

Adding two angular momentum

$|j_1 m_1\rangle \otimes |j_2 m_2\rangle$ can be expanded in terms of $|j_1 j_2; m_1 m_2\rangle$ or $|j_1 j_2; j m\rangle$ basis vector

$$\begin{aligned}
 |j_1 j_2; j m\rangle &= \sum_{m_1 m_2} |j_1 j_2; m_1 m_2\rangle \langle j_1 j_2; m_1 m_2 | j_1 j_2; j m\rangle \\
 &= \sum_{m_1 m_2} C_{m_1 m_2}^{j m} |j_1 j_2; m_1 m_2\rangle
 \end{aligned}$$

↑ Clebsch-Gordan coefficient

or

$$|j_1 j_2 m_1 m_2\rangle = \sum_j C_{m_1 m_2}^{j m} |j_1 j_2; j m\rangle$$

$$j = j_1 + j_2$$

Act J_z on both sides of the equation.

$$\Rightarrow m |j_1 j_2; j m\rangle = \sum_{m_1 m_2} C_{m_1 m_2}^{j m} (m_1 + m_2) |j_1 j_2; m_1 m_2\rangle$$

Times $\langle j_1 j_2 m_1 m_2 |$ on both side

$$\Rightarrow m C_{m_1 m_2}^{j m} = (m_1 + m_2) C_{m_1 m_2}^{j m}$$

$$\Rightarrow \boxed{\text{If } m - m_1 - m_2 \neq 0 \quad C_{m_1 m_2}^{j m} = 0}$$

Apply J_+ on both sides.

$$\sqrt{j(j+1) - m(m+1)} |j_1 j_2; j m+1\rangle$$

$$= \sum_{m_1 m_2} C_{m_1 m_2}^j \left\{ \begin{aligned} &\sqrt{j_1(j_1+1) - m_1(m_1+1)} |j_1 j_2; m_1+1, m_2\rangle \\ &+ \sqrt{j_2(j_2+1) - m_2(m_2+1)} |j_1 j_2; m_1, m_2+1\rangle \end{aligned} \right\}$$

times $\langle j_1, j_2, m_1, m_2 |$ from the left

$$\Rightarrow \sqrt{j_1(j_1+1) - m_1(m_1+1)} C_{m_1, m_2}^{j_1 m_1+1} = \sqrt{j_1(j_1+1) - m_1(m_1-1)} C_{m_1-1, m_2}^{j_1 m_1} + \sqrt{j_2(j_2+1) - m_2(m_2-1)} C_{m_1, m_2-1}^{j_1 m_1}$$

similarly, by acting J_- on the equation, we have

$$\Rightarrow \sqrt{j_1(j_1+1) - m_1(m_1-1)} C_{m_1, m_2}^{j_1 m_1-1} = \sqrt{j_1(j_1+1) - m_1(m_1+1)} C_{m_1+1, m_2}^{j_1 m_1} + \sqrt{j_2(j_2+1) - m_2(m_2+1)} C_{m_1, m_2+1}^{j_1 m_1}$$

Proof of the Wigner-Eckart theorem

* From $[J_z, T_k^g] = g T_k^g$

$$\langle \alpha'; j', m' | J_z T_k^g | \alpha; j, m \rangle = \langle \alpha'; j', m' | T_k^g J_z | \alpha; j, m \rangle = g \langle \alpha'; j', m' | T_k^g | \alpha; j, m \rangle$$

$$\Rightarrow (m' - m - g) \langle \alpha'; j', m' | T_k^g | \alpha; j, m \rangle = 0$$

namely if $m' - m - g \neq 0$

$$\langle \alpha'; j', m' | T_k^g | \alpha; j, m \rangle = 0$$

* From $[J_-, T_k^g] = \sqrt{k(k+1) - g(g-1)} T_k^{g-1}$

$$\langle \alpha'; j', m' | J_- T_k^g | \alpha; j, m \rangle = \langle \alpha'; j', m' | T_k^g J_- | \alpha; j, m \rangle = \sqrt{k(k+1) - g(g-1)} \langle \alpha'; j', m' | T_k^{g-1} | \alpha; j, m \rangle$$

$$\Rightarrow \sqrt{k(k+1) - \xi(\xi-1)} \langle \alpha' j' m' | T_k^{\xi-1} | \alpha j m \rangle$$

$$+ \sqrt{j(j+1) - m(m-1)} \langle \alpha' j' m' | T_k^{\xi} | \alpha j m-1 \rangle$$

$$= \sqrt{j'(j'+1) - m'(m'+1)} \langle \alpha' j' m'+1 | T_k^{\xi} | \alpha j m \rangle$$

(The recursion relation is same as the CG from J_+)

★ From $[J_+, T_k^{\xi}] = \sqrt{k(k+1) - \xi(\xi+1)} T_k^{\xi+1}$

$$\Rightarrow \langle \alpha' j' m' | J_+ T_k^{\xi} | \alpha j m \rangle - \langle \alpha' j' m' | T_k^{\xi} J_+ | \alpha j m \rangle$$

$$= \sqrt{k(k+1) - \xi(\xi+1)} \langle \alpha' j' m' | T_k^{\xi+1} | \alpha j m \rangle$$

$$\Rightarrow \sqrt{j'(j'+1) - m'(m'-1)} \langle \alpha' j' m'-1 | T_k^{\xi} | \alpha j m \rangle$$

$$= \sqrt{j(j+1) - m(m+1)} \langle \alpha' j' m' | T_k^{\xi} | \alpha j m+1 \rangle$$

$$+ \sqrt{k(k+1) - \xi(\xi+1)} \langle \alpha' j' m' | T_k^{\xi+1} | \alpha j m \rangle$$

