Doubly Charged Higgs Bosons at Hadron Colliders

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Abstract:

Ongoing searches for doubly charged Higgs bosons $(H^{\pm\pm})$ at the Tevatron assume that $H^{\pm\pm}$ are produced in pairs via $q\bar{q} \rightarrow H^{++}H^{--}$. We show that single $H^{\pm\pm}$ production via $q\bar{q} \rightarrow H^{\pm\pm}H^{\mp}$ can be comparable in size and consider its impact on present and future searches at Hadron Colliders. A.A and Mayumi Aoki (KEK), Phys. Rev. D72 (2005) 035011

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Outline

- Imminent start of Large Hadron Collider (LHC)
- Neutrino mass and mixing
- TeV scale mechanisms for neutrino mass generation
- Higgs Triplet Model and Left-Right Symmetric Model
- $H_L^{\pm\pm}$ phenomenology at hadron colliders

Large Hadron Collider

LHC (CERN) due to commence operation in summer 2007

- Proton-Proton collisions at $\sqrt{s} = 14$ TeV
- Highest energy collider ever built
- ATLAS and CMS optimized for Higgs boson search
- New Physics discovery potential up to TeV scale

Strong evidence for neutrino masses and mixings from both terrestrial and celestial sources Mixing angles are being probed by oscillation experiments: i) Atmospheric angle almost maximal: $\sin^2 \theta_{23} \sim 1$ ii) Solar angle close to maximal: $\sin^2 \theta_{12} \sim 0.8$ iii) Reactor angle not measured: $\sin^2 \theta_{13} < 0.16$ θ_{23} and θ_{12} much larger than mixing angles in the quark sector

Oscillation experiments also sensitive to neutrino mass differences:

i) $\Delta M_{atm}^2 \sim 10^{-3} eV^2 \ (=M)$ ii) $\Delta M_{sol}^2 \sim 10^{-5} eV^2 \ (=m)$

Much lighter than other fermion masses ($m_e = 0.5 \text{ MeV}$) Three possibilities:

- 1) Hierarchical $\sim (0, m, M)$.
- 2) Inverted hierarchical $\sim (M, M + m, 0)$
- 3) Quasi-degenerate $\sim (M, M, M)$

Present and Future neutrino experiments (e.g. T2K, MINOS, OPERA) i) Measure Δm^2 , θ_{12} , θ_{23} with higher precision ii) Possible first measurement of θ_{13} , CP Phase

Mechanism of mass and mixing?

Models which can be probed at the Tevatron/LHC are of *immediate phenomenological interest*

Many models for neutrino mass generation!

Models with a specific signature at High Energy Colliders

(Tevatron/LHC) are phenomenologically appealing

 $H_L^{\pm\pm}$ with coupling to W, Z and leptons

- Two such models are:
- i) Higgs Triplet Model (HTM) Gelmini/Roncadelli 80
- ii) Left-Right Symmetric Model (LR Model) Mohapatra/Senjanov

Higgs Triplet Model (HTM)

SM Lagrangian with one $SU(2)_L$ I = 1, Y = 2 Higgs triplet

$$\Delta = \begin{pmatrix} H^+/\sqrt{2} & H^{++} \\ H^0 & -H^+/\sqrt{2} \end{pmatrix}$$

Higgs potential:

$$V = m^{2}(\Phi^{\dagger}\Phi) + \lambda_{1}(\Phi^{\dagger}\Phi)^{2} + M^{2}\mathrm{Tr}(\Delta^{\dagger}\Delta)$$

$$+\lambda_i (\text{quartic terms}) + \frac{1}{\sqrt{2}}\mu(\Phi^T i\tau_2 \Delta^{\dagger} \Phi) + h.c$$

Triplet vacuum expectation value:

$$< H^{0} > = v_{L} \sim \mu / M$$
 (1 $eV < v_{L} < 8GeV$)

