

① quick review of the previous lecture.

②. $|1\rangle, |2\rangle, \dots$ $\langle s'|s\rangle = \delta_{ss'}$

are orthonormal basis

② $|s\rangle$ is just a naming protocol, if like ^{Arabic number}.

they can be called $|p\rangle, |e\rangle, |m\rangle, \dots$ or $|\frac{1}{2}\rangle, |\frac{1}{3}\rangle, |\frac{1}{4}\rangle, \dots$

bottom line is the ordering $\rightarrow$ integer system $100, 101, 102, \dots$

③. Zuff additional to $|s\rangle$,

$|k_n\rangle$ is also a good basis for the N -dim Hilbert space.

$$|k_0\rangle = A(|1\rangle + |2\rangle + |3\rangle) \quad k_n = \frac{2\pi n}{3}$$

$$|k_1\rangle = A\left(e^{\frac{2\pi i}{3}}|1\rangle + e^{\frac{4\pi i}{3}}|2\rangle + e^{\frac{6\pi i}{3}}|3\rangle\right)$$

$$|k_2\rangle = A\left(e^{\frac{4\pi i}{3}}|1\rangle + e^{\frac{8\pi i}{3}}|2\rangle + e^{\frac{12\pi i}{3}}|3\rangle\right)$$

$$\langle k_n | k_{n'} \rangle = A^2 \delta_{n,n'}$$

④. you can obtain $|s\rangle$ from $|k_n\rangle$

$$\text{e.g. } |1\rangle = \frac{1}{3A} \left[|k_0\rangle + e^{-\frac{2\pi i}{3}} |k_1\rangle + e^{-\frac{4\pi i}{3}} |k_2\rangle \right]$$

$$|2\rangle = \frac{1}{3A} \left[|k_0\rangle + e^{-\frac{4\pi i}{3}} |k_1\rangle + e^{-\frac{8\pi i}{3}} |k_2\rangle \right]$$

$$|3\rangle = \frac{1}{3A} \left[|k_0\rangle + |k_1\rangle + |k_2\rangle \right]$$

In the $N \rightarrow \infty$ limit

$$|x\rangle \propto \int dk e^{-ikx} |k\rangle$$

$$\text{or } |x\rangle = \int dk |k\rangle \langle k | x \rangle = \frac{1}{\sqrt{2\pi}} \int dk e^{-ikx} |k\rangle$$

2. F & H 的自伴性.

3. uncertainty principle.

$$A_1 = A - \bar{A}, \quad B_1 = B - \bar{B}, \quad A^\dagger = A, \quad B^\dagger = B.$$

Therefore $A_1^\dagger = A_1, \quad B_1^\dagger = B_1$.

let $|\alpha\rangle = A_1|\psi\rangle, \quad |\beta\rangle = B_1|\psi\rangle$

by Schwarz inequality

$$\langle A_1^2 \rangle = \langle A^2 - 2A\bar{A} + \bar{A}^2 \rangle = \langle A^2 \rangle - \langle \bar{A}^2 \rangle = \sigma_A^2$$

$$\underbrace{\langle \alpha | \alpha \rangle}_{\sigma_A^2} \underbrace{\langle \beta | \beta \rangle}_{\sigma_B^2} \geq |\langle \alpha | \beta \rangle|^2$$

$$A_1 B_1 = \frac{A_1 B_1 + B_1 A_1}{2} + \frac{A_1 B_1 - B_1 A_1}{2} = \frac{1}{2} \{A_1, B_1\} + \frac{1}{2} [A_1, B_1]$$

↑
anti commutator

↑
commutator

↓
Hermitian

↓
anti Hermitian

Real

↓
purely imaginary

$$\overline{\langle \psi | O | \psi \rangle} = \langle \psi | O^\dagger | \psi \rangle = \langle \psi | O | \psi \rangle$$

$$\overline{\langle \psi | A | \psi \rangle} = \langle \psi | A^\dagger | \psi \rangle = -\langle \psi | A | \psi \rangle$$

$$\Rightarrow |\langle \psi | A_1 B_1 | \psi \rangle|^2 = \left| \underbrace{\frac{1}{2} \langle \psi | \{A_1, B_1\} | \psi \rangle}_{\text{Real}} + \underbrace{\frac{1}{2} \langle \psi | [A_1, B_1] | \psi \rangle}_{\text{Imaginary}} \right|^2$$

$$= \frac{1}{4} |\langle \{ \} \rangle|^2 + \frac{1}{4} |\langle [\] \rangle|^2 \geq \frac{1}{4} |\langle [\] \rangle|^2$$

$$[A_1, B_1] = [A - \bar{A}, B - \bar{B}] = (AB - A\bar{B} - \bar{A}B + \bar{A}\bar{B}) - (BA - B\bar{A} - \bar{B}A + \bar{B}\bar{A}) = [A, B]$$

$$\Rightarrow \boxed{\sigma_A^2 \sigma_B^2 \geq \frac{1}{4} |\langle [A, B] \rangle|^2}$$

$$\sigma_x^2 \sigma_p^2 \geq \frac{1}{4} \left| \underbrace{\langle [X, P] \rangle}_{i\hbar} \right|^2 = \left(\frac{\hbar}{2} \right)^2$$

$$\Rightarrow \boxed{\sigma_x \sigma_p \geq \frac{\hbar}{2}}$$

It works for any pair of operators!

→ It's mathematical outcome, not put in by hand

if $[A, B] = 0$, A, B are called compatible or commuting observables.

★ An important property $\Rightarrow$ without going into the proof, see appendix

A, B admit complete sets of simultaneous eigenstate.

They can be simultaneously diagonalized;

an example $[L_z, L^2] = 0$,

any state can be labeled by (l, m)

★ comment $\Delta E \Delta t \geq \frac{\hbar}{2}$ is not fundamental, because t is NOT an operator in QM. The relation is derivative —

★ 考題有時間，3 1/2 分 Lecture 很多人犯錯的概念

① current vs density ρ, j

but not $m(X) = \langle p \rangle$

② time evolution vs. Ehrenfest relation, $\psi = \sum c_n e^{-\frac{iE_n t}{\hbar}} \phi_n$, $m \frac{d}{dt} \langle X \rangle = \langle p \rangle$

③ W.T. matching, ψ continuous, ψ' may/may not —

④ scattering state 不能歸一 vs. bound state