



Biophysics

Lecture at National Center for Theoretical Sciences
2007/4/27

What do people study in this discipline?

Virtual Journal of Biological Physics Research

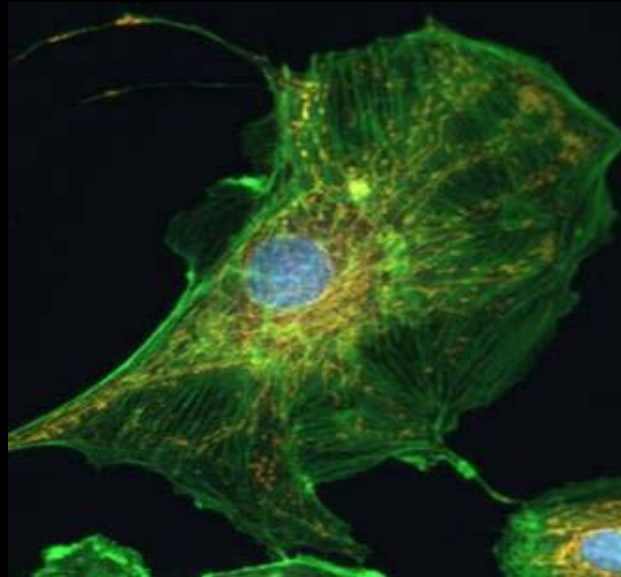
<http://www.vjbio.org/bio/>

- quantum mechanical dynamics
- physics of water and hydrogen-bonded solvents
- membrane biophysics
- fundamental polymer statics/dynamics
- protein conformational dynamics/folding
- DNA conformational dynamics
- single molecule dynamics
- intermolecular interactions
- physical studies of cell mechanics
- information transfer in biological systems
- multicellular phenomena
- biological networks
- quantitative genomics
- statistical and nonlinear physics
- instrumentation development
- miscellaneous

Outline

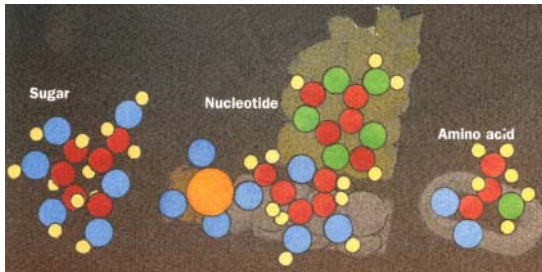
1. Physics in a cell
2. Self-assembling (protein folding)
3. Mechanics (membrane & filaments)
4. Energy transduction (biological motors)
5. Electrostatics (neuron cells & networks)
6. Thermodynamics (non-equilibrium biosystems)

1. Physics in a Cell



Size of different ingredients in a cell

Basic chemical elements in life: H, C, N, O, P, S



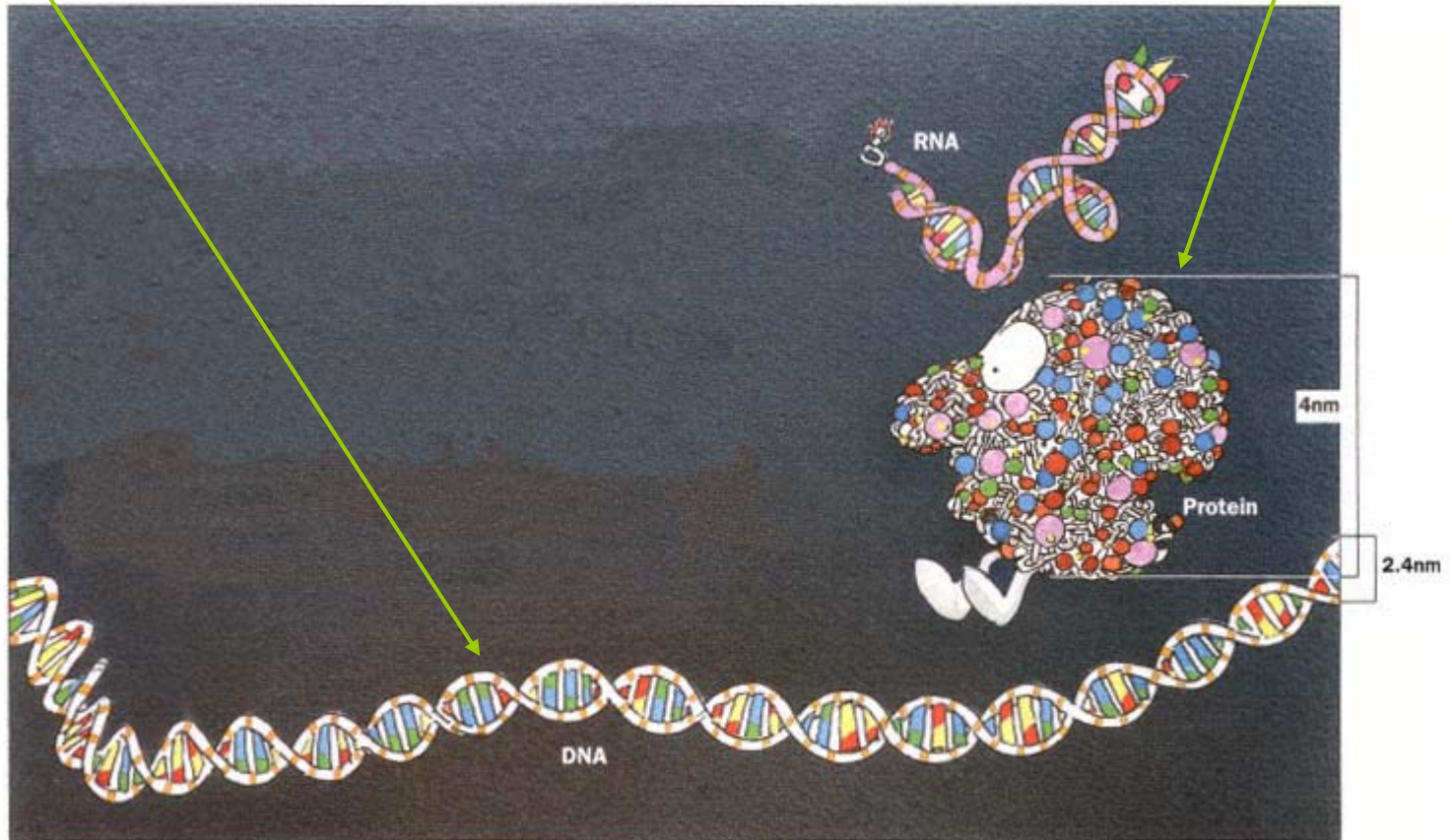
Small molecules
(2-50 atoms)

Big molecules
(300-400 atoms?)



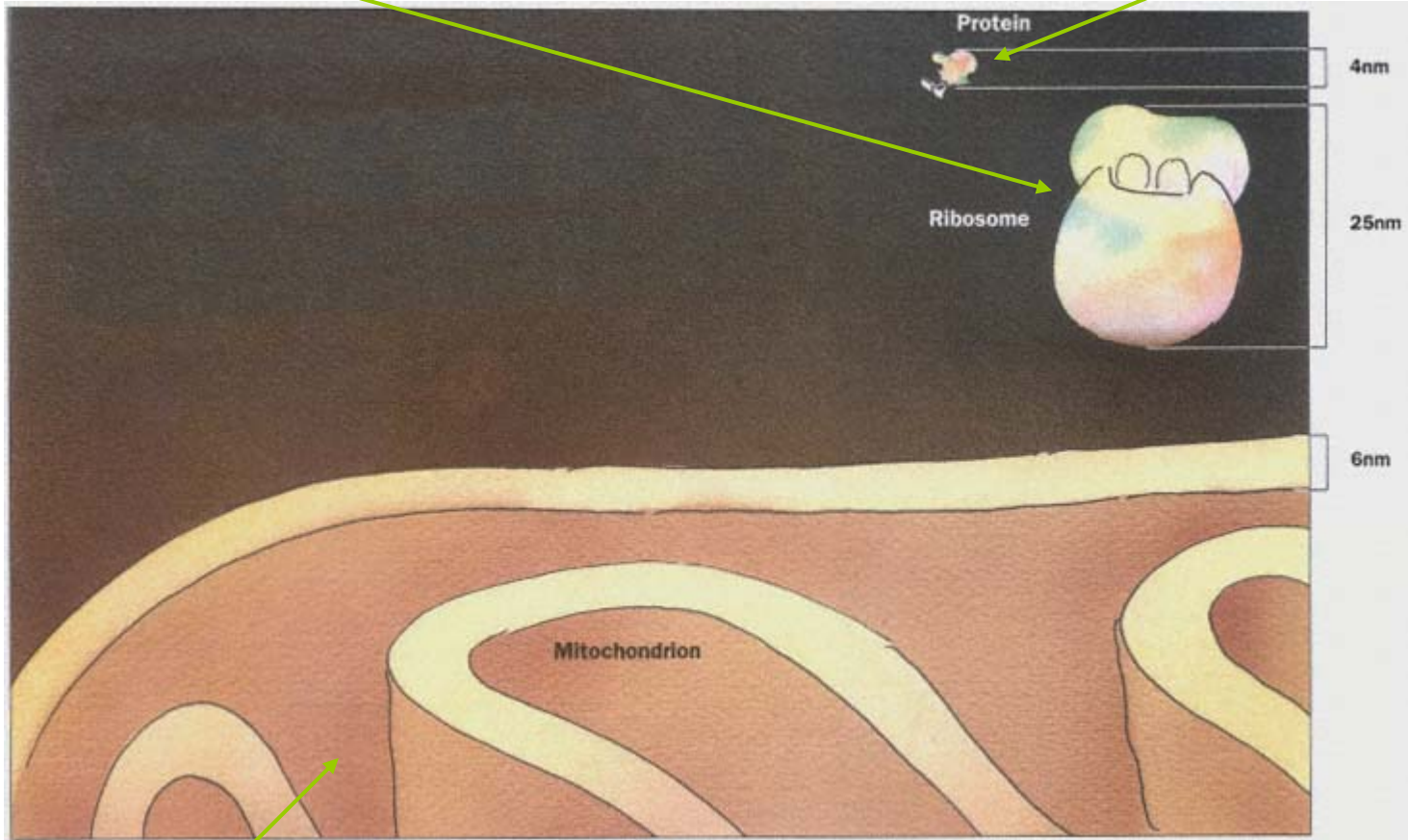
Size of different ingredients in a cell

DNA is a book describing how to produce proteins.



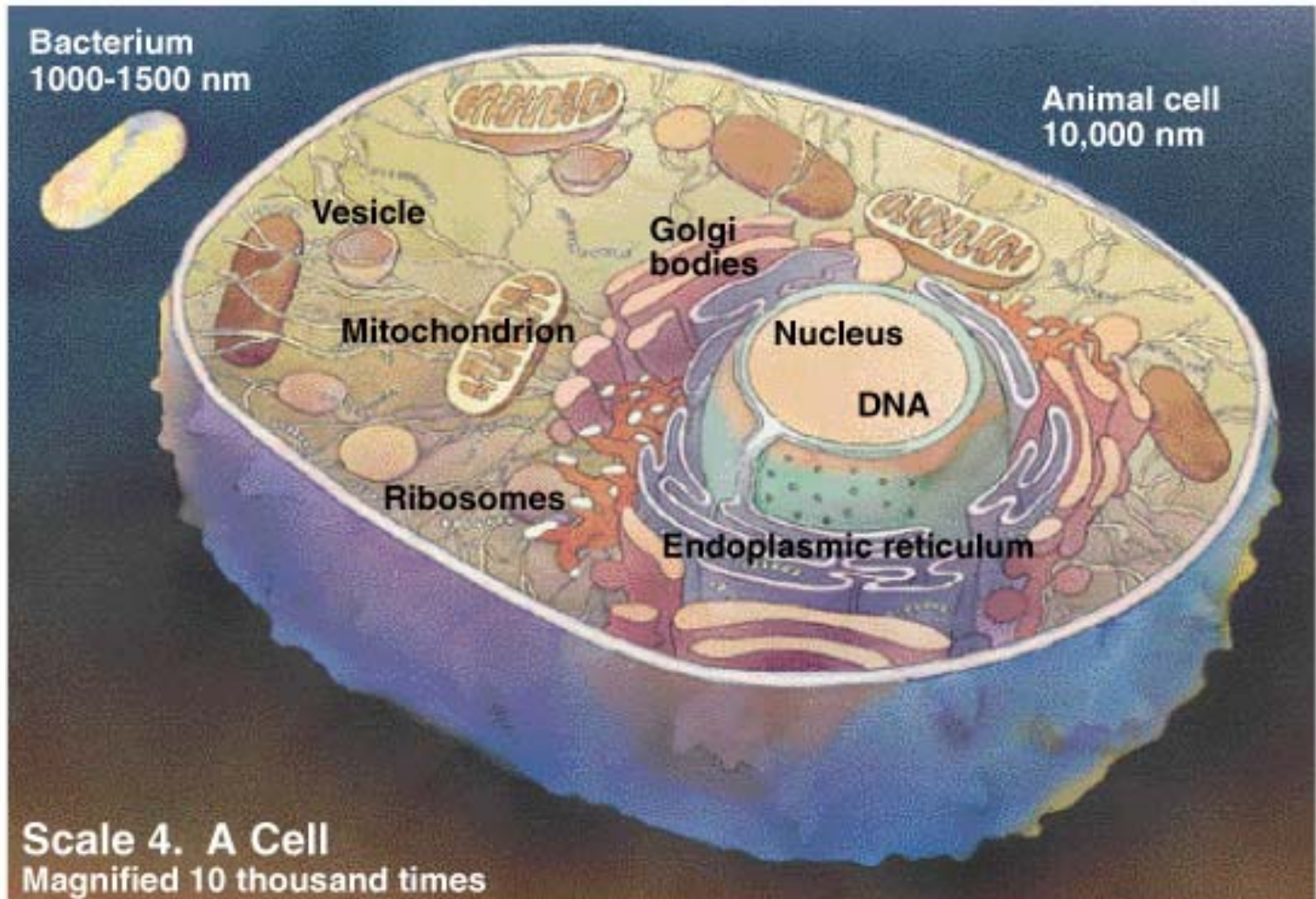
Size of different ingredients in a cell

Ribosome: an machine to produce proteins.

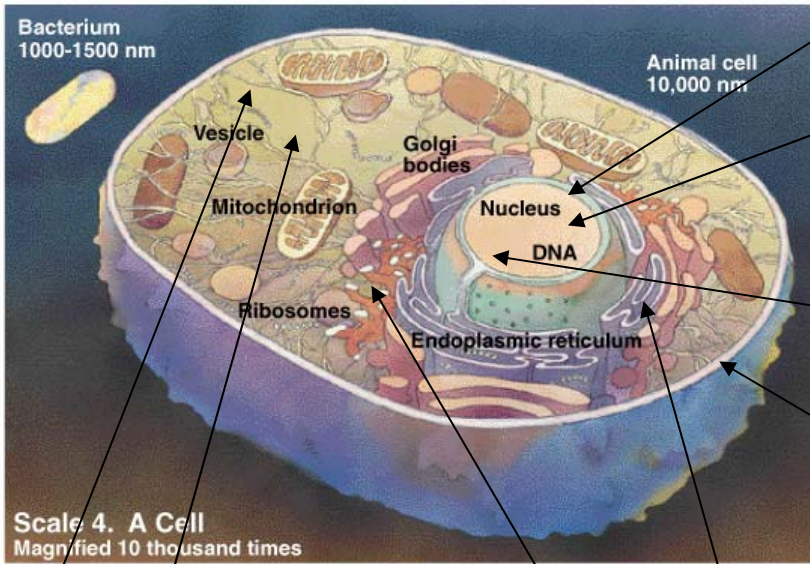


Mitochondrion: an factory to produce energy.

Size of different ingredients in a cell



Components in an animal cell



Nucleus(細胞核): the repository of a cell's genetic information.

Nucleolus(核仁): large, spherical structure in the nucleus; the site of ribosomal assembly.

Chromosome(染色體): the protein-complexed DNA molecules that become observable during cell divisions.

Membrane(細胞膜)

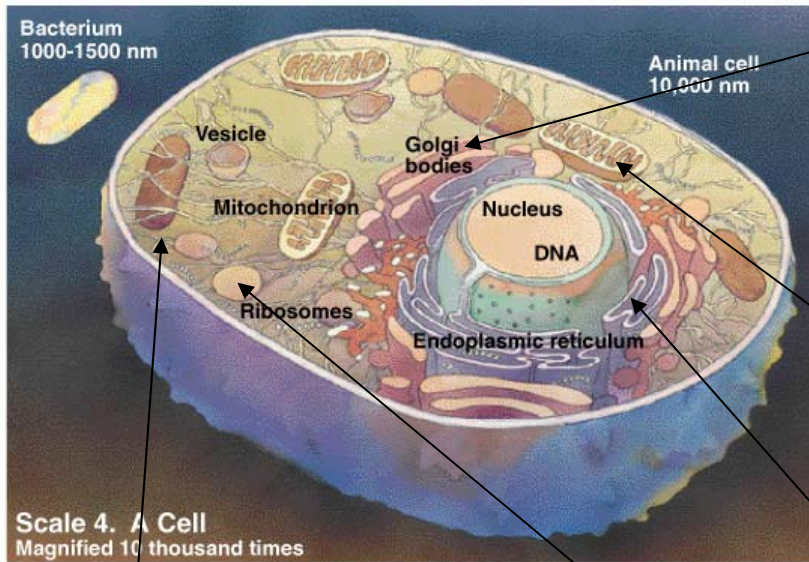
Cytosol(細胞液): the semifluid substance in which the organelles of the cytoplasm are suspended.

Cytoplasm(細胞質): the portion of interior of a cell that is not occupied by the nucleus.

Endoplasmic reticulum (ER) (內質網): network of interconnected membranes distributed throughout the cytoplasm.

Rough ER: ER studded with ribosomes on its side; involved in protein synthesis.

Smooth ER: ER that has no attached ribosomes; involved in packaging of secretory proteins and synthesis of lipids.



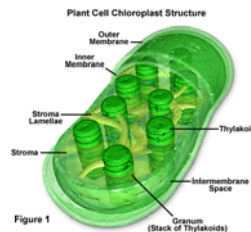
Golgi apparatus(高基氏體): stacks of flattened, disk-shaped membrane cisternae; involved in the processing & packaging of secretory proteins and in the synthesis of complex polysaccharides.

Mitochondrion(粒線體): double membrane-enclosed cytoplasmic organelle; the site of aerobic respiration.

Ribosome(核糖體): molecular machinery for protein synthesis.

Cytoskeleton(骨架): 3-dim network of filaments (microtubules, microfilaments, and intermediate filaments) that provides structure to the cytoplasm.

Chloroplast(葉綠體): double membrane-enclosed organelle of plants & algae that contains chlorophyll(葉綠素) and the enzymes necessary to carry out photosynthesis.