Higgs boson spectrum

The HTM has 7 Higgs bosons:

$$H^{\pm\pm}, H^{\pm}, H^{0}, A^{0}, h^{0}$$

In general H^{\pm} , H^{0} , A^{0} eigenstates are mixtures of doublet and triplet fields However, mixing $\sim v_L/v$ and $v_L << v$: i) h^{0} is *SM Higgs boson* (essentially I = 1/2 doublet) ii) $H^{\pm\pm}$ purely triplet and H^{\pm} , H^{0} , A^{0} essentially triplet Higgs boson masses

Close to degenerate with splittings caused by λ_5

$$m_{H^{\pm\pm}}^2 \simeq M^2 + 2 \frac{(\lambda_4 - \lambda_5)}{g^2} M_W^2$$

$$m_{H^{\pm}}^2 \simeq m_{H^{\pm\pm}}^2 + 2 \frac{\lambda_5}{g^2} M_W^2 \quad m_{H^0,A^0}^2 \simeq m_{H^{\pm}}^2 + 2 \frac{\lambda_5}{g^2} M_W^2$$

For $\lambda_5 > 0 \ (\lambda_5 < 0)$:

 $m_{H^{\pm\pm}} < m_{H^{\pm}} < m_{H^{0},A^{0}} \ (m_{H^{\pm\pm}} > m_{H^{\pm}} > m_{H^{0},A^{0}})$

For $H^{\pm\pm}$ in range at LHC require M < 1 TeV.

Neutrino mass in Higgs Triplet Model (HTM)

No right handed neutrino:

Neutrino mass from triplet-lepton-lepton coupling (h_{ij}) :

$$h_{ij}\left[\sqrt{2}\,\overline{l}_i^c P_L l_j H_L^{++} + (\overline{l}_i^c P_L \nu_j + \overline{l}_j^c P_L \nu_i) H_L^{+} - \sqrt{2}\,\overline{\nu}_i^c P_L \nu_j H_L^{0}\right]$$

Light neutrinos receive a Majorana mass:

$$\mathcal{M}_{\nu} \sim v_L h_{ij}$$

$$h_{ij} = \frac{1}{\sqrt{2}v_L} V_{\text{MNS}} diag(m_1, m_2, m_3) V_{\text{MNS}}^T$$

Left-Right Symmetric Model

- Left-Right Symmetric Model Pati/Salam 74 can provide low-energy seesaw mechanism Main differences with SM :
- i) Introduce new gauge group $SU(2)_R$
- ii) Right Handed fermions in doublets (u_R, d_R), (u_R, l_R)
- iii) Existence of ν_R is required by gauge symmetry
- iv) New (heavy) gauge bosons W_R^{\pm}, Z_R
- v) Extended Higgs sector for symmetry breaking
- $SU(2)_R \otimes SU(2)_L \otimes U(1)_{B-L} \rightarrow SU(2)_L \otimes U(1)_Y \rightarrow U(1)_Q$

Virtues of LR model?

1) Arbitrary hypercharge Y has been replaced by theoretically attractive B-L2) Parity conservation of weak interactions i.e. $\psi_L \rightarrow \psi_R$ symmetry restored at energies $> v_R$ and broken spontaneously 3) Low energy seesaw mechanism for neutrino mass can be implemented with appropriate choice of Higgs representations (Isospin I=1 Triplets) Mohapatra/Senjanovic 80 4) New particles (Higgs bosons, heavy RH neutrinos, gauge bosons) at v_R scale possibly accessible to LHC

Higgs boson content

$$\begin{array}{l} \text{bidoublet} : \phi = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix} : <\phi_1^0 > = \kappa_1, <\phi_2^0 > = \kappa_2 \\ \\ \kappa_1^2 + \kappa_2^2 = \kappa^2 = 246^2 GeV^2 \\ \\ \begin{array}{l} \text{triplets} : \ \Delta_{L,R} = \begin{pmatrix} H_{L,R}^+ / \sqrt{2} & H_{L,R}^{+++} \\ H_{L,R}^0 & -H_{L,R}^+ / \sqrt{2} \end{pmatrix} \end{array}$$

$$< H_{L,R}^{0} > = v_{L,R}$$

20 d.o.f give 14 physical scalars (6 neutral, 4 charged,4 doubly charged) and 6 goldstone bosons