(Same as the recursion relations for CG from J_-)

This is the same recursion relations for the CG coefficients with $(j, j', m, m') \Rightarrow (j, k, m, \xi)$

\Rightarrow Therefore the m behavior of $\langle \alpha' j' m' | T_k^{\xi} | \alpha j m \rangle$ must be proportional to the CG coefficient, We write

$$\langle \alpha' j' m' | T_k^{\xi} | \alpha j m \rangle = \underbrace{\langle j k ; m \xi | j k j' m' \rangle}_{\text{(terms independent of } m', m \text{ and } \xi)}}$$

This part is called $\frac{1}{\sqrt{2j+1}}^k$ (reduced matrix)
 convention

Recap: Wigner-Eckart Theorem

$$\langle j' m' | T_L^M | j m \rangle = \langle j L ; m M | j' m' \rangle \frac{\langle j' || T_L || j \rangle}{\sqrt{2j'+1}}$$

$\langle || \cdot || \rangle$ is called the "reduced matrix element"

It depends on j, j', L . But it is independent of the quantum numbers m, m' and M .

* The WE Theorem: The matrix element can be separated into a pure geometrical factor (the CG coeff) and a factor containing the dynamics (the reduced matrix element)

⇒ Selection rules:

$$\langle j' m' | T_L^M | j m \rangle \neq 0 \Rightarrow \boxed{m' = M + m}$$

also the "triangle rule",

$$j' = j + L, j + L - 1, \dots, |j - L|$$

* A tensor of zero rank can only connect states of the same total j .

* A tensor of first rank can only connect states

$$j' = j \text{ or } j' = j \pm 1 \text{ (However, it can't connect } j=0 \text{ to } j'=0 \text{)}$$

* A tensor of the second rank can only connect states

$$j' = j, j \pm 1, j \pm 2$$

however, it can't connect states with $j' = j = 0$

$$j' = j = \frac{1}{2}, \text{ or } j=1, j'=0 \text{ (or vice versa)}$$

An example of WE theorem application.

$$\begin{aligned}
& \langle 2, 0 | Y_{10} | 1, 0 \rangle \\
&= \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi Y_{20}^*(\theta, \phi) Y_{10}(\theta, \phi) Y_{10}(\theta, \phi) \\
&= \int_{-1}^1 d(\cos\theta) \int_0^{2\pi} d\phi \sqrt{\frac{5}{16\pi}} (3\cos^2\theta - 1) \left(\sqrt{\frac{3}{4\pi}} \cos\theta\right)^2 \\
&= \frac{2\pi}{\pi} \frac{\sqrt{15}}{8} \sqrt{\frac{3}{4\pi}} \int_{-1}^1 dx (3x^2 - 1) x^2 = \frac{1}{\sqrt{15\pi}}
\end{aligned}$$

By WE theorem

$$\frac{1}{\sqrt{15\pi}} = \underbrace{\langle 11; 00 | 2, 0 \rangle}_{\frac{1}{\sqrt{3}}} \frac{\langle 2 || Y_1 || 1 \rangle}{\sqrt{2 \cdot 1 + 1}}$$

from the CG coeff table

$$\Rightarrow \langle 2 || Y_1 || 1 \rangle = \frac{1}{\sqrt{15\pi}} \sqrt{3} \cdot \sqrt{\frac{8}{2}} = \frac{3}{\sqrt{10\pi}}$$

then by WE theorem and the CG table, we can quickly obtain the following ~~result~~ matrix elements:

$$\langle 2, 1 | Y_{10} | 1, 1 \rangle = \underbrace{\langle 11; 01 | 2, 1 \rangle}_{\frac{1}{\sqrt{2}}} \frac{\langle 1 || Y_1 || 1 \rangle}{\sqrt{3}} = \frac{1}{\sqrt{15\pi}} \frac{\sqrt{\frac{1}{2}}}{\sqrt{\frac{2}{3}}} = \frac{1}{2} \sqrt{\frac{3}{15\pi}}$$

$$\langle 2, 2 | Y_{11} | 1, 1 \rangle = \underbrace{\langle 11; 11 | 2, 2 \rangle}_1 \frac{\langle 1 || Y_1 || 1 \rangle}{\sqrt{3}} = \sqrt{\frac{3}{10\pi}}$$

$$\langle 2, -1 | Y_{1-1} | 1, 0 \rangle = \underbrace{\langle 11; -10 | 2, -1 \rangle}_{\frac{1}{\sqrt{2}}} \frac{\langle 1 || Y_1 || 1 \rangle}{\sqrt{3}} = \frac{1}{2} \sqrt{\frac{3}{15\pi}}$$

$$\langle 2, 0 | Y_{1-1} | 1, 1 \rangle = \underbrace{\langle 11; -1, 1 | 2, 0 \rangle}_{\frac{1}{\sqrt{6}}} \frac{\langle 1 || Y_1 || 1 \rangle}{\sqrt{3}} = \frac{1}{2} \frac{1}{\sqrt{15\pi}}$$