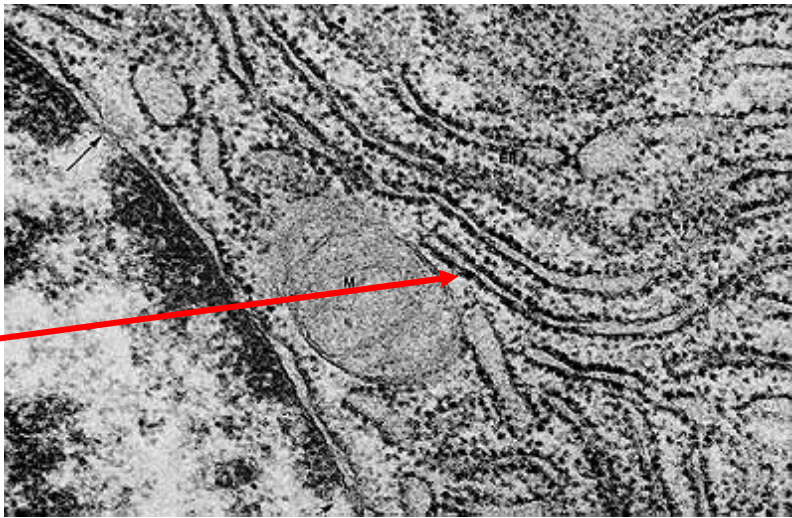
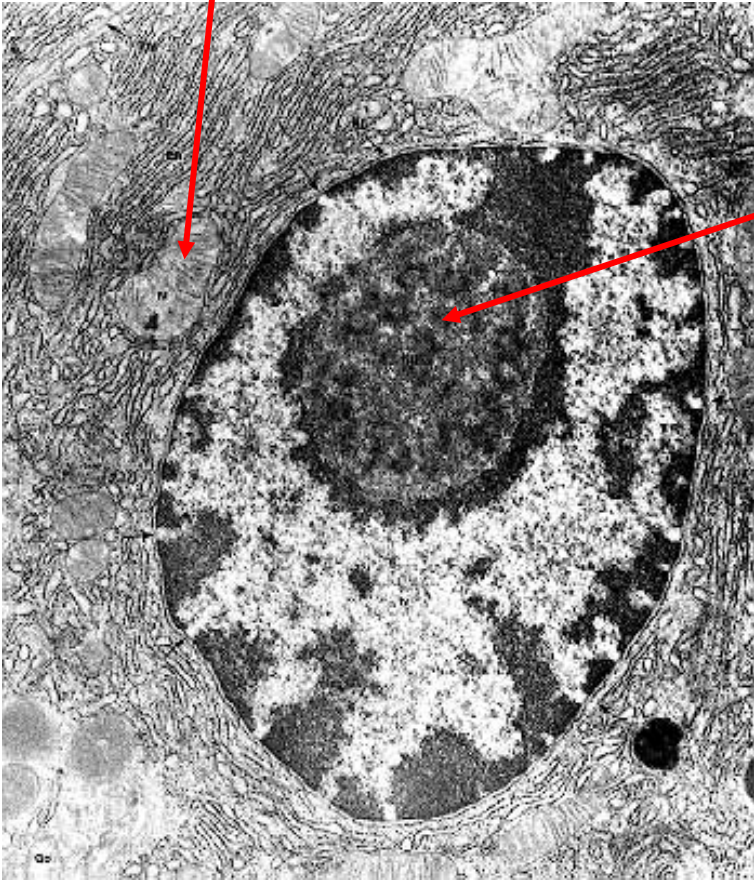


Lysosome(溶酶體): membrane-bounded organelle containing digestive enzymes capable of degrading all the major classes of biological macromolecules; formed by budding from the Golgi apparatus.

- (a) Mitochondrion
- (b) Rough ER
- (c) Nucleolus
- (d) Cell membrane
- (e) Nucleus

(a)

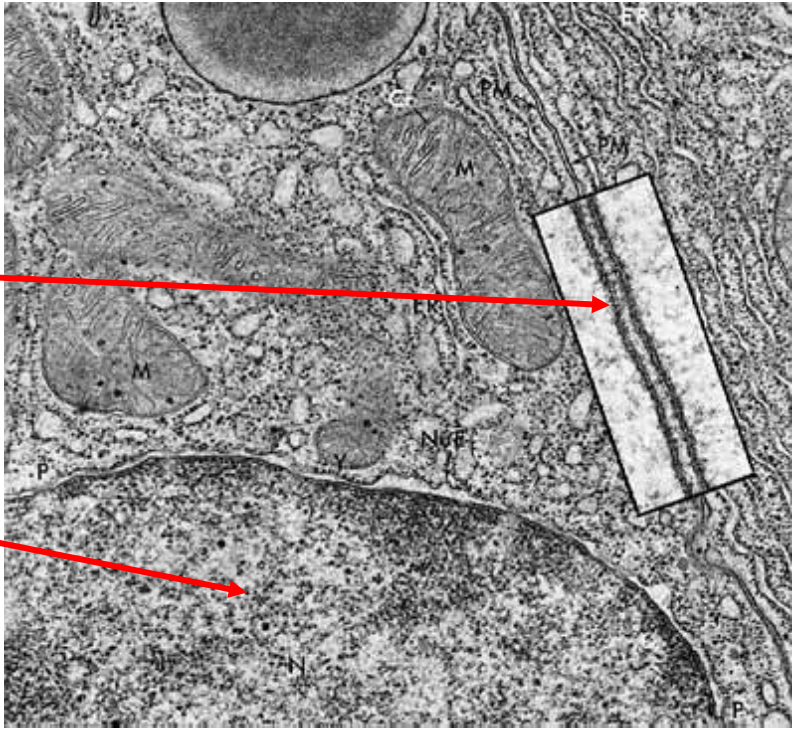
(b) black points



(c)

(d)

(e)

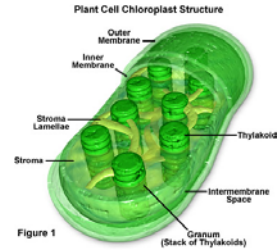
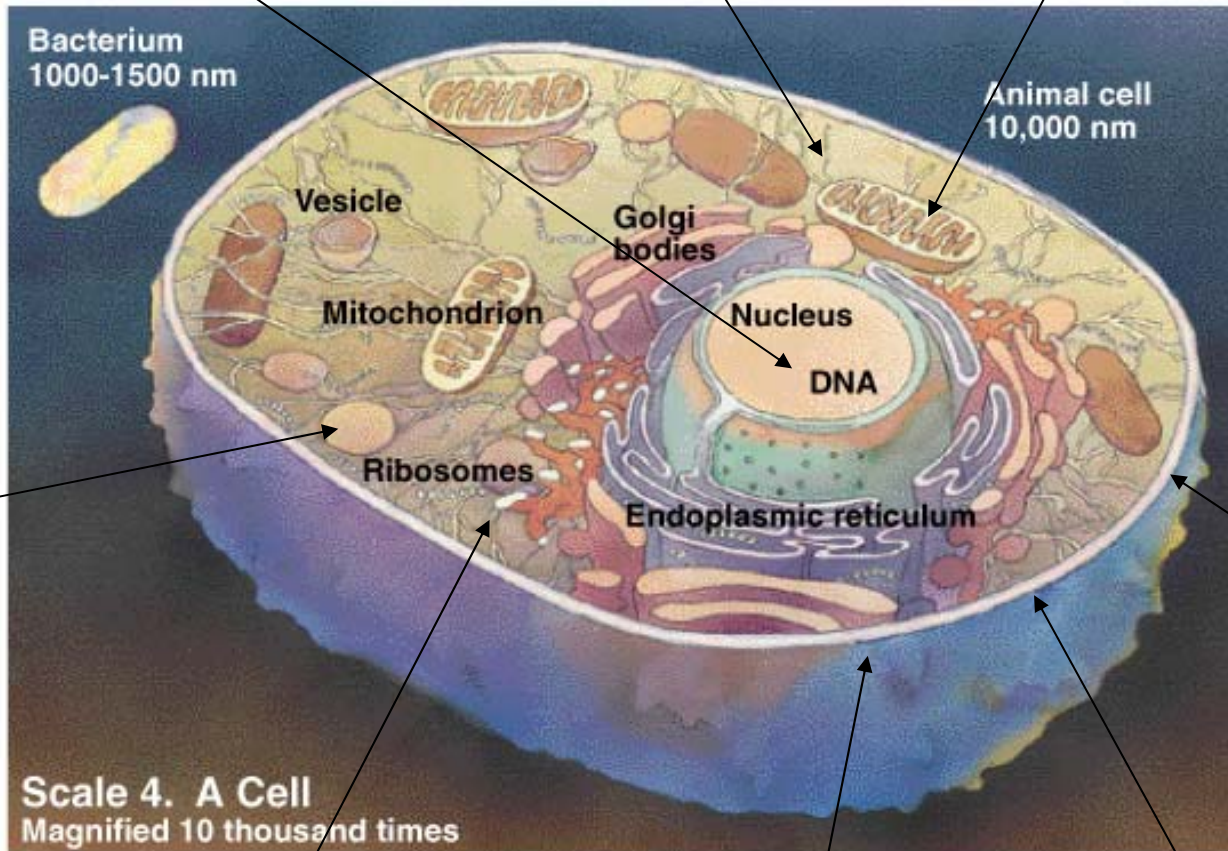


Physical problems inside a cell

Polymer: DNA

Transport: filaments

Energy: ATP, Chloroplast



Structure: protein

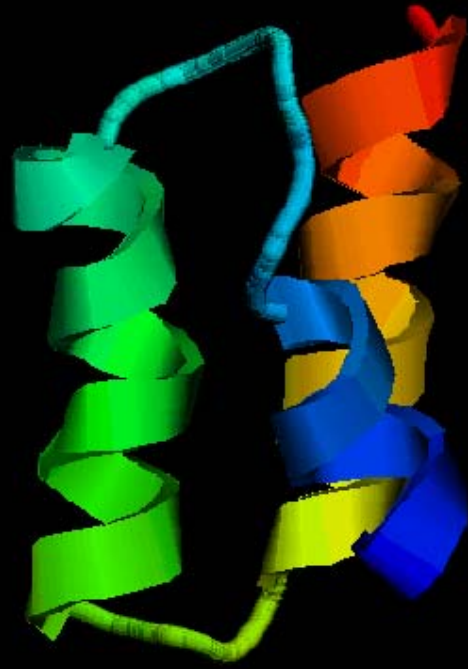
Ion channel

Colloid: cytoplasm

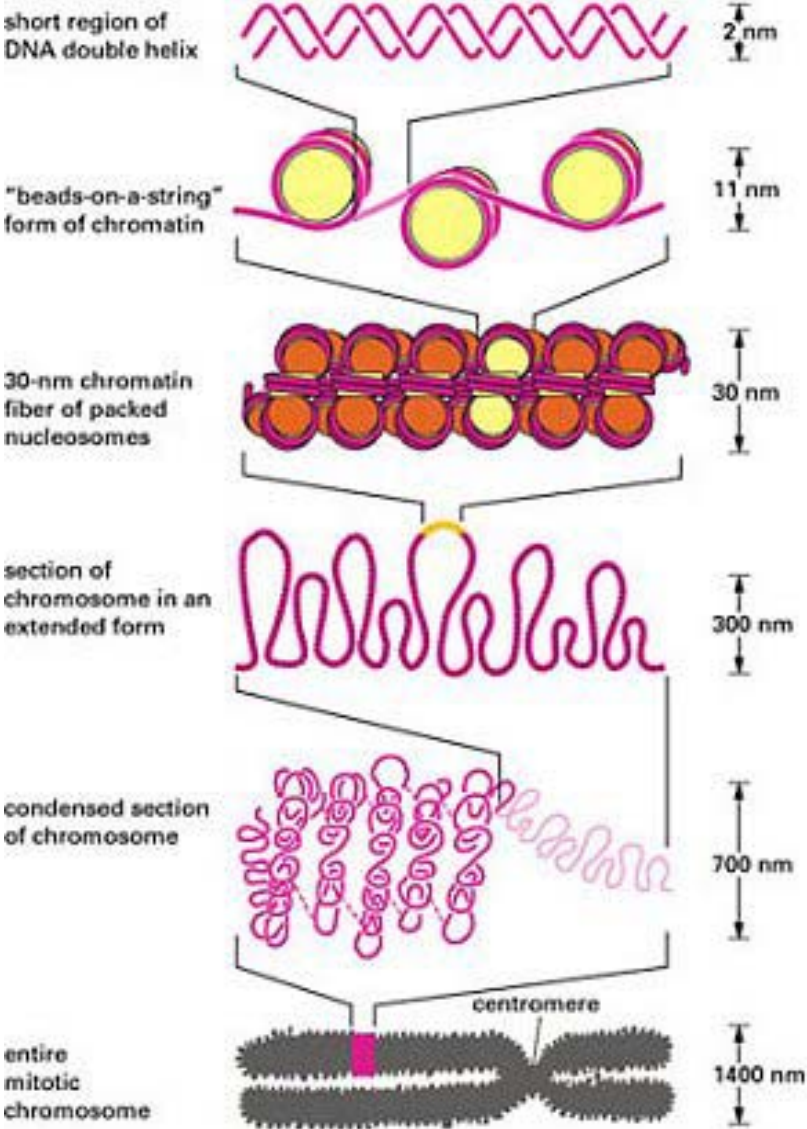
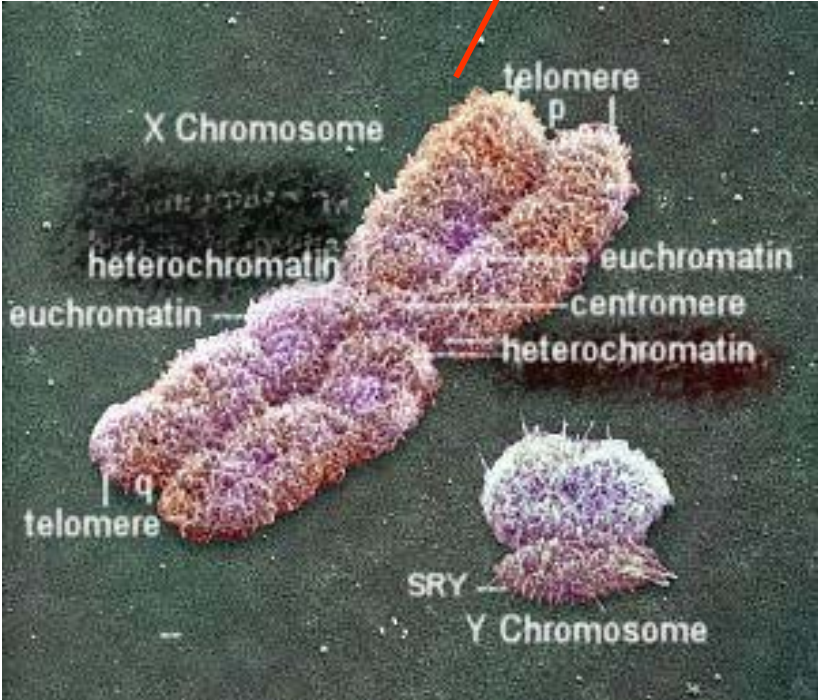
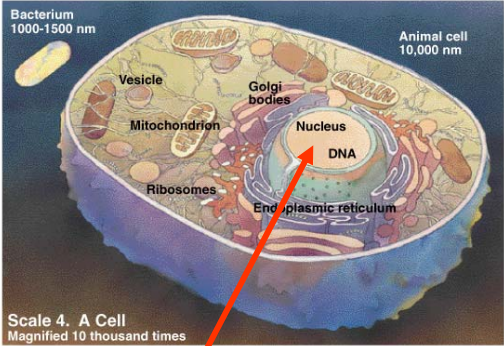
Membrane: bilayer

Motors: ion pump

2. Self-assembly (protein folding)

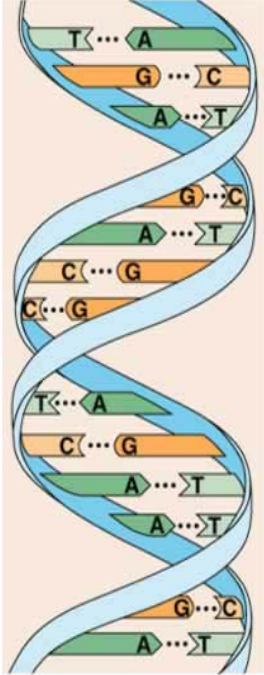
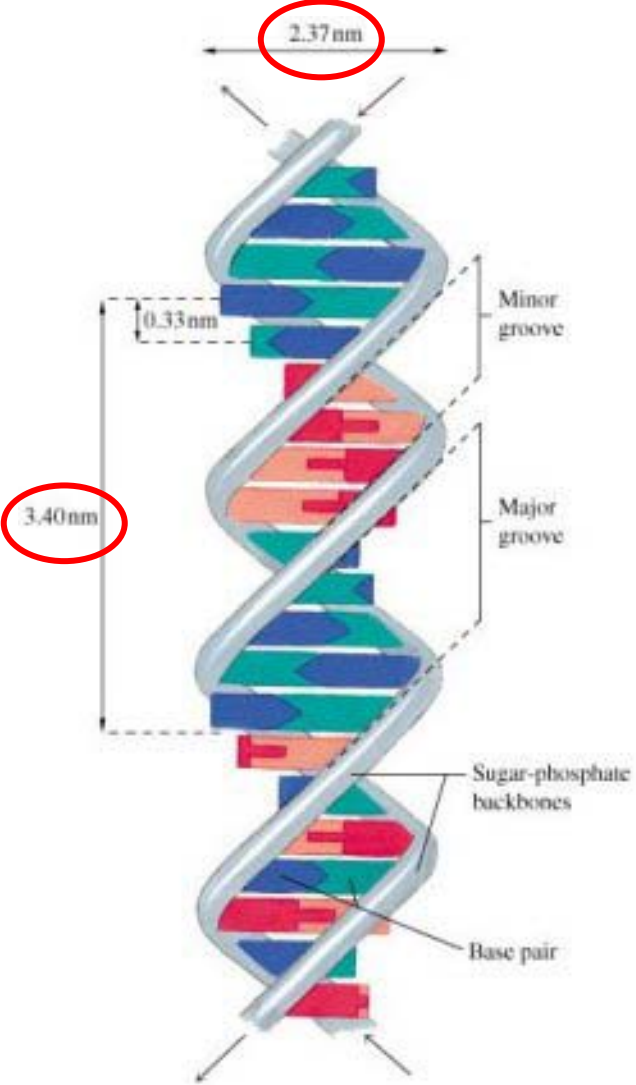


Chromosome and DNA

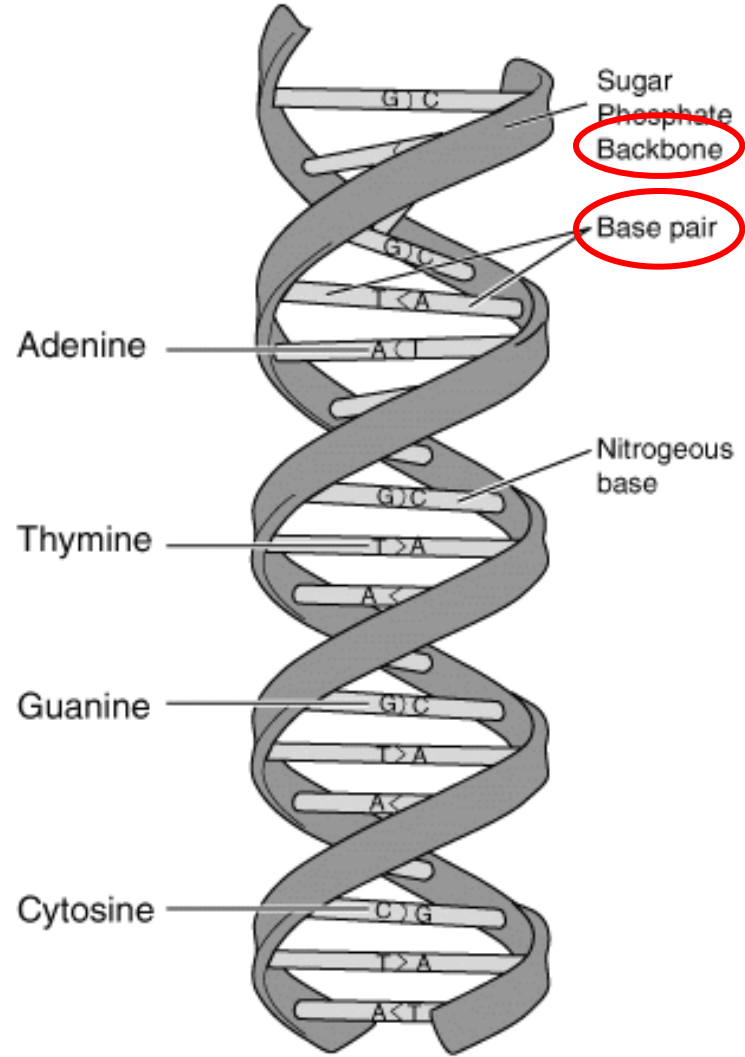
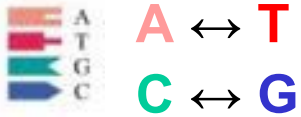


NET RESULT: EACH DNA MOLECULE HAS BEEN PACKAGED INTO A MITOTIC CHROMOSOME THAT IS 50,000x SHORTER THAN ITS EXTENDED LENGTH

DNA double helix



©Addison Wesley Longman, Inc.



