  tensor, it has two arguments:  $g(\cdot, \cdot)$

If we fix  $\vec{v} \in V$ , then  $g(\vec{v}, \cdot)$  is a functional on  $V$  and we may call this 1-form as  $\tilde{v}$ :

$$\tilde{v}(\vec{\omega}) \triangleq g(\vec{v}, \vec{\omega}) \quad (3.4.1)$$

Components of  $\tilde{v}$  can be worked out as follows:

$$v_\alpha \triangleq \tilde{v}(\vec{e}_\alpha) = g(\vec{v}, \vec{e}_\alpha) = v^\beta g(\vec{e}_\beta, \vec{e}_\alpha) \quad \vec{v} = v^\beta \vec{e}_\beta,$$

$$\therefore v_\alpha = v^\beta g_{\beta\alpha} \quad (3.4.2)$$

- Inversely, we could also go from a 1-form to a vector.

Let  $\tilde{v} \in \tilde{V}$  be a 1-form, and  $\tilde{g}$  a  $(^2_0)$  tensor,

then  $\tilde{g}(\tilde{v}, \cdot)$  is a linear functional on  $\tilde{V}$ .

Denote this by  $\tilde{v}$ , then

$$\tilde{v}(\tilde{\omega}) \triangleq \tilde{g}(\tilde{v}, \tilde{\omega}) \quad (3.4.3)$$

In terms of components,

$$v^\alpha = g^{\alpha\beta} v_\beta \quad \text{where } g^{\alpha\beta} \triangleq \tilde{g}(w^\alpha, w^\beta) \quad (3.4.4)$$

In general,  $\tilde{g}, \tilde{\tilde{g}}$  are tensors of type  $(\frac{0}{2})$  and  $(\frac{2}{0})$ .

In Specific problem, the system may admit a metric



a  $(\frac{0}{2})$  tensor which is non-degenerate.

In this case, as  $g$  is non-degenerate, its inverse exists and is a  $(\frac{2}{0})$  tensor.

In components, we denote the metric as  $g_{\alpha\beta}$  and the inverse as  $g^{\alpha\beta}$ .

### 3.5 Basic manipulations of tensors

Refer to P-26 & P-27 (§2.3)