Neutrino mass in LR Symmetric Model

Right handed neutrino required by $SU(2)_L \otimes SU(2)_R \otimes U(1)$ Majorana neutrino mass obtained via seesaw mechanism:

$$M_{\nu} = \begin{pmatrix} h_L v_L & M_D \\ M_D^T & h_R v_R \end{pmatrix}, \quad M_D = \frac{1}{\sqrt{2}} \left(h_l \kappa_1 + \tilde{h}_l \kappa_2 \right)$$

If $h_L = h_R = h = \mathcal{O}(1)$, $v_R = 1$ TeV, $v_L = 0$, $M_D \sim \text{MeV}$

$$M_{\nu} = \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix},$$

 \rightarrow obtain neutrino mass (M_D^2/M_R) of eV scale

Summary

- HTM: Majorana mass for light neutrinos ν_L (No v_R).
- LR symmetric model : Majorana mass for ν_L via seesaw mechanism (v_R needed)
- Both models can have $H_L^{\pm\pm}$ with mass < 1 TeV.

Phenomenology of $H_L^{\pm\pm}$ at Hadron Colliders

Summary of $H^{\pm\pm}$ searches

Searches at LEP, HERA, Tevatron:



Production of $H_L^{\pm\pm}$ at Tevatron

First searches at a Hadron collider in 2003 CDF,D0



- $\bullet \ \sigma_{H^{++}H^{--}}$ is a simple function of $M_{H^{\pm\pm}}$
- $\sigma_{H^{++}H^{--}}$ has no dependence on h_{ij}
- $H^{\pm\pm}$ decays via h_{ij} to same charge $ee, e\mu, \mu\mu, e\tau, \mu\tau, \tau\tau$

Search strategy

Many possible searches!

From pair produced $H^{++}H^{--}$ one can select:

2 leptons (2l), 3 leptons (3l), and 4 leptons (4l)

Signal yield $S_{2l} > S_{3l} > S_{4l}$, Background $B_{2l} > B_{3l} > B_{4l}$

For discovery want to maximize S/B

Current status of Tevatron searches:

	ee	$e\mu$	$\mu\mu$	e au	μau
21	< 133 GeV	< 113 GeV	< 136 GeV	Х	X
31				< 114 GeV	< 112 GeV
41				< 114 GeV	< 112 GeV

Tevatron search for $p\overline{p} \to H^{++}H^{--}$, $H^{\pm\pm} \to e^{\pm}e^{\pm}, e^{\pm}\mu^{\pm}, \mu^{\pm}\mu^{\pm}$

Mass limits assume BR $(H^{\pm\pm} \rightarrow l_i^{\pm} l_j^{\pm}) = 100\%$ in a given channel.



Strongest mass limits for any Higgs boson! CDF 03

3*l* and 4*l* searches give essentially same mass limit CDF 06



Single $H^{\pm\pm}$ production via $qq' \rightarrow H^{\pm\pm}H^{\mp}$

Any additional production mechanisms for $H^{\pm\pm}$? For $H_L^{\pm\pm}$ one can exploit the vertex $W^{\pm}H_L^{\pm\pm}H_L^{\mp}$ Dion 98, Gunion 98



- Not possible at e^+e^- colliders
- $H_L^{\pm} \rightarrow l^{\pm} \nu$ decay expected

Impact of
$$H_L^{\pm\pm}H_L^{\mp}$$

Current searches are already sensitive to $qq' \rightarrow H_L^{\pm\pm} H_L^{\pm}!$

- 2*l* search: sensitive to $H_L^{\pm\pm}H_L^{\mp}$ irrespective of H_L^{\pm} decay
- 3*l* search: sensitive to $H_L^{\pm\pm}H_L^{\mp}$ if $H_L^{\pm} \to l^{\pm}\nu$
- 4*l* search: *insensitive* to $H_L^{\pm\pm}H_L^{\mp}$ Define inclusive single $H_L^{\pm\pm}$ cross-section for 2*l*,3*l* search:

$$\sigma_{H^{\pm\pm}} = \sigma(p\overline{p} \to H_L^{\pm\pm} H_L^{--}) + \sigma(p\overline{p} \to H_L^{\pm\pm} H_L^{\mp})$$