4 nucleotides (核甘酸)

Gene

DNA =

(i) genes & (ii) junk DNA

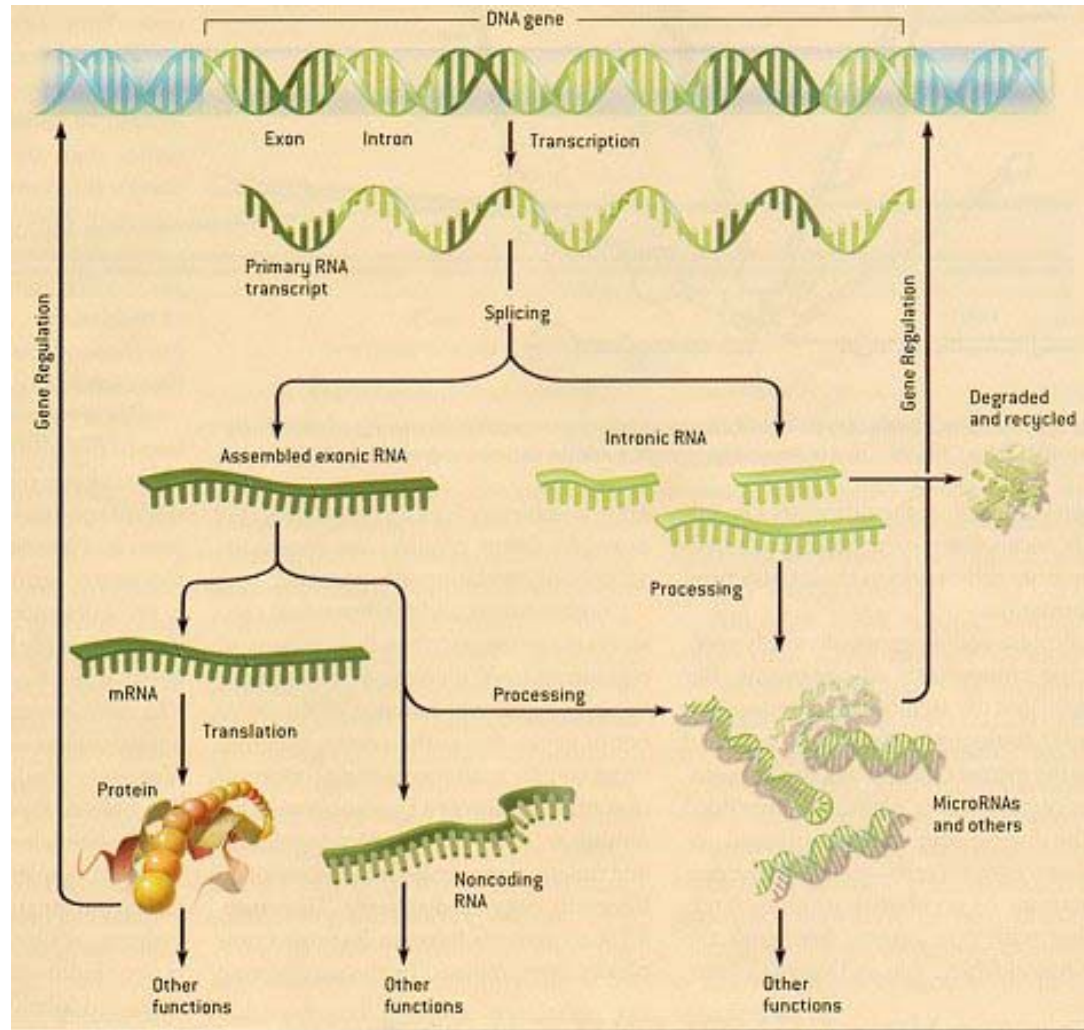
- (i) contains protein information
- (ii) non-coding DNA with repeated sequences inserted within a region of coding gene.

The purpose of the noncoding DNA, if any, is not understood. As much as 97% of human DNA is noncoding.

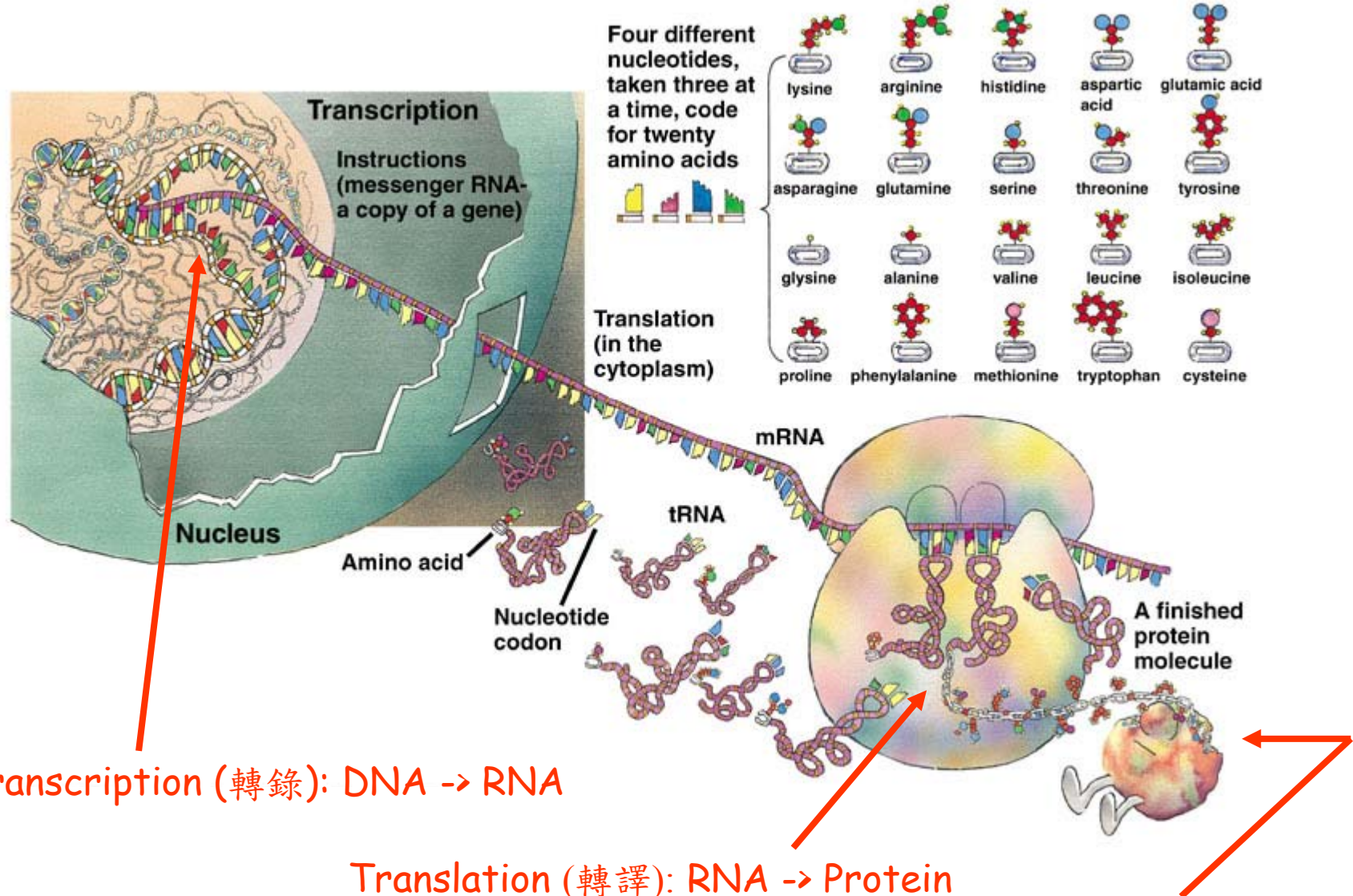
Some researches show that they might be used as testing site for genetic mutation.

Genome

the sum of all genes that make up the genetic code of an individual.



Gene expression: DNA → RNA → Protein



Transcription (轉錄): DNA → RNA

Translation (轉譯): RNA → Protein

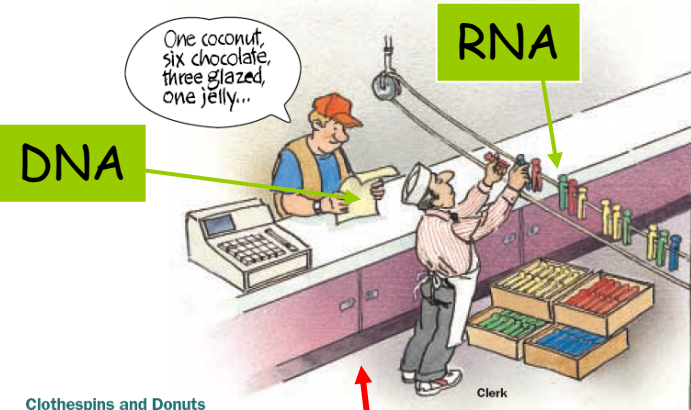
Protein folding

Transcription (轉錄): DNA → RNA

Nucleotides

Amino acids

5.4 How Orders Translate into Assembled Boxes of Donuts

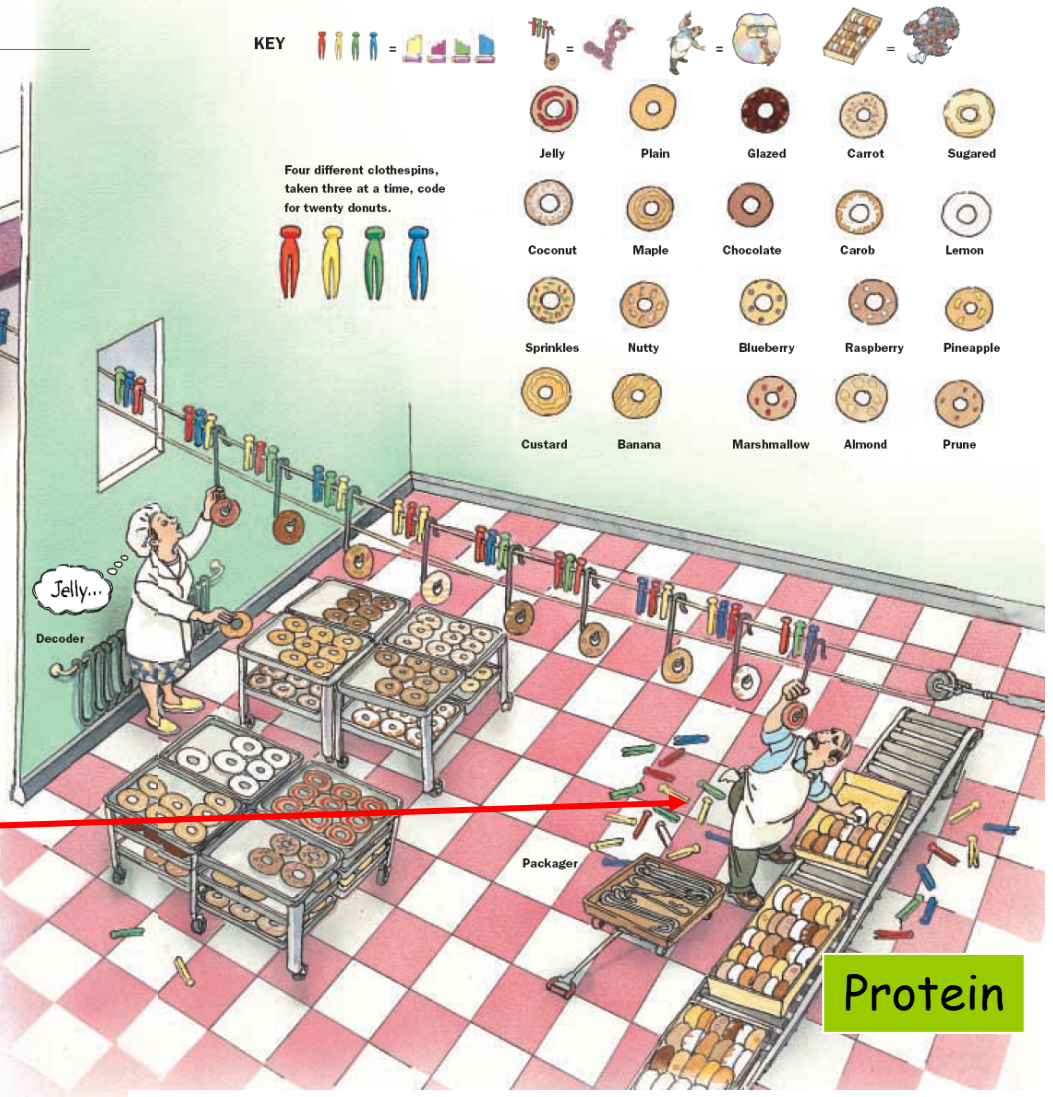
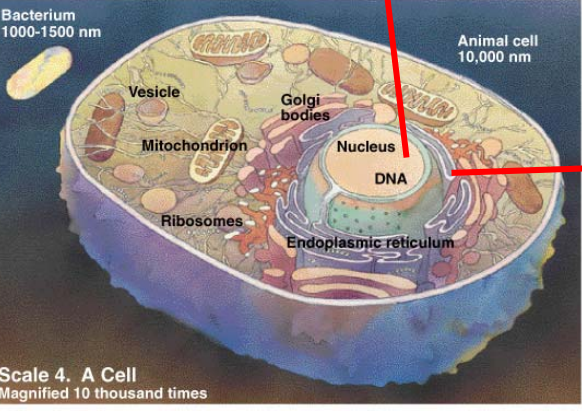


Clothespins and Donuts

KEY: = =

Four different clothespins, taken three at a time, code for twenty donuts.

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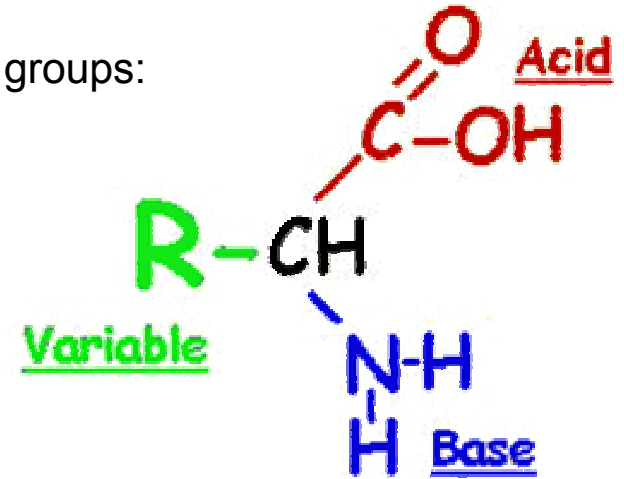
Translation (轉譯): RNA → Protein

Amino acid (氨基酸)

Proteins are necklaces of **amino acids**

A **amino acid** consists of a central carbon C_{α} linked by 4 groups:

- **carboxyl group** **COOH** (acidic). (羧基) \rightarrow COO^{-}
- **amino group** **NH₂** (basic). (氨基) \rightarrow NH_3^{+}
- hydrogen **H**.
- **residue R**.



A theoretical amino acid

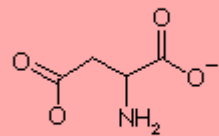
All 20 different amino acids have this same structure.

But their **side chain R** may vary in size, shape, charge, hydrophobicity, and reactivity.

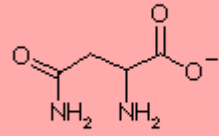
Different combinations of **R** give different proteins.

The side chains of 20 amino acids:

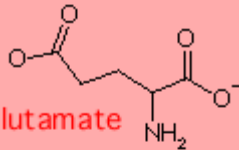
Acidic and amide side chains



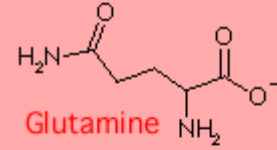
Aspartate



Asparagine

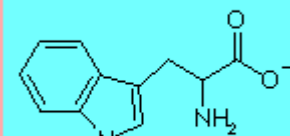


Glutamate

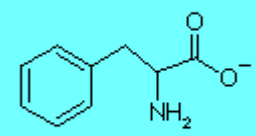


Glutamine

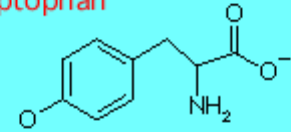
Aromatic side chains



Tryptophan

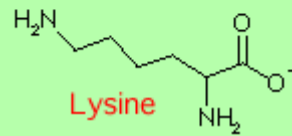


Phenylalanine

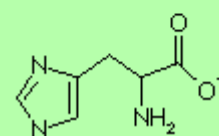


Tyrosine

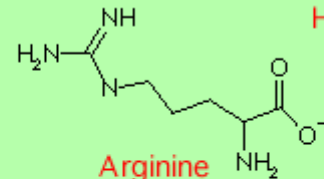
Basic side chains



Lysine

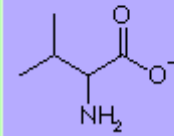


Histidine

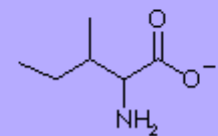


Arginine

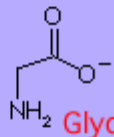
Aliphatic side chains



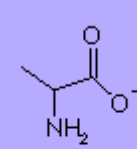
Valine



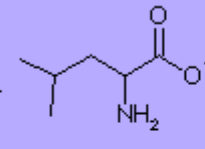
Isoleucine



Glycine

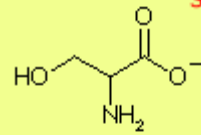


Alanine

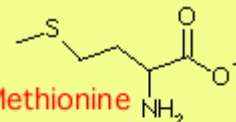


Leucine

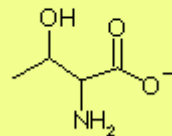
Hydroxyl or sulfur-containing side chains



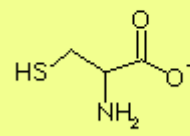
Serine



Methionine

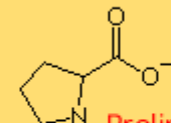


Threonine



Cysteine

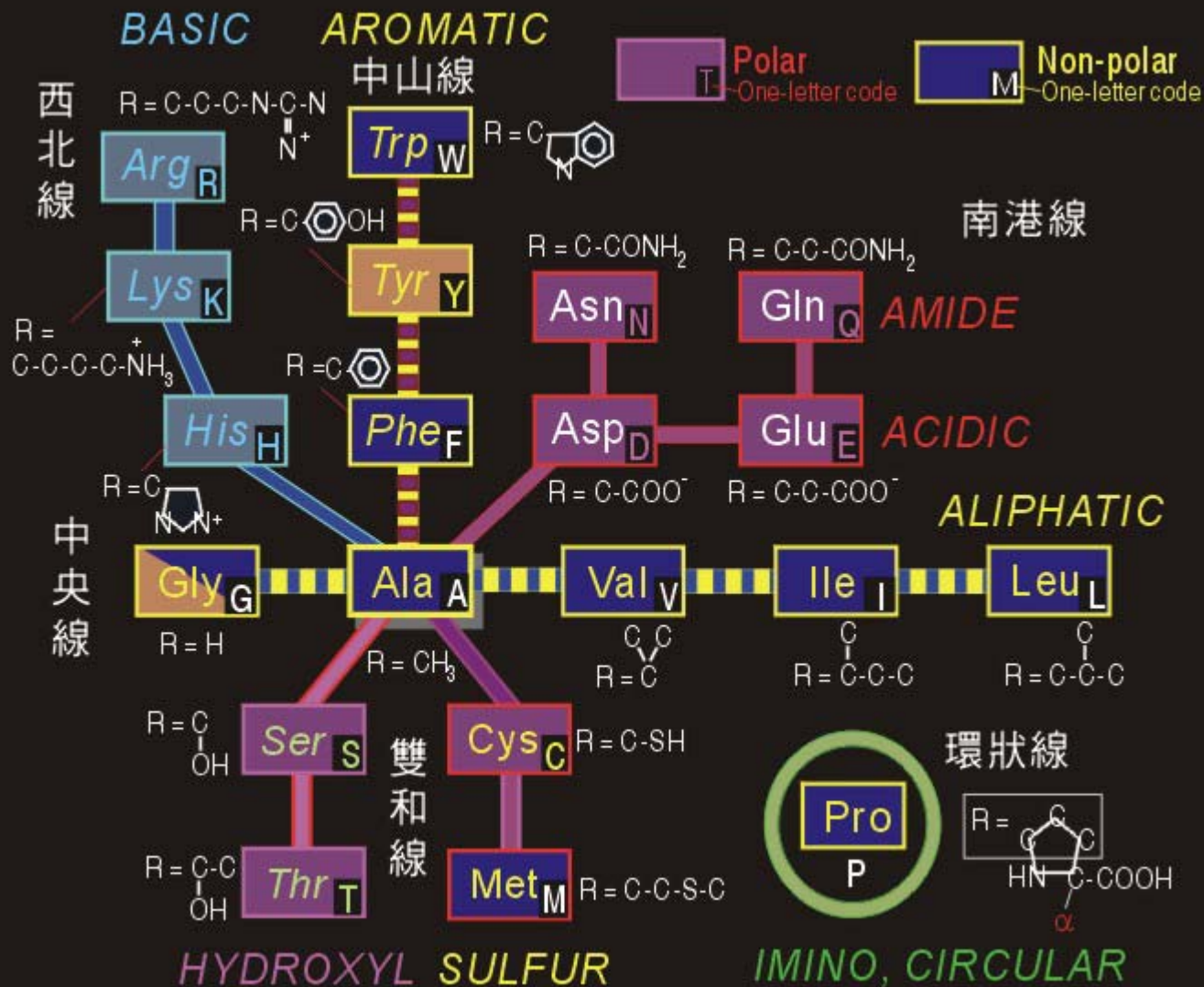
Cyclic side chain



Proline

http://www.bact.wisc.edu/Microtextbook/index.php?module=Book&func=displayarticle&art_id=40&theme=Printer

胺基酸地下鐵道圖

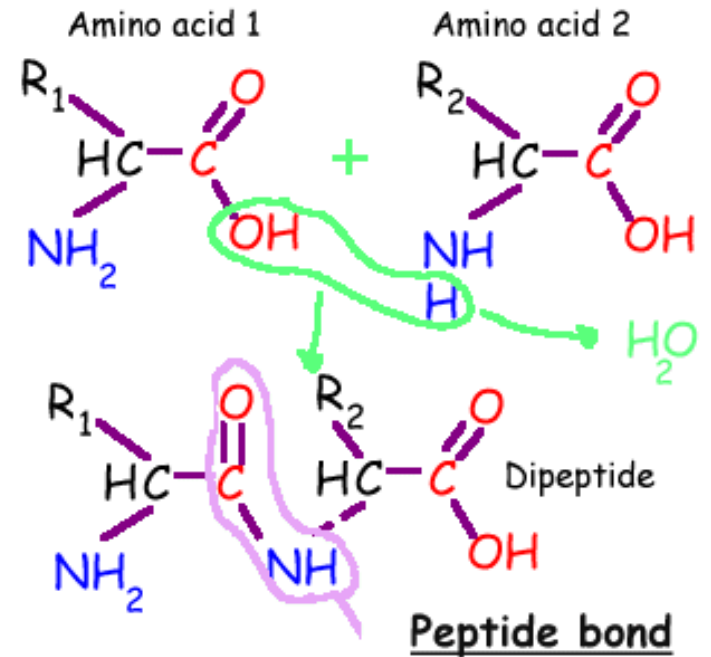


Peptide (胜肽, 缩氨酸)

Two amino acids can react together and create a dipeptide (basic reaction involved in the synthesis of proteins).

A molecule of water is released in the process.

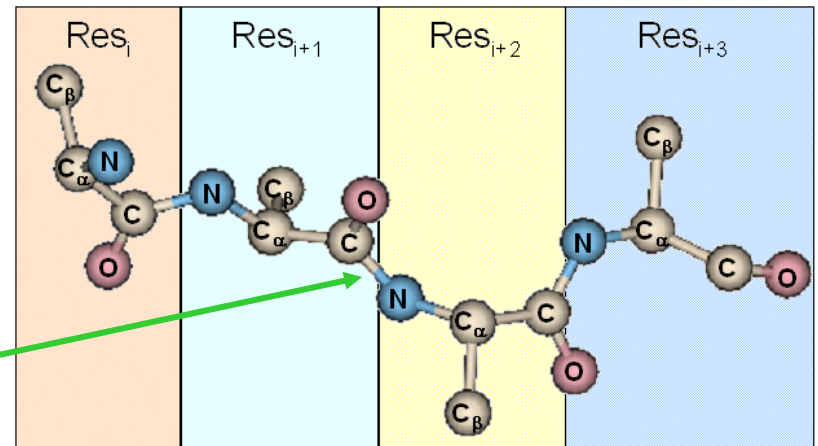
Peptides = {dipeptide, tripeptide, ...}.



Polypeptide

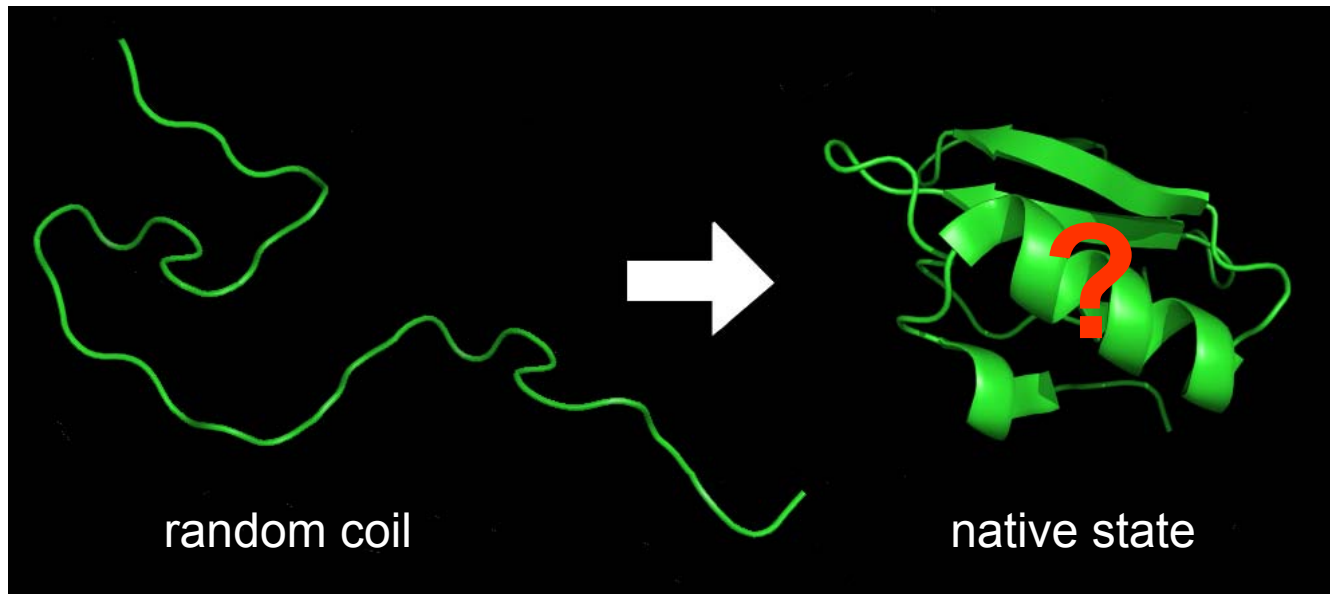
A series of applications of the peptide bond give rise to the **polypeptide** chain.

backbond



Protein folding:

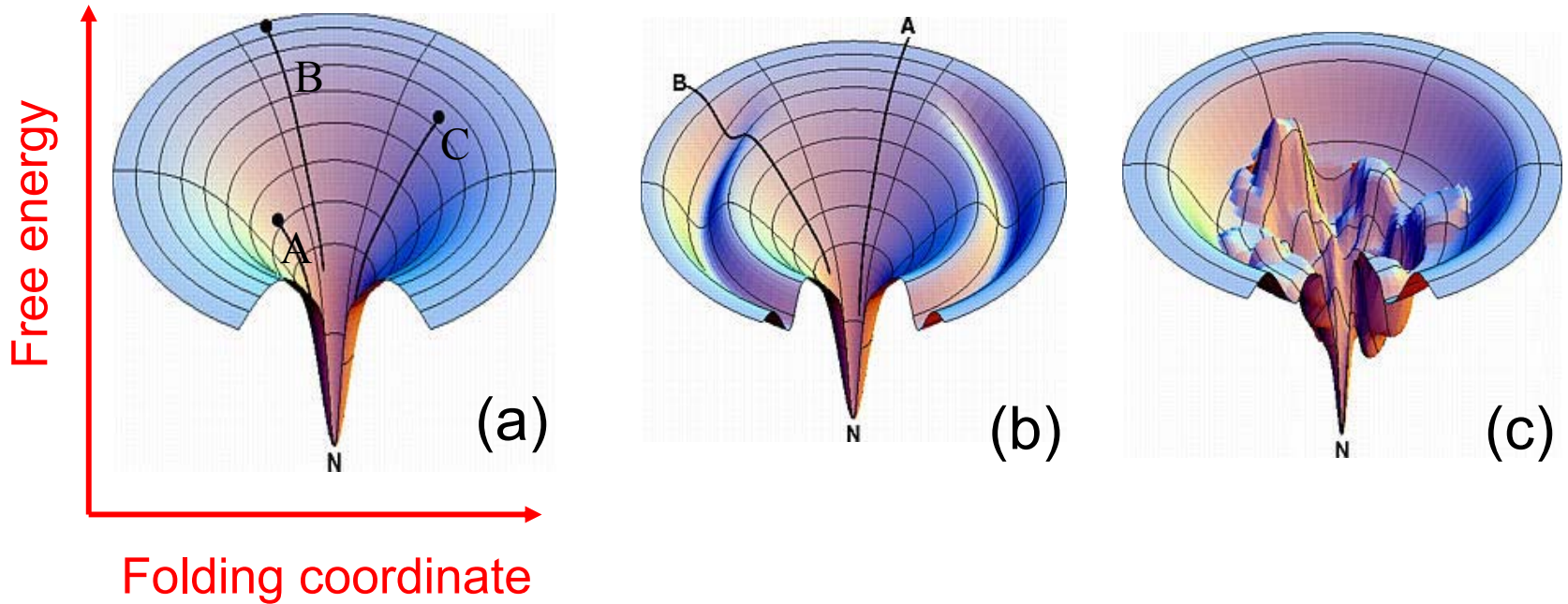
Put a polypeptide chain into water. Which structure would it form?



Anfinsen's dogma:

The native structure corresponds to the state with the lowest free energy of the protein-solvent system.

Funnel theory



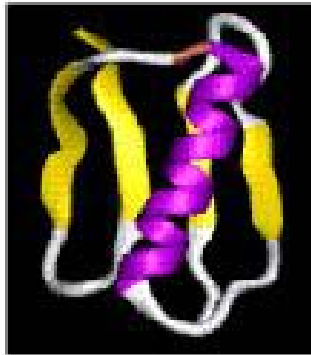
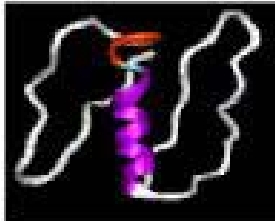
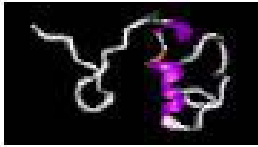
For a multidimensional smooth landscape like (a), to find its **minimum** is simple.

A real landscape is much more complex, with multiple local minima (folding traps) like (c).

Funnel theory

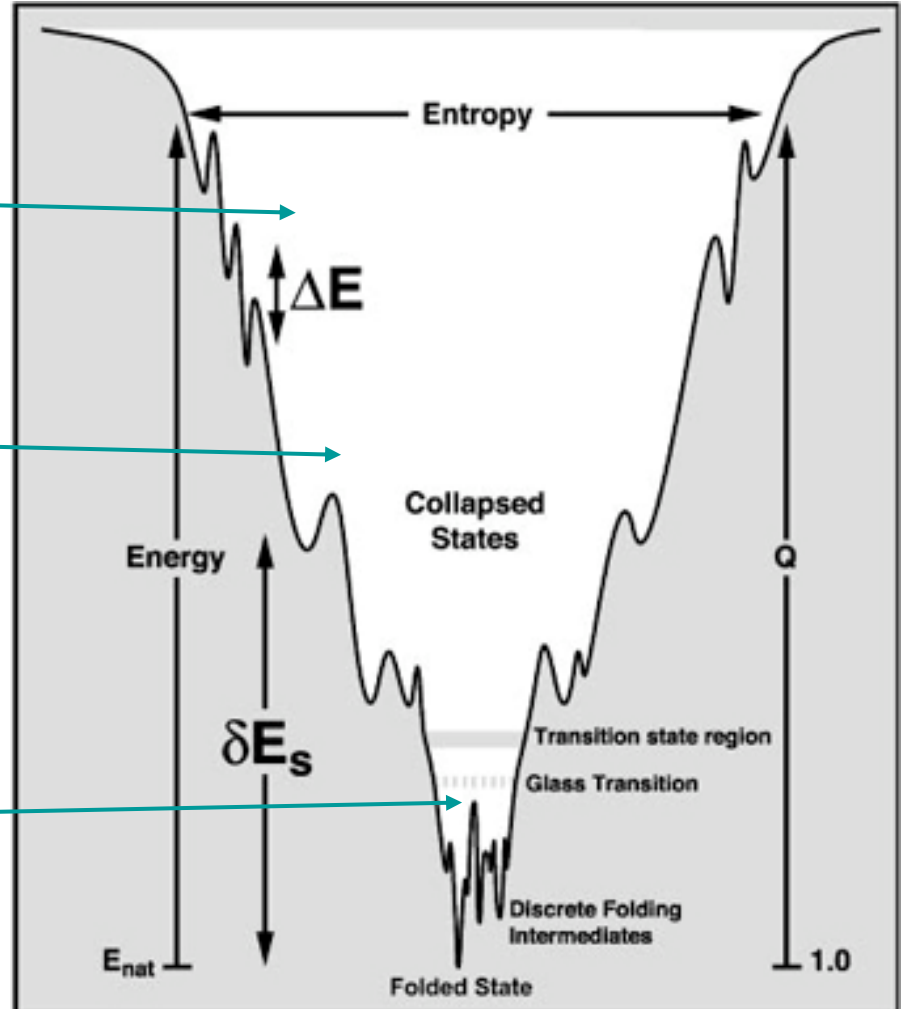
Gain of binding **energy** pulls peptide down.
Reduction of **entropy** keeps peptide up.

(高能/無序)



(低能/有序)

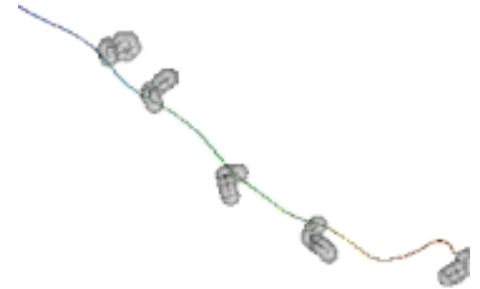
Free energy



Levinthal Paradox:

Suppose

- (i) each amino acid has 3 conformations.
- (ii) a protein consists of 100 amino acid residues with a total conformation number $3^{100} \approx 5 \times 10^{47}$.
- (iii) 100 psec (10^{-10} sec) are required for a conformational change.



Then a random search of all conformations would require around **10^{30} years**.

Nevertheless, folding of a real protein takes place in msec to sec order.

Protein folding **cannot be via a random search**.

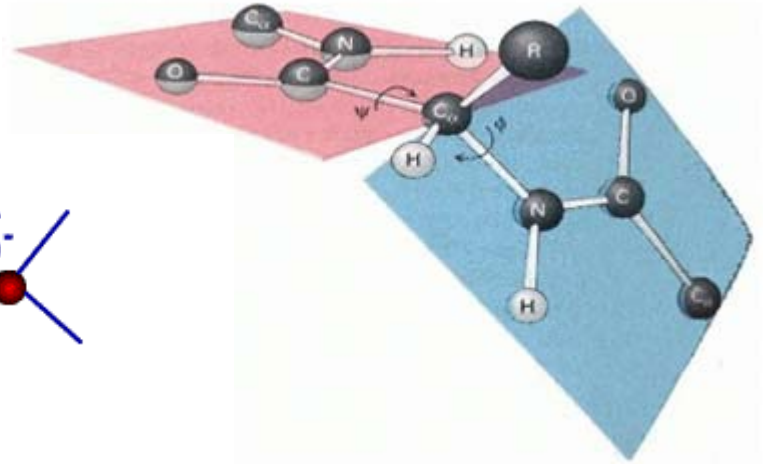
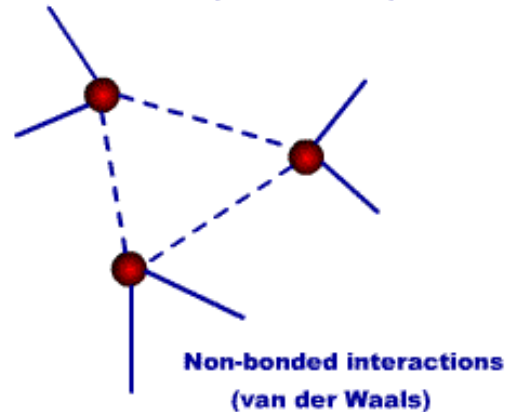
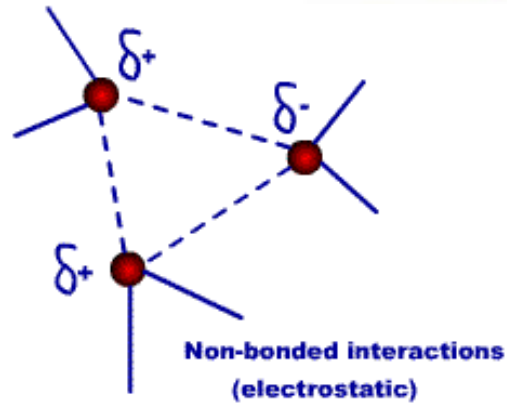
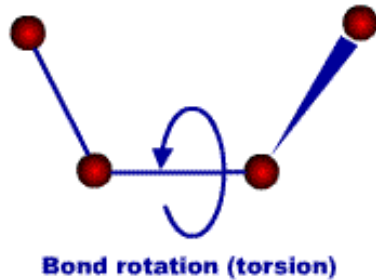
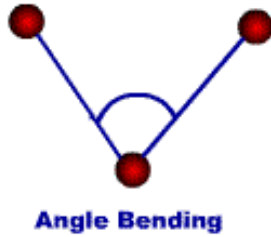
Dynamics of Protein Structures

| | |
|---------------------------------------------------|------------------------------------------------------------|
| Electron transport: | $10^{-15} - 10^{-10} \text{ s}$ |
| Bond vibration: | $10^{-14} - 10^{-12} \text{ s}$ |
| Bond rotation: | $10^{-13} - 10^{-11} \text{ s}$ |
| Rotation aromatic ring: | $10^{-9} - 10^{-3} \text{ s}$ |
| Hinge movement: | $10^{-12} - 10^{-4} \text{ s}$ |
| Folding/unfolding: | $10^{-3} - 10^0 \text{ s}$ ($10^{-9} - 10^4 \text{ s}$ *) |
| Rotational relaxation time of water: | 10^{-11} s |
| of Mb: | 10^{-8} s |
| Rotation of γ unit of F_1 : | $10^1 - 10^0 \text{ s}$ |
| Exchange of labile H: | $10^{-5} - 10^5 \text{ s}$ |
| Backbone motion during enzyme Catalysis (NMR): | $10^{-1} - 10^1 \text{ s}$ |
| Myosin movement: | 10^{-8} m s^{-1} |
| Diffusion 1 nm particle: | $25 \times 10^{-6} \text{ ms}^{-1}$ |
| Diffusion CO in Mb: | $10^{-9} - 10^2 \text{ s}$ |

* <http://pcwww.liv.ac.uk/~volk/folding/previous.htm> (Fast events in protein folding) and mutated proteins.

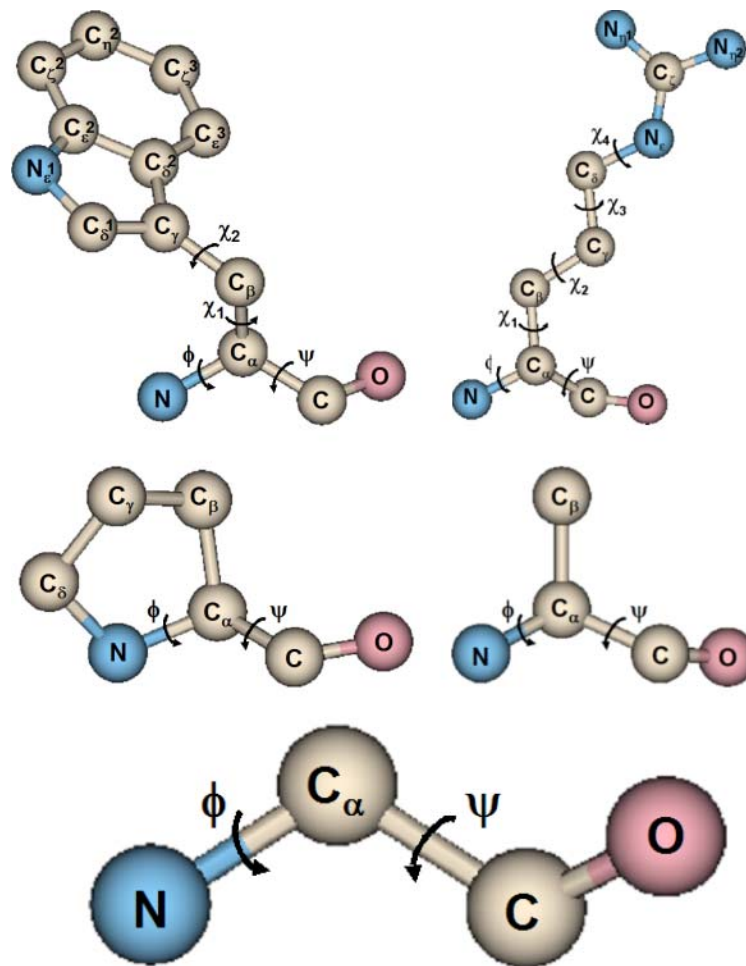
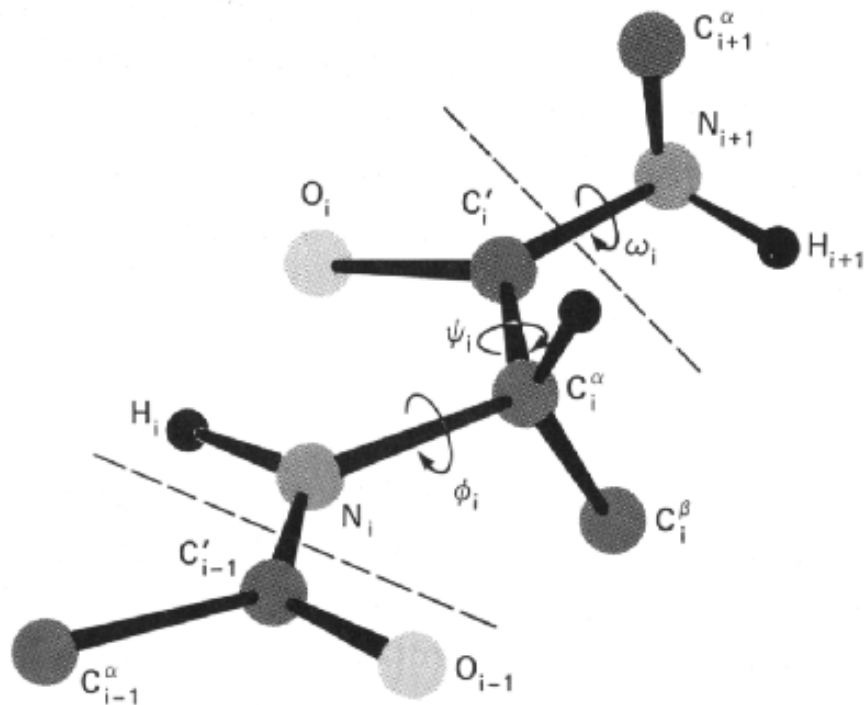
Force and energy in proteins

Degrees of freedom in a protein

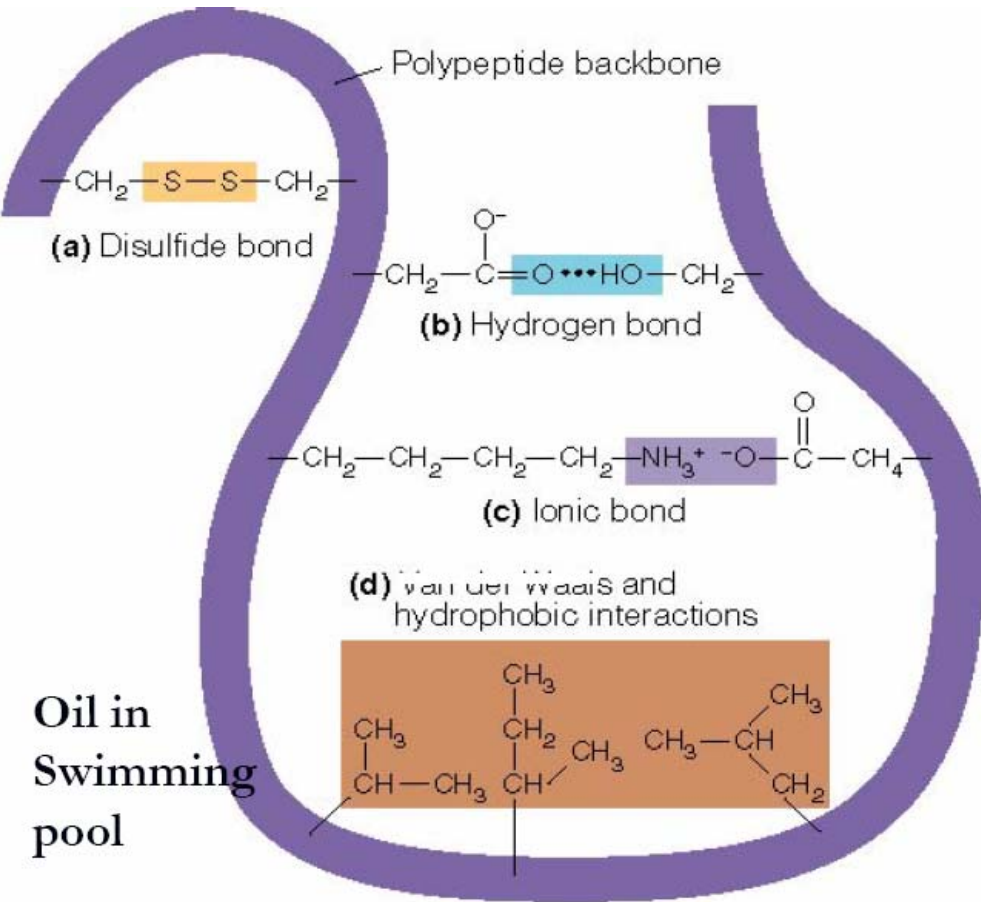


Most flexible degrees of freedom:

torsion Φ between N-C $_{\alpha}$
torsion Ψ between C-C $_{\alpha}$



Forces acting on Proteins



- (a) Disulfide bond
- (b) Hydrogen bond
- (c) Electrostatic interaction
- (d) van der Waals interaction
- (e) Hydrophobic interaction
- (f) Intrinsic properties
(conformational preference)

Hydrophobicity:

the dominant force in protein folding (Dill, 1990).
The presence of water is important for this interaction.

Force scales

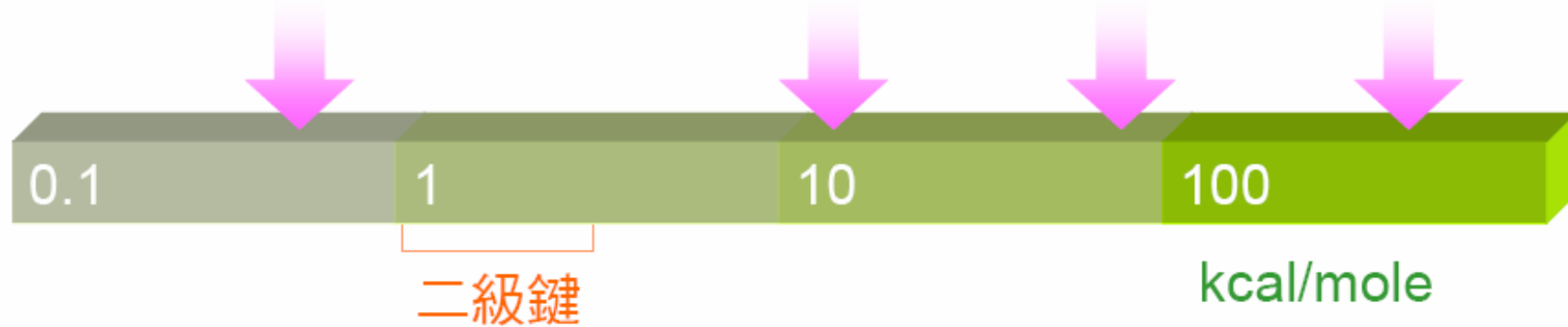
Adapted from Alberts et al (2002) Molecular Biology of the Cell (4e) p.53

室溫

ATP 水解

共價鍵

葡萄糖氧化



凡得瓦爾力

0.1 kcal/mole

疏水鍵

1 kcal/mole

氫鍵

1 kcal/mole

離子鍵

3 kcal/mole

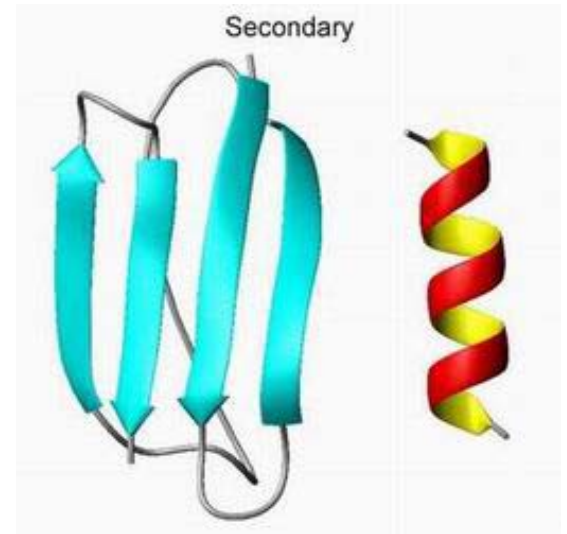
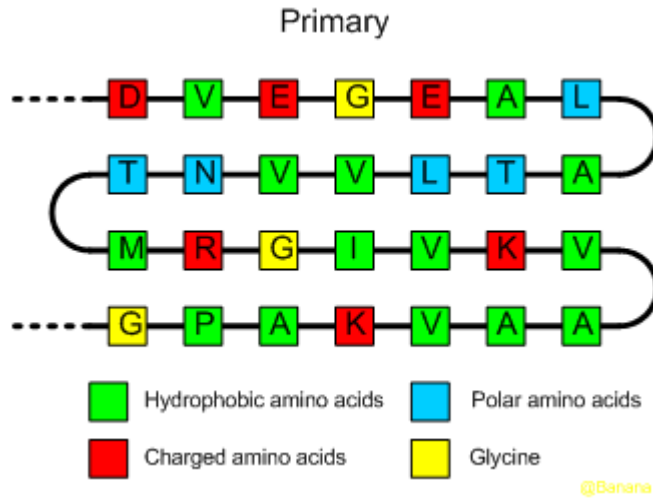
共價鍵

90 kcal/mole

For room temperature 300K:
 $1kT = 1.38 \times 10^{-23} \text{ J/K} \cdot 300 \text{ K}$
 $\approx 4.1 \times 10^{-21} \text{ J}$
 $\approx 2.5 \times 10^{-19} \text{ eV}$
 $\approx 0.6 \text{ kcal / mole}$
 $\approx 2.5 \text{ kJ / mole}$

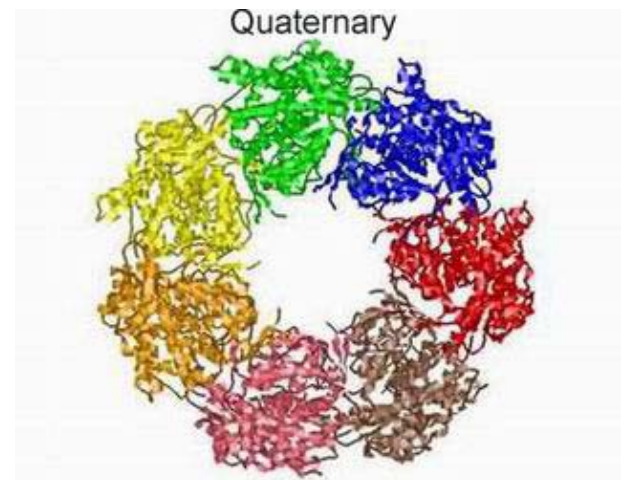
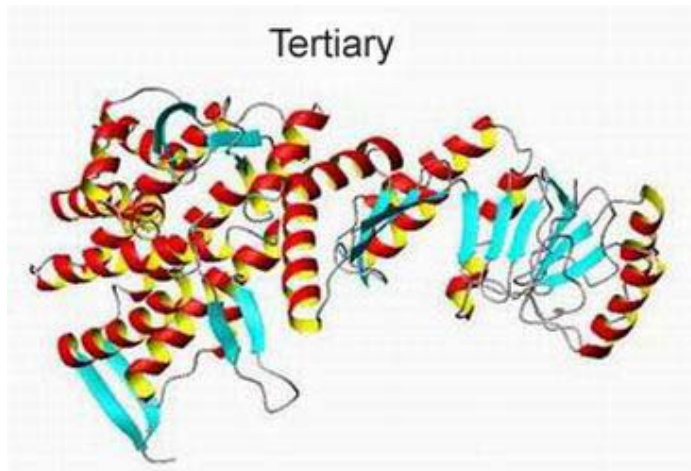
Models and Methods

Hierarchy of protein structure

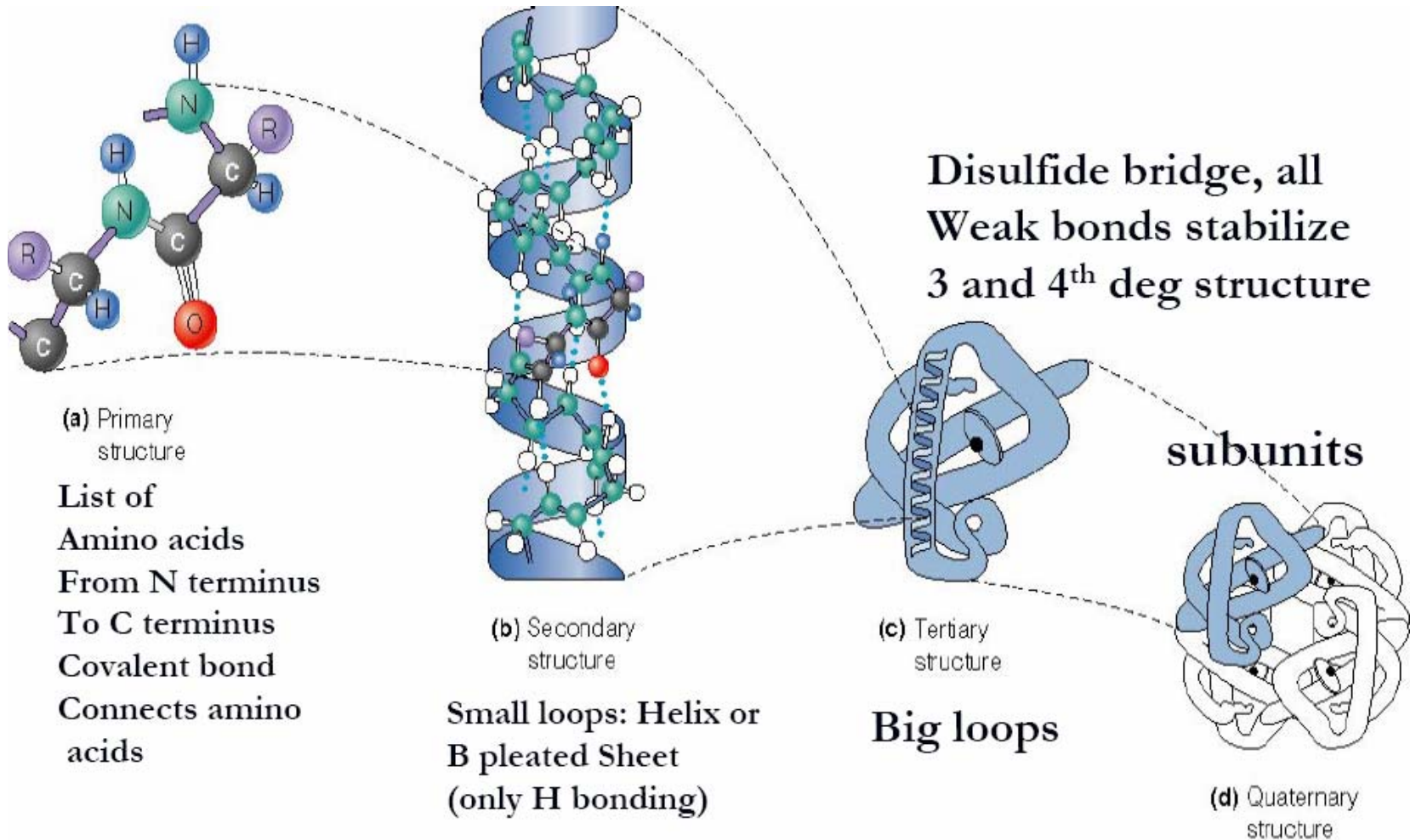


β sheets

α helix



Hierarchy of protein structure



Model for protein folding:

(i) **Framework model** describes a step-wise mechanism to greatly narrow the conformational search. This involves a **hierarchical assembly** whereby local elements of secondary structure are formed according to the primary sequence, but independent from tertiary structure. These elements then **diffuse** until they collide, whereupon they coalesce to form the tertiary structure.

(ii) **Nucleation model** suggests that tertiary structure forms as an **immediate consequence** of the formation of secondary structure. Nucleation occurs through the formation of native secondary structure by only a few residues (e.g. a beta-turn, or the first turn of an alpha-helix), and structure propagates out from this nucleus.

(iii) **Hydrophobic collapse model** hypothesises that the native protein conformation forms by rearrangement of a compact collapsed structure. Hydrophobic collapse to form a molten globule therefore constitutes an early step in the folding pathway.

(iv) **Nucleation-condensation model** suggests that a diffuse folding nucleus is formed in (ii) and consolidated through the **transition state**, concomitant with tertiary structure formation.

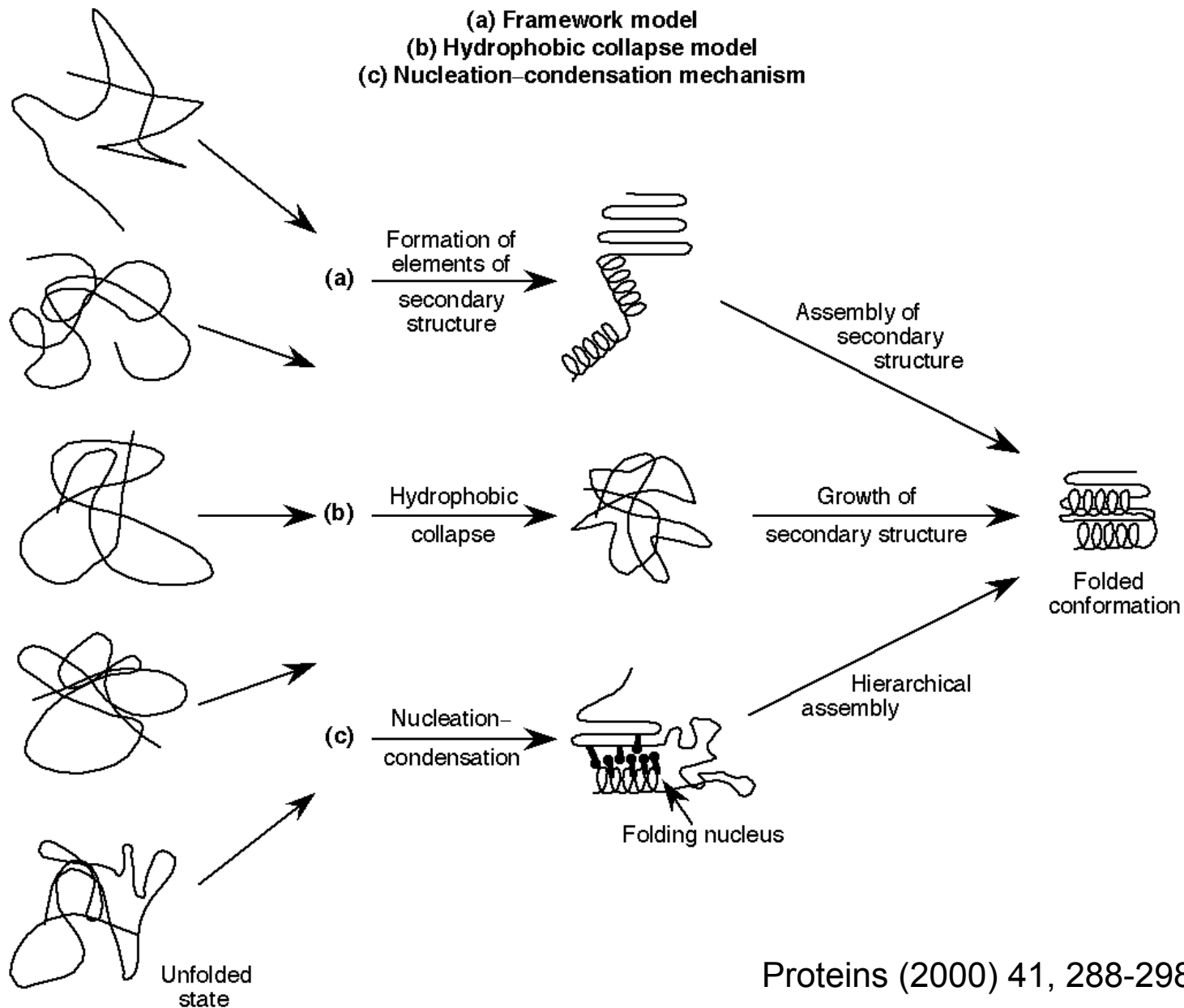
(i), (ii), and (iv) suggest the formation of **kinetic intermediates**, whereas (iii) does not. (A. R. Fersht. *Curr. Opin. Struc. Biol.*, 1997, 7, 1, 3-9.)

Models for protein folding:

(a) Framework model

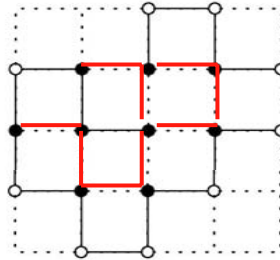
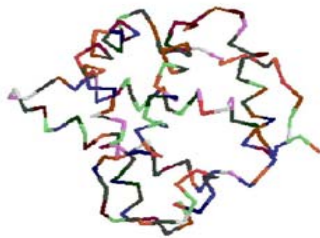
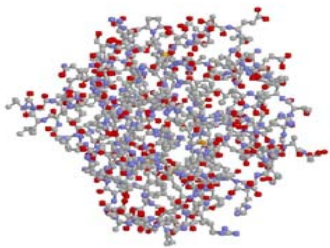
(b) Hydrophobic collapse model

(c) Nucleation–condensation mechanism



HP Lattice Models

- Put each residue in the polypeptide to a site of a 2D (or 3D) lattice. A protein structure is represented as a path through lattice points. Empty lattice cells represent water.
- The dominant interaction in protein folding is residue with water: **hydrophobic** (厭水) and **hydrophilic** (親水), represented by **H** and **P**. In the native conformation of a protein, hydrophobic residues tend to be **buried**, hydrophilic residues tend to be **exposed** to water.
- Dividing residues into H and P groups, a polypeptide becomes a binary sequence.



- hydrophobic amino acid
- hydrophilic amino acid
- Covalent bond
- H-H contact

Goal: maximize the number of H-H contacts

Molecular Dynamics (MD)

Newton's equation of N atoms of mass m_i at \mathbf{r}_i under force $\mathbf{F}_i(t)$:

$$m_i \frac{d^2 \mathbf{r}_i(t)}{dt^2} = \mathbf{F}_i(t), \quad (i = 1, 2, \dots, N).$$

$$\mathbf{r}_i(t + \Delta t) = 2\mathbf{r}_i(t) - \mathbf{r}_i(t - \Delta t) + \frac{1}{m_i} \mathbf{F}_i(t) \Delta t^2 + O(\Delta t^4).$$

(finite difference method)

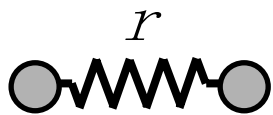
$$\mathbf{F}_i = -\nabla_i V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N),$$

Δt should be approximately 1/10 the time of the fastest motion in the system. The fastest motion in a protein is the stretching of light atoms, e.g., O-H, C-H, with periods around 10^{-14} sec. Thus Δt is usually taken as 1 fsec.

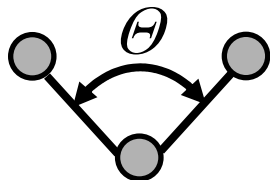
Huge number of water molecules are considered in MD simulations. Their non-bonded interactions (van der Waals, electrostatic interactions) increases in order of N^2 (N is the number of atoms).

Usually only a few tens of nanoseconds simulation is performed.

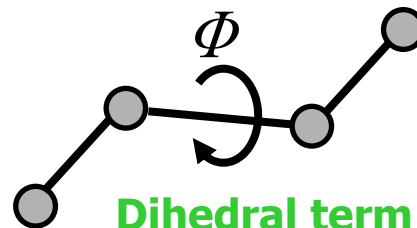
Energy function in Molecular Dynamics



stretching term



bending term



Dihedral term

$$V_{\text{total}} = \sum_{\text{bonds}} K_b (r - r_0)^2 + \sum_{\text{angles}} K_\theta (\theta - \theta_0)^2 + \sum_{\text{dihedrals}} K_\phi [1 + \cos(n\phi - \gamma)]$$

$$+ \sum_{\text{Hbonds}} \left(\frac{C_{ij}}{r_{ij}^{12}} - \frac{D_{ij}}{r_{ij}^{10}} \right)$$

H-bonding term

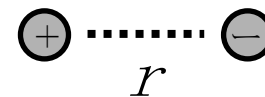
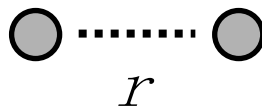
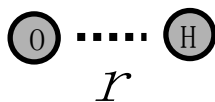
$$+ \sum_{\text{van der Waals } i, j \text{ pairs}} \left(\frac{A_{ij}}{r_{ij}^{12}} - \frac{B_{ij}}{r_{ij}^6} \right)$$

Van der Waals term

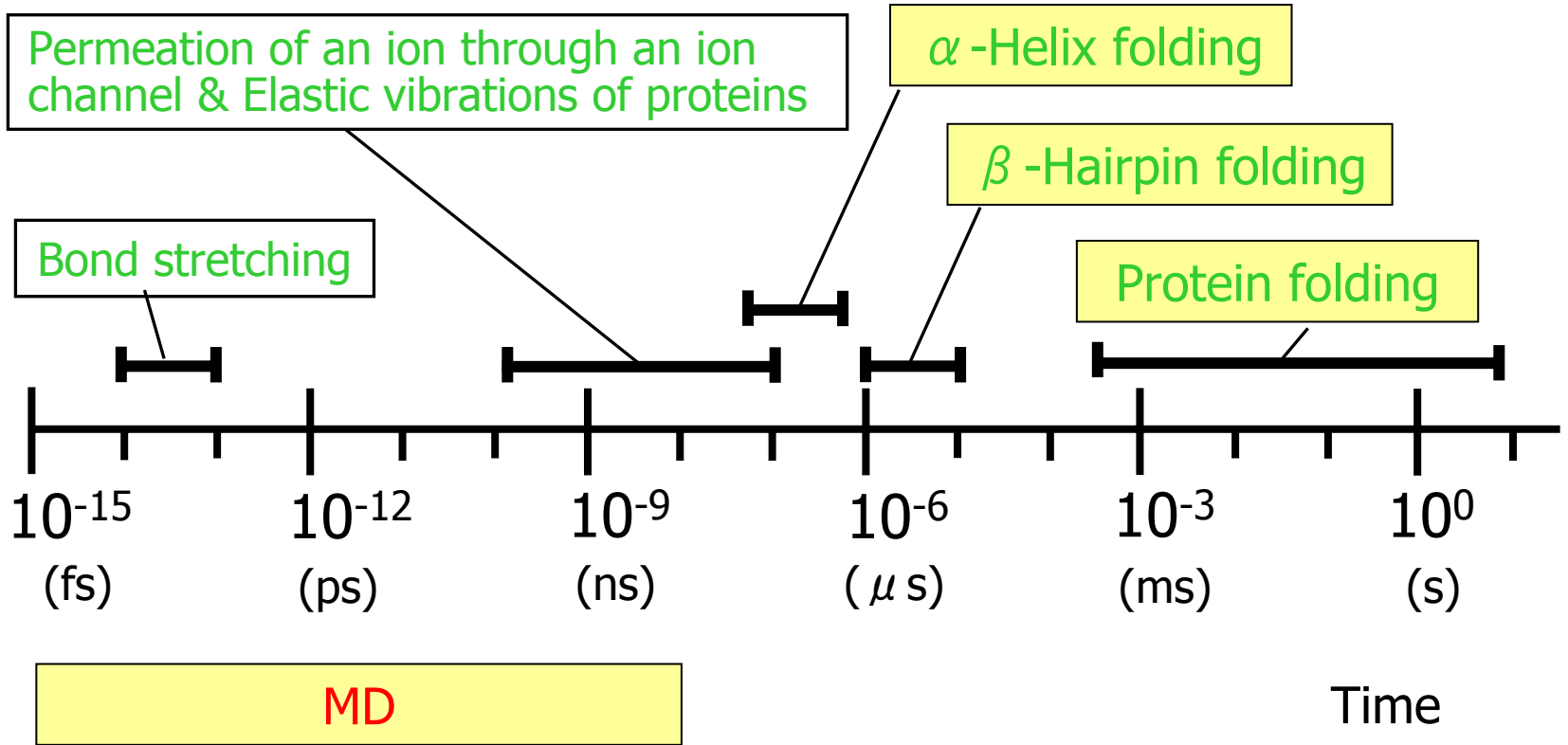
$$+ \sum_{\text{electrostatic } i, j \text{ pairs}} \frac{q_i q_j}{\epsilon r_{ij}}$$

Electrostatic term

The most time demanding part.



Time scales of MD and protein motions

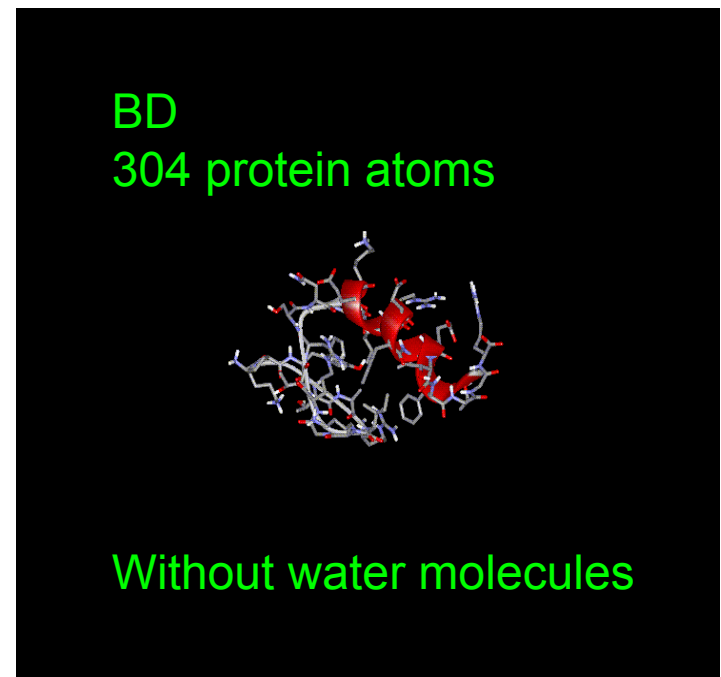
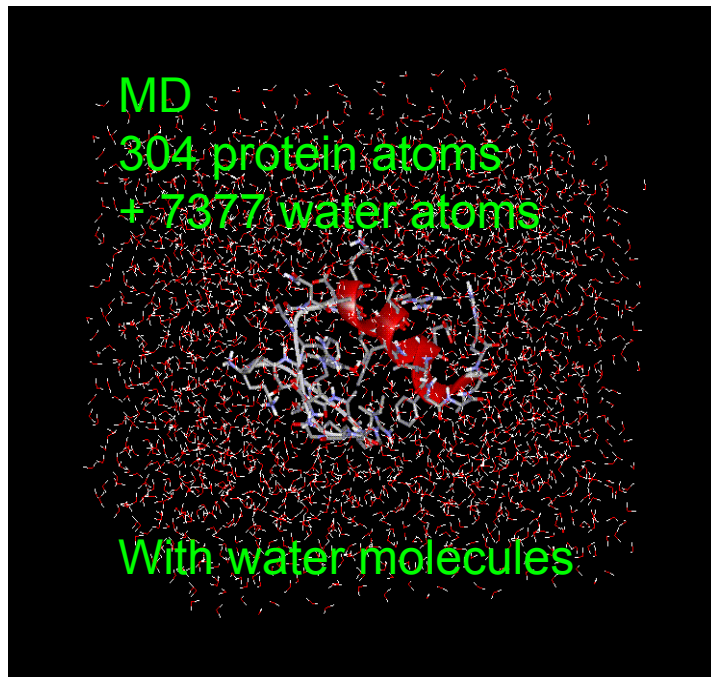


MD is hard to simulate a whole process of a protein folding.

Brownian Dynamics (BD)

The dynamic contributions of the solvent are incorporated as a **dissipative random force**.

Thus, water molecules are not treated explicitly.



Brownian Dynamics

The Langevin equation can be expressed as

$$m_i \frac{d^2 \mathbf{r}_i}{dt^2} = -\zeta_i \frac{d\mathbf{r}_i}{dt} + \mathbf{F}_i + \mathbf{R}_i$$

with the viscosity η of water, the frictional coeff. $\zeta_i = 6 \pi a_i \eta$ (determined by Stokes' law), a Stokes radius a_i of atom i , the systematic force \mathbf{F}_i on atom i , a random force \mathbf{R}_i on atom i with zero mean $\langle \mathbf{R}_i(t) \rangle = 0$ and a variance $\langle \mathbf{R}_i(t) \mathbf{R}_j(t') \rangle = 6 \zeta_i kT \delta_{ij} \delta(t-t')$ which accounts for the solvent effects.

In the overdamped limit, we have

$$\zeta_i \frac{d\mathbf{r}_i}{dt} = \mathbf{F}_i + \mathbf{R}_i$$

Integrating this equation yields the Brownian Dynamics:

$$\mathbf{r}_i(t + \Delta t) = \mathbf{r}_i(t) + \frac{\mathbf{F}_i(t)}{\zeta_i} \Delta t + \sqrt{\frac{2k_B T}{\zeta_i}} \Delta t \boldsymbol{\omega}_i$$

with a random noise vector $\boldsymbol{\omega}_i$ obtained from Gaussian distribution.

Difficulties in Protein folding

- Very time-consuming
- Not clear if the energy function V used is precise enough.

General problems in biophysics:

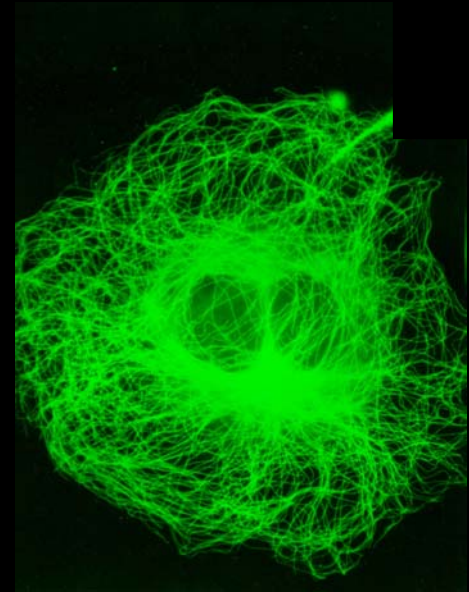
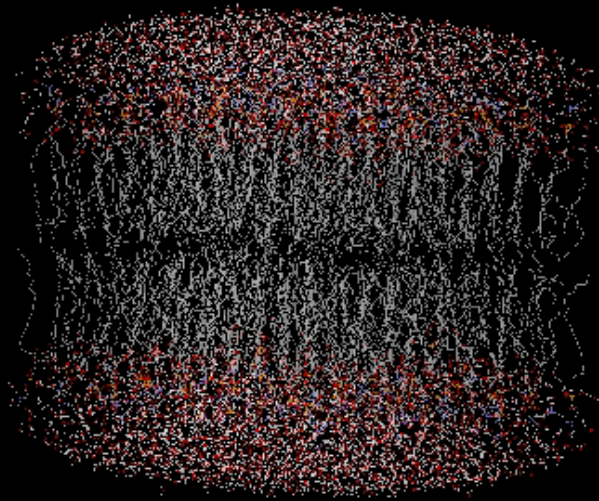
Physics is an art of reduction which reduces systems of high degree of freedom to (effective) models with low degree of freedom, via **symmetry** or **constrains**.

However, symmetry and constrains are **weak in biosystems**.

Hard to truncate or reduce problems to effective simple models.

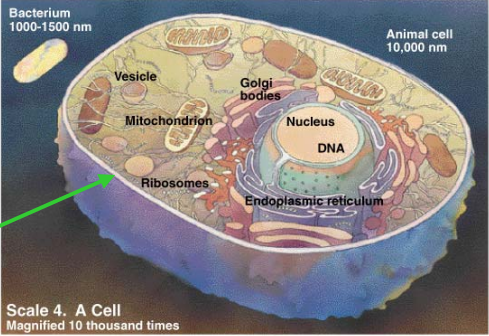
Example: potential V used in protein folding.

3. Mechanics (membrane & filaments)

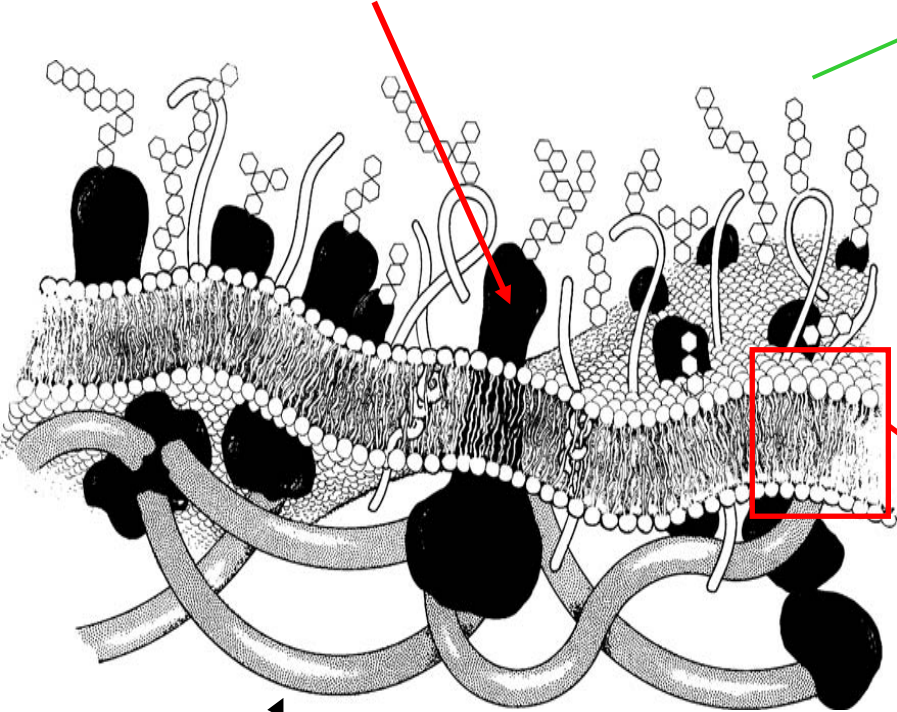


Cell membranes

Structure of a biological membrane



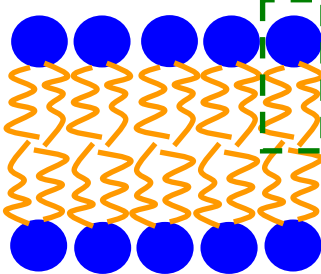
Inclusion



Primary building blocks

- Lipid (脂質), or fat
- Protein (蛋白質)

lipid bilayers



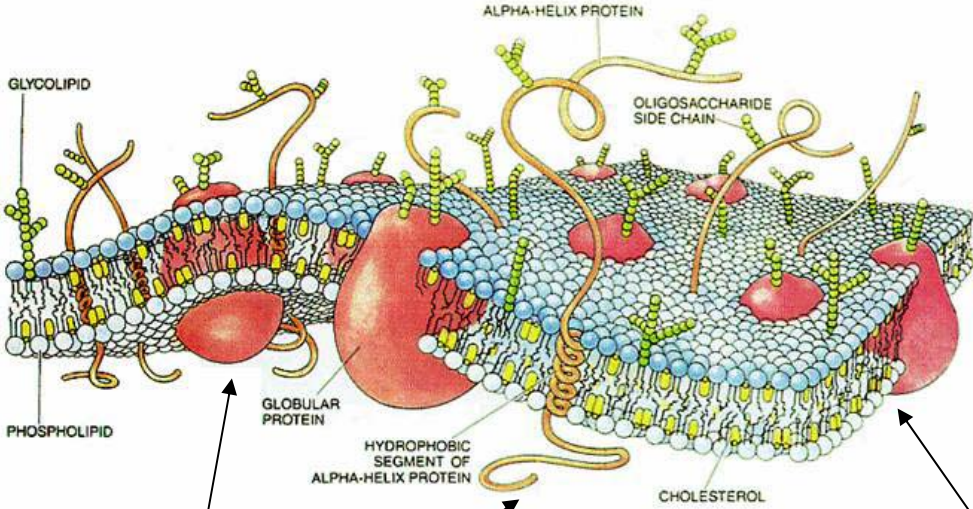
~ 4nm

Filament (骨架)

Lipids are in liquid crystalline (2D liquid) phase

Structure of an artificial membrane

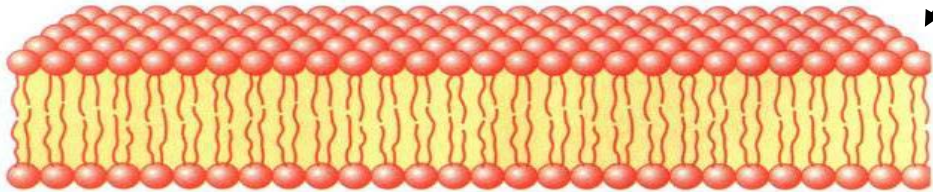
Biological Membrane



Inclusion

Lipid bilayer

Artificial Membrane

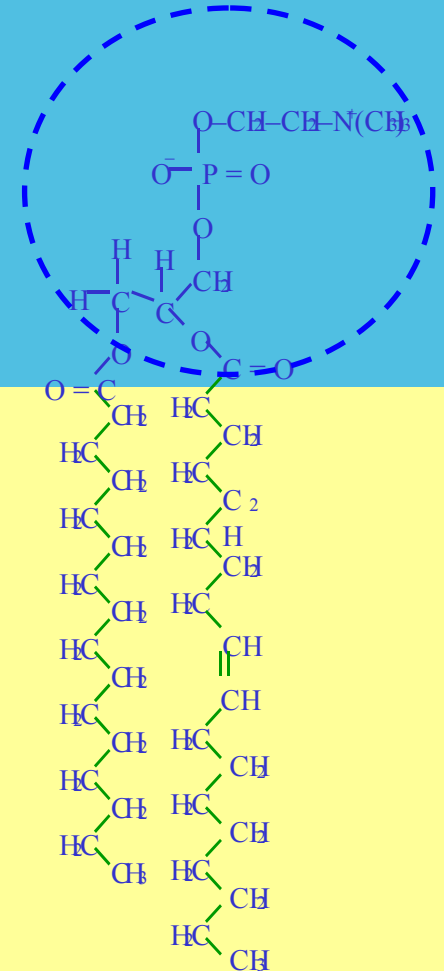
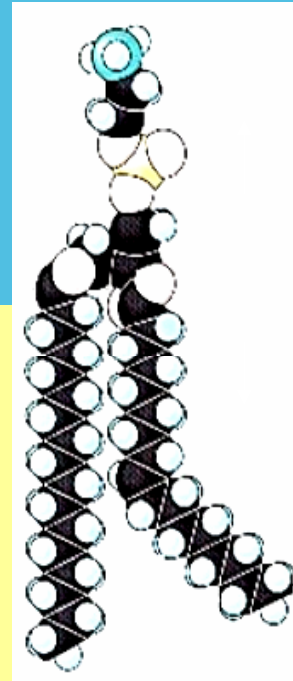
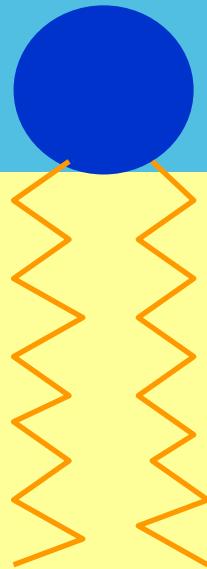


Lipid is an amphipathic molecules (molecules containing both polar and nonpolar parts).

水

Headgroup:
polar
hydrophilic 親水
(water-loving)

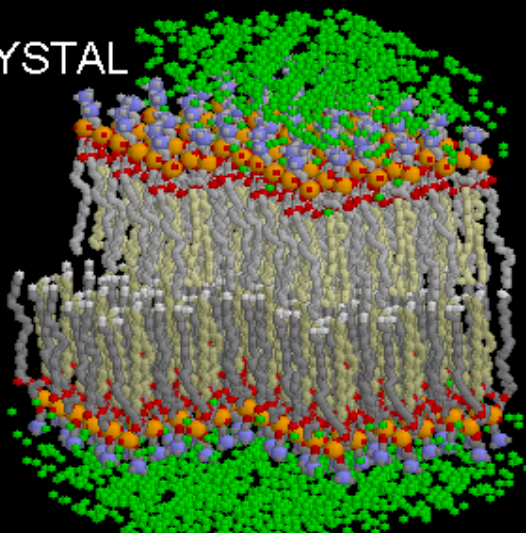
Tail (or acyl chain):
non-polar
Hydrophobic 厭水
(water-fearing)



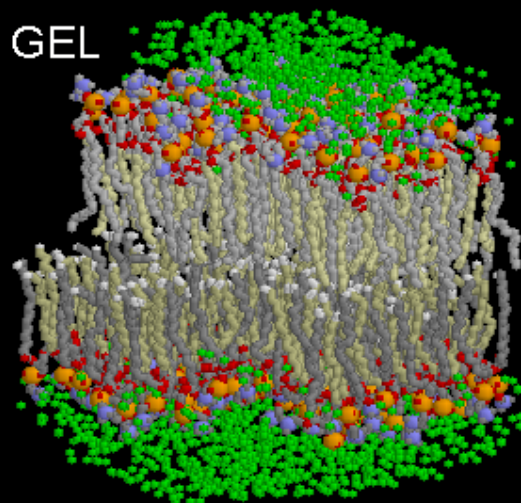
Different (microscopic) phases

(These microscopic 'phases' are differently defined from the assembly 'phases' in the previous page)

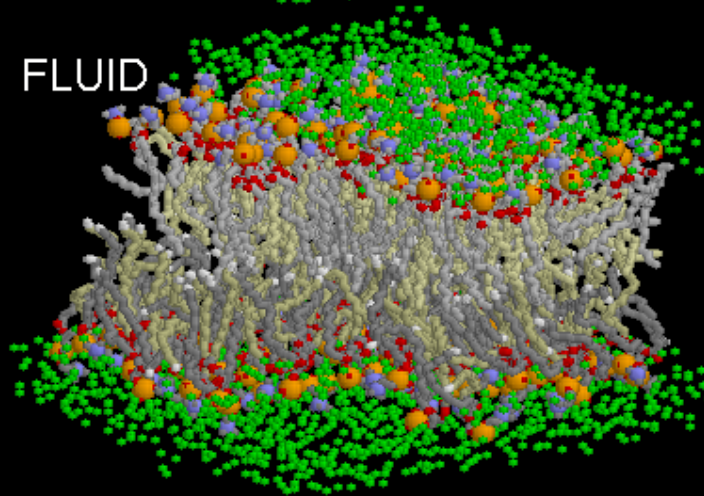
CRYSTAL



GEL



FLUID



of Phosphatidyl Choline Bilayer

Carbon/Palmitic Oleic

Nitrogen Oxygen Phosphorus

Water Oxygens

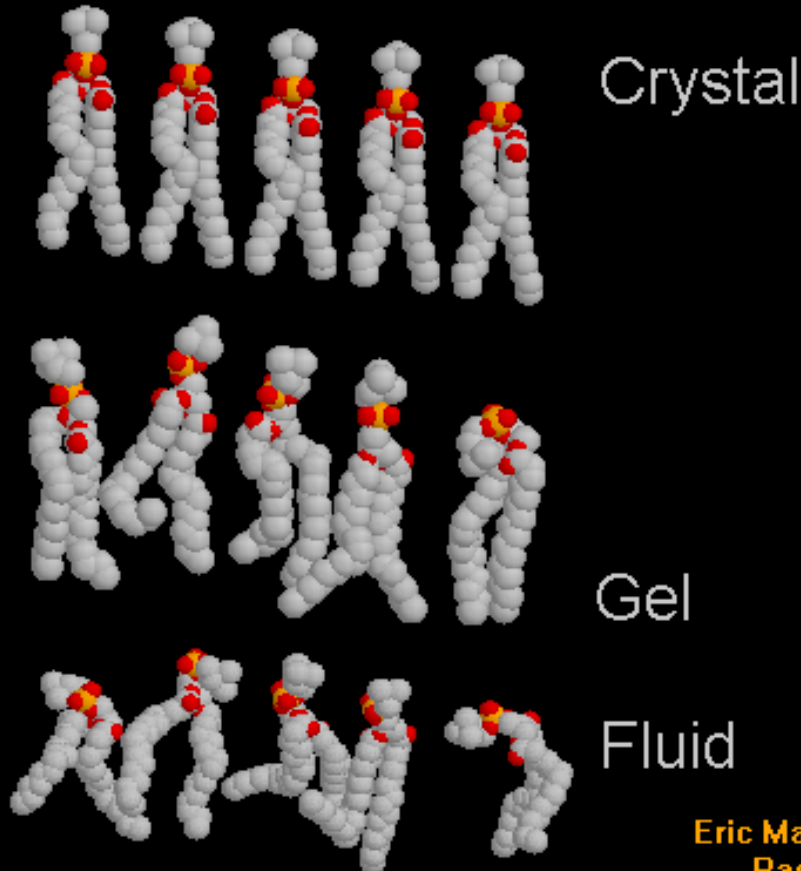
H Heller, M Schaefer, K Schulten,

J Phys Chem 97:8343, 1993.

RasMol Image by E Martz

Phase transition

gel → liquid crystal transition
lattice melting & chain melting



Gel phase

- solid-ordered (so) phase
- not important to cell function
- ordered (all-trans) chains arranged on a crystalline lattice

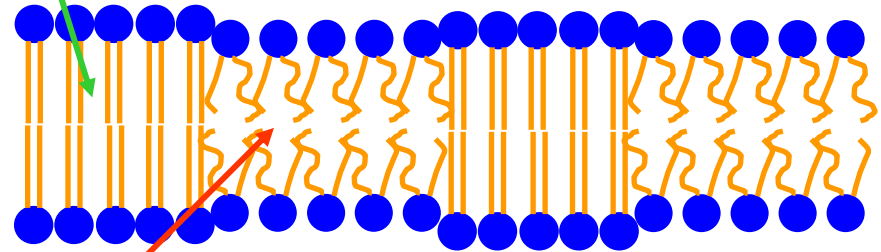
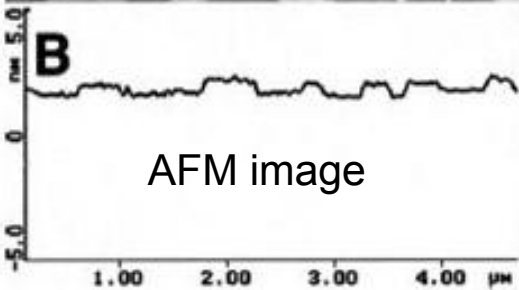
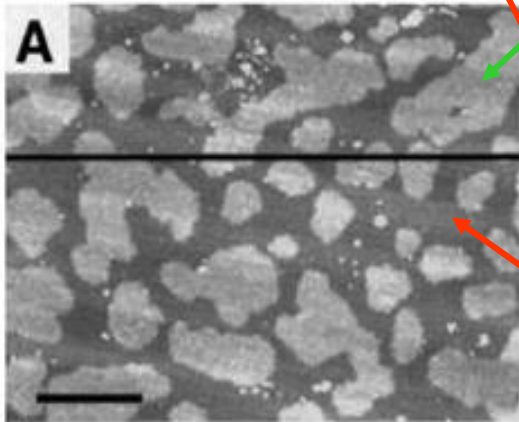
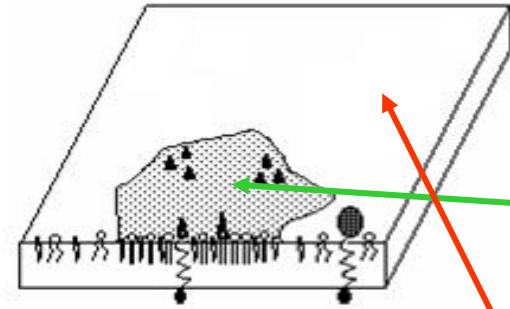
Liquid crystalline phase

- liquid-disorder (ld) phase
- important to cell function
- disordered (many gauches) chains with random lateral arrangement.
- Lateral diffusion constant:
 $D \sim 10^{-12} \text{ m}^2/\text{sec}$

Eric Martz with
RasMol

Lipid rafts (浮冰)

- Liquid-ordered (lo) phase (like 冰山):
- Enriched in Cholesterol and sphingolipids
 - Cholesterol(膽固醇)-associated phase
 - lateral motion: slower than Id (half of Id's)
 - highly ordered chain (only a few gauches)



- Liquid-disorder (ld) phase (like 海洋):
- disordered with many gauches chains
 - with random lateral arrangement

Intrinsic curvature

Elastic free energy density

$$H_{el} = \int dA \left(\gamma + \frac{\kappa}{2} [c - c_0]^2 \right)$$

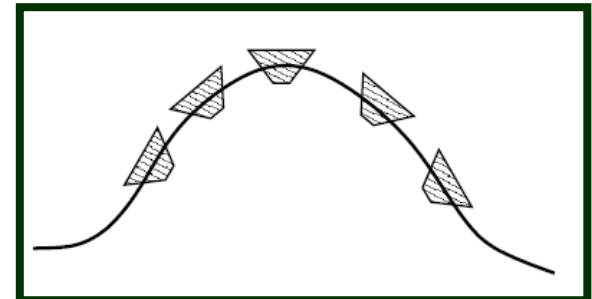
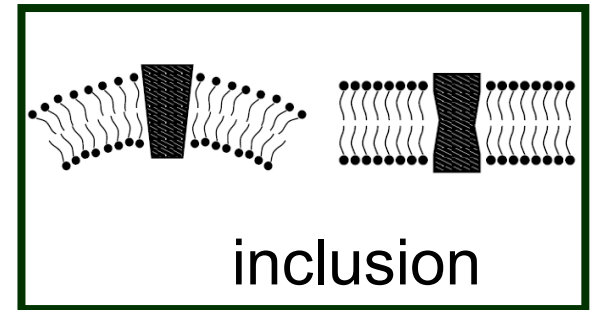
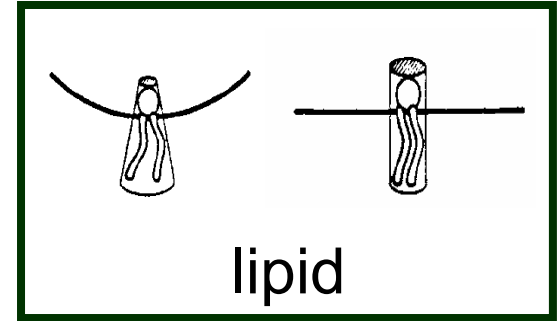
γ : surface energy $\sim kT/(nm)^2$

κ : curvature elastic modulus $\sim 10 kT$

c : curvature of the membrane

c_0 : intrinsic curvature

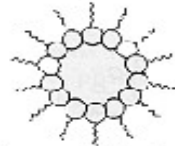
Curvature-induced interaction



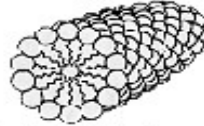
Different (macroscopic) phases



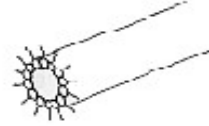
Micelle



Inverse Micelle



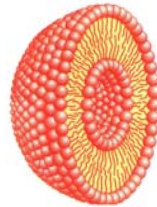
Prolate Micelle



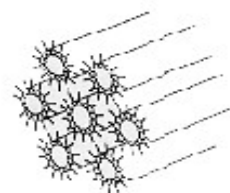
Inverse Prolate Micelle



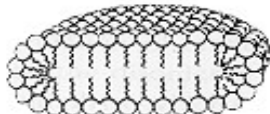
Hexagonal phase
Normal



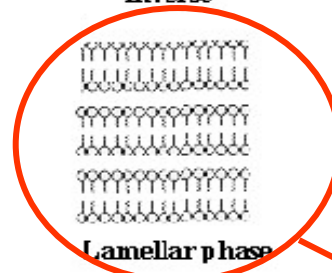
Lamellar vesicle
(multi- or uni-layer)



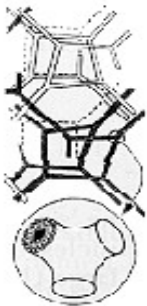
Hexagonal phase
Inverse



Oblate Micelle,
bilayered fragments



Lamellar phase

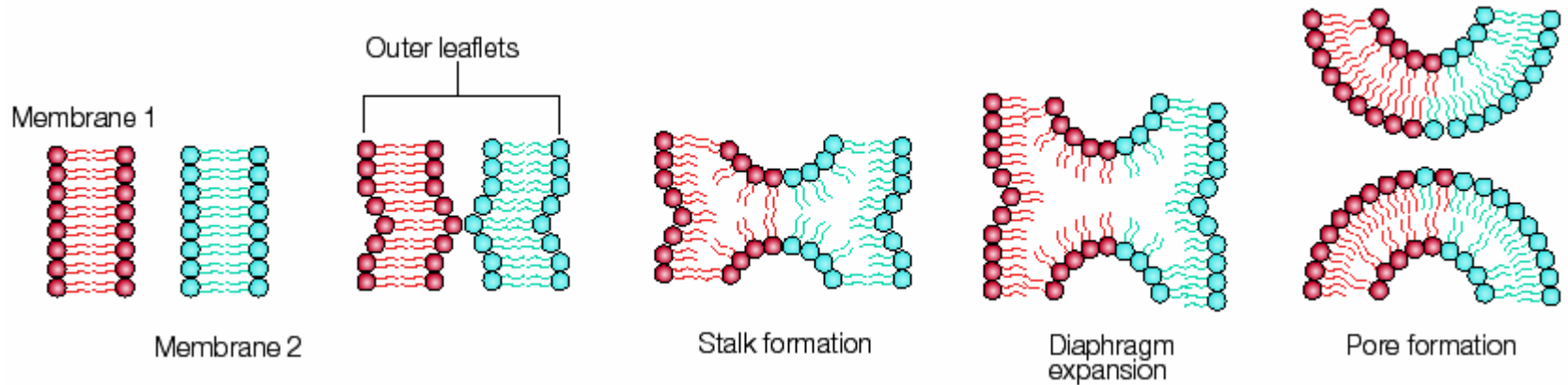


Cubic phases



Lamellar phase
important for
cell membranes.

The stalk-pore model for cell fusion (Example: viruses with cells³⁷)



The polar heads of phospholipids in the outer leaflets of plasma membranes of two cells approach one another.

The energy required for this approach of negatively charged leaflets might be facilitated by receptor–ligand interactions (not shown).

The outer leaflets fuse to form a stalk-like structure, which expands to create a diaphragm.

Forces generated from the extension of the diaphragm promote fusion of the inner leaflets and the creation of a fusion pore.

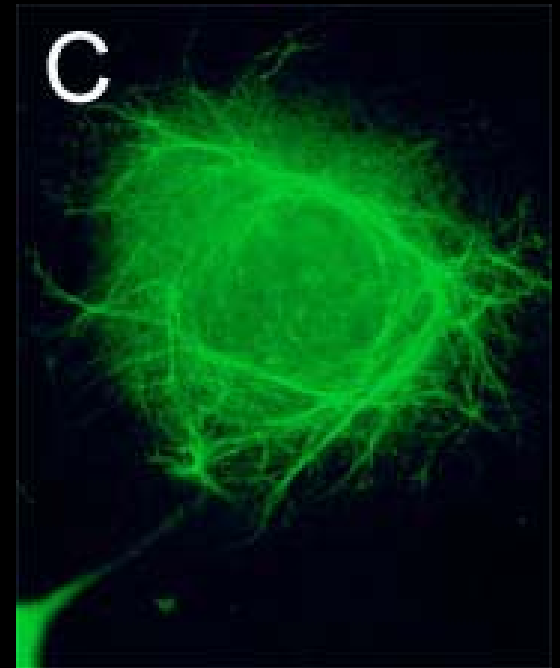
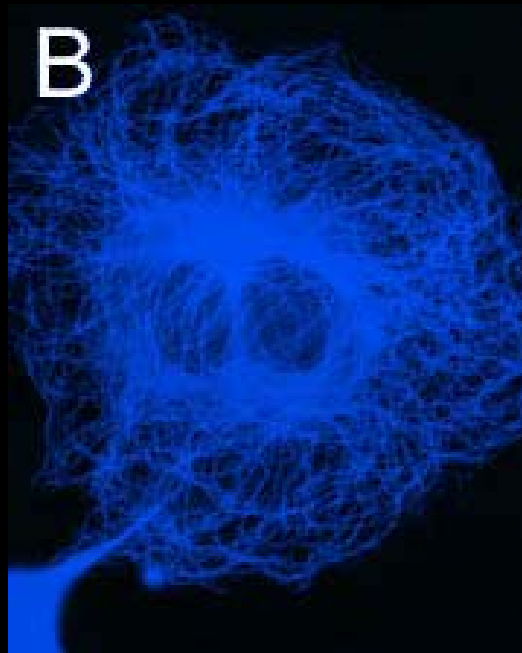
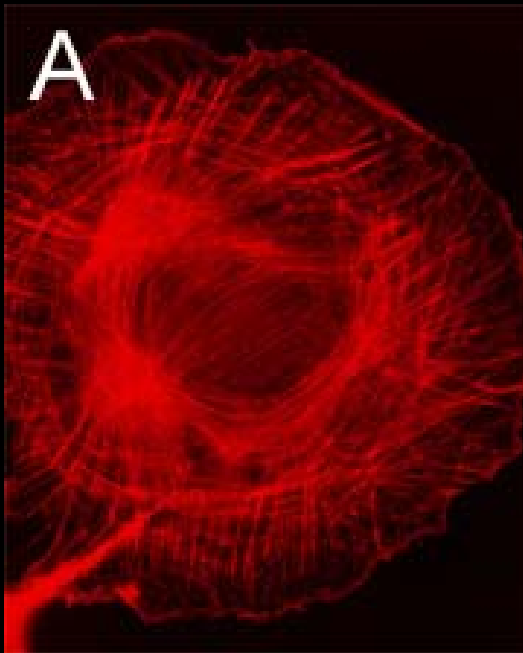
Functions of filaments

The three polymer systems of a fibroblast cytoskeleton

A, the actin cytoskeleton

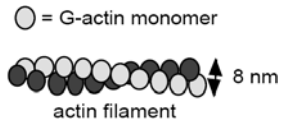
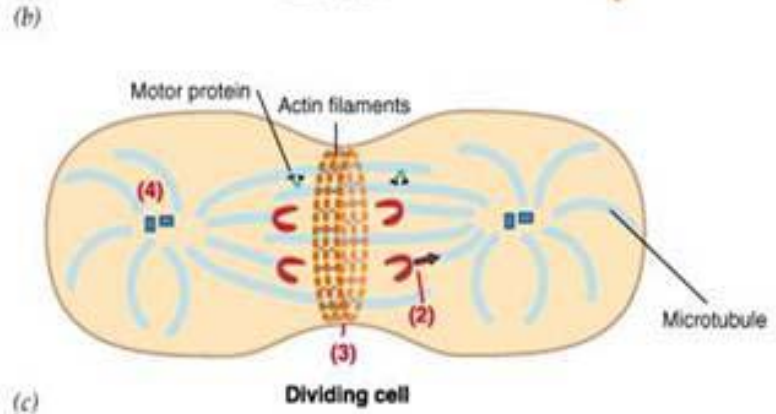
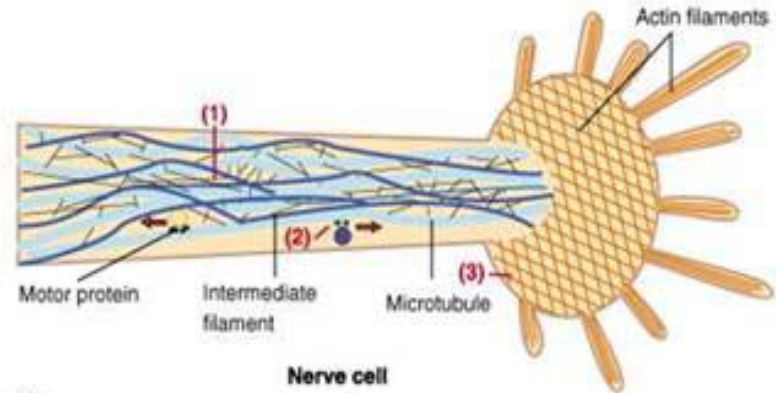
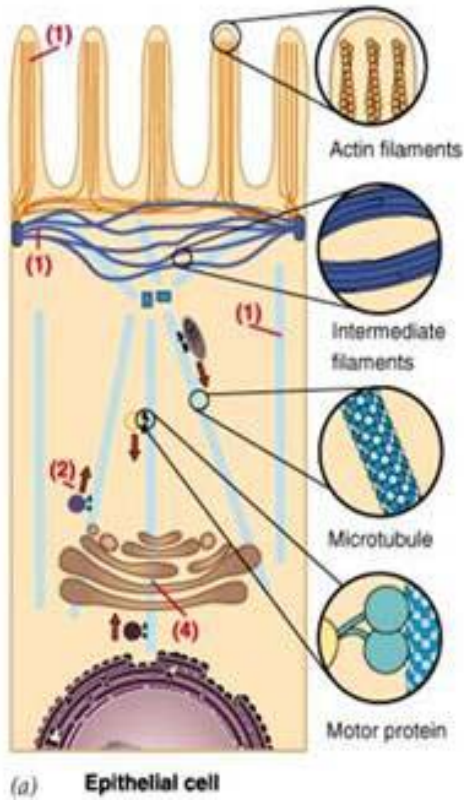
B, the microtubule cytoskeleton

C, the intermediate filament cytoskeleton

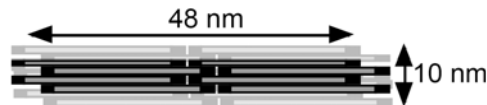


Different kinds of filaments

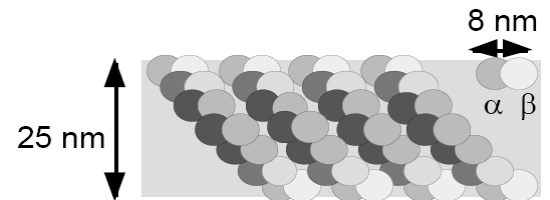
- (1) Structure and Support (2) Intracellular Transport (3) Contractility and Motility (4) Spatial Organization



Actin



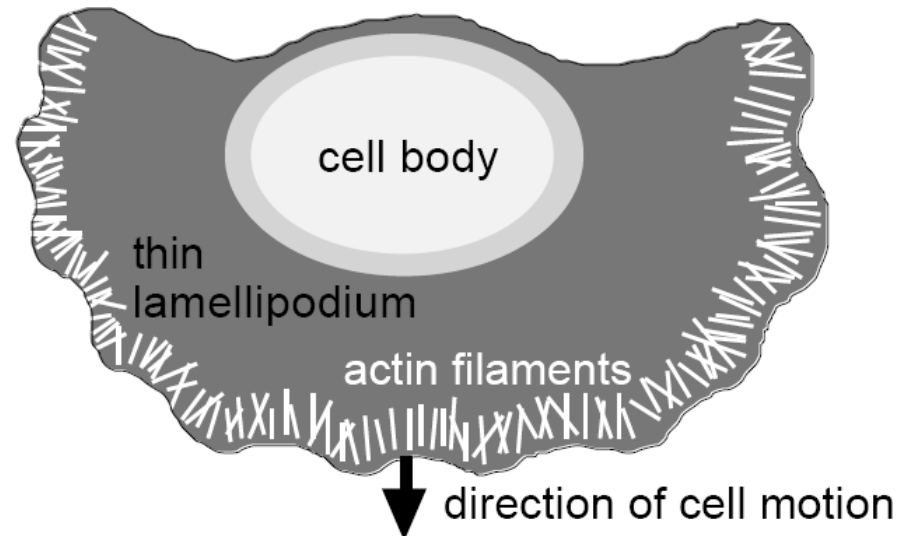
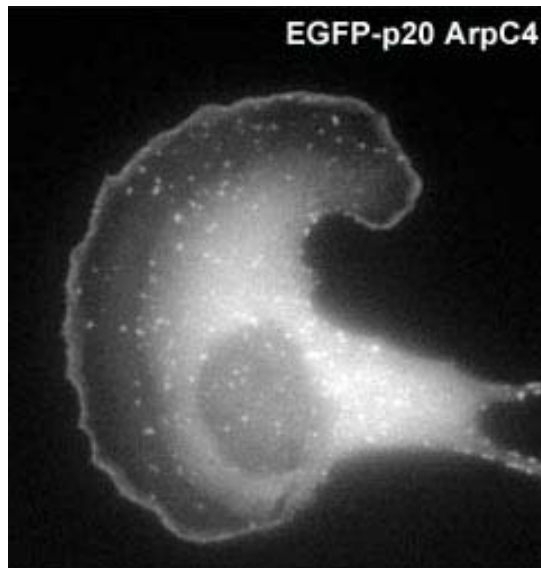
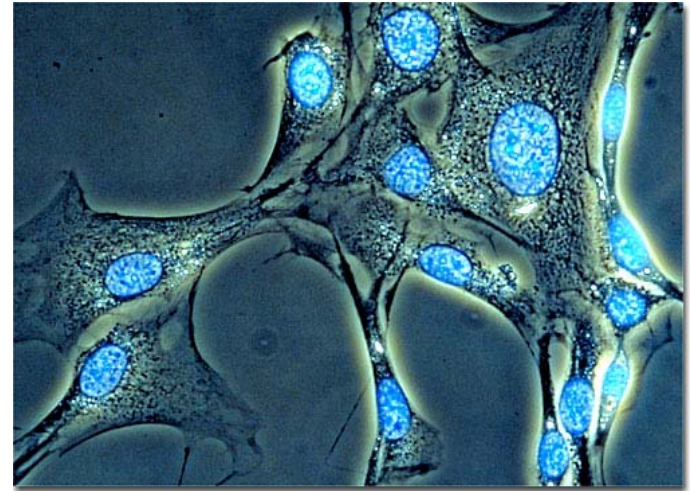
Actin



Microtubules

Actin – cell movement

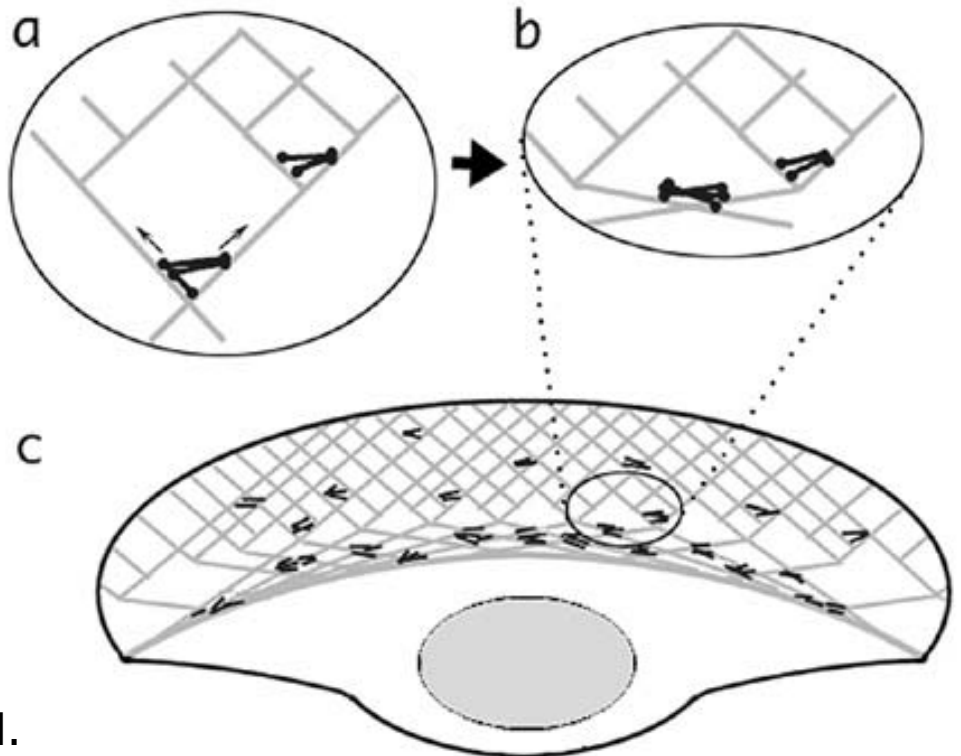
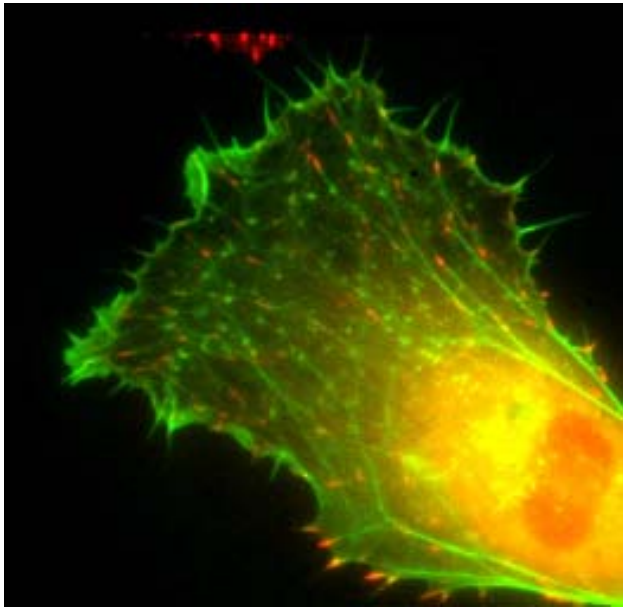
- Fibroblasts (纖維母細胞) move along a substrate, adhering by the sheet-like **lamellipodium** (片足) direction of cell motion.
- The leading edge is actin-rich, with plus ends at the cell boundary.



Actin – cell movement

In keratocytes (角膜細胞), actin filaments move back through the cell body at roughly the speed of the cell, such that a given position on a filament remains roughly stationary with respect to the substrate (Theriot et al, 1991).

Compression of an actin-myosin network into a bundle at the lamellipodial of a fish keratocyte is coupled to forward translocation.

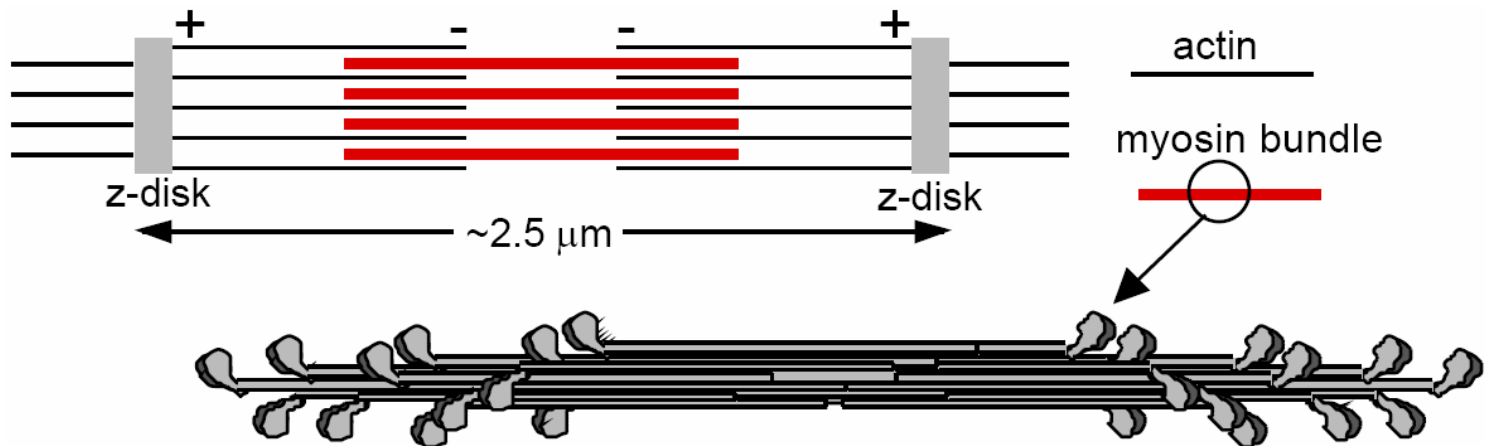


Adhesion points are colored red.

Actin - track of motor protein

Myosins move on actin (to plus end)

- Actin and myosin may be organized into highly cooperative structures in our muscles.
- Thick filaments are bundles of more than 100 myosins.
- Walking towards the plus end of the actin, myosin pulls the minus ends of the filament towards one another, contracting the muscle along the horizontal direction.



Microtubules – highways in cells

Cells need a mechanism for directed transport of materials from production to consumption site.

NOT diffusion via $\langle x^2 \rangle = 2Dt$ with $D = k_B T / 6\pi\eta R$.

For 10 nm proteins, $D = 10^{-10} \text{ m}^2/\text{s}$ in water and $10^{-14} \text{ m}^2/\text{s}$ in lipids.

Taking $D = 10^{-12} \text{ m}^2/\text{s}$,

if $\langle x^2 \rangle = 1 \text{ } \mu\text{m}^2$

→ $t \sim 1$ second (ok for local transport)

if $\langle x^2 \rangle = 1 \text{ m}^2$

→ $t \sim 10^{12}$ secs = 30,000 years

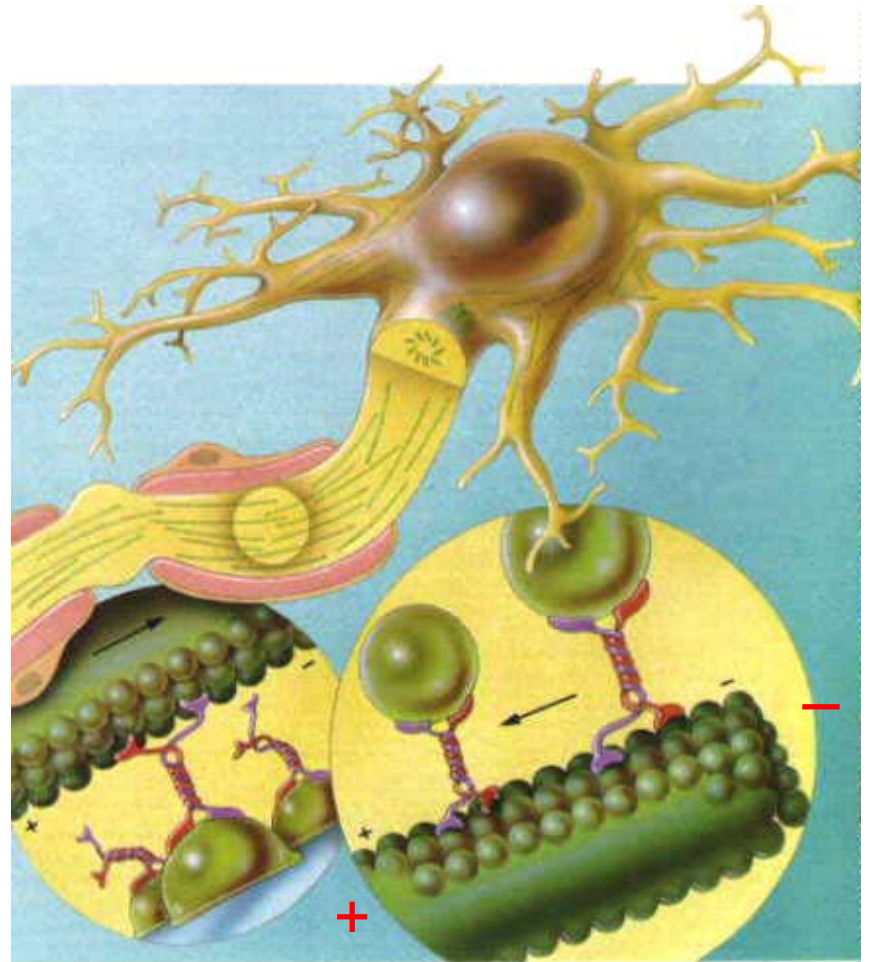