Akeroyd, Aoki 05

Inclusive single $H_L^{\pm\pm}$ production at Tevatron

$$\sigma_{H^{\pm\pm}} = \sigma(p\overline{p} \to H_L^{++}H_L^{--}) + 2\sigma(p\overline{p} \to H_L^{++}H_L^{-})$$



Increases search potential of Tevatron in 21,31 channels Akeroyd, Aoki 05

Inclusive single $H_L^{\pm\pm}$ production at LHC



Discovery up to $M_{H^{\pm\pm}} < 1$ TeV for luminosity=300 fb⁻¹ Precision studies possible if light $M_{H^{\pm\pm}}$ (< 250 GeV) found at Tevatron

Decay branching ratios of $H^{\pm\pm}$

Mass limits assume BR $(H^{\pm\pm} \rightarrow l_i l_j)=100\%$ for specific i, jActual BR $(H^{\pm\pm} \rightarrow l_i l_j)$ depends on relative values of h_{ij} LR Model: h_{ij} arbitrary. No prediction for BRs HTM: h_{ij} directly related to neutrino mass matrix

$$h_{ij} = \frac{1}{\sqrt{2}v_L} V_{\text{MNS}} diag(m_1, m_2, m_3) V_{\text{MNS}}^T$$

Prediction for BR($H^{\pm\pm} \rightarrow l_i l_j$) depends on which neutrino mass hierarchy is realized!

Inclusive single
$$\sigma(p\overline{p} \to H_L^{\pm\pm} + X \to e^{\pm}e^{\pm}, e^{\pm}\mu^{\pm}, \mu^{\pm}\mu^{\pm})$$

Normal neutrino mass hierarchy: $\sigma_{\mu\mu} >> \sigma_{ee}, \sigma_{e\mu}$



Inclusive single
$$\sigma(p\overline{p} \to H_L^{\pm\pm} + X \to e^{\pm}e^{\pm}, e^{\pm}\mu^{\pm}, \mu^{\pm}\mu^{\pm})$$

Inverted neutrino mass hierarchy: $\sigma_{ee} > \sigma_{\mu\mu} >> \sigma_{e\mu}$



Inclusive single
$$\sigma(p\overline{p} \to H_L^{\pm\pm} + X \to e^{\pm}e^{\pm}, e^{\pm}\mu^{\pm}, \mu^{\pm}\mu^{\pm})$$

Quasi-degenerate neutrinos: $\sigma_{\mu\mu} \sim \sigma_{ee} >> \sigma_{e\mu}$



Summary of $H^{\pm\pm}$ discovery propapects at Hadron collider

• Tevatron searches are ongoing:

 $0.24 \text{ fb}^{-1}/0.35 \text{ fb}^{-1} \text{ used } (> 1 \text{ fb}^{-1} \text{ accumulated})$

- Tevatron can discover $H^{\pm\pm}$ up to 250 GeV in $ee, e\mu, \mu\mu$ channels (smaller masses in $e\tau, e\mu, \tau\tau$)
- HTM gives predictions for $H^{\pm\pm} \rightarrow l_i l_j$ which are sensitive to the neutrino mass hierarchy
- LHC has i) discovery capability up to 1 TeV or ii) precision measurements of $H^{\pm\pm} \rightarrow l_i l_j$ if $H^{\pm\pm}$ discovered at Tevatron

Conclusions

- Higgs triplets can generate neutrino mass via triplet vev
- Arise in HTM and LR symmetric model
- $H_L^{\pm\pm} \rightarrow l^{\pm}l^{\pm}$ distinctive signal
- Produced at Hadron Colliders via $q\overline{q} \rightarrow H_L^{++}H_L^{--}$
- Inclusion of $qq' \rightarrow H_L^{\pm\pm} H_L^{\pm}$ increases cross-section
- Tevatron have shown interest in including $qq' \rightarrow H_L^{\pm\pm} H_L^{\mp}$

in future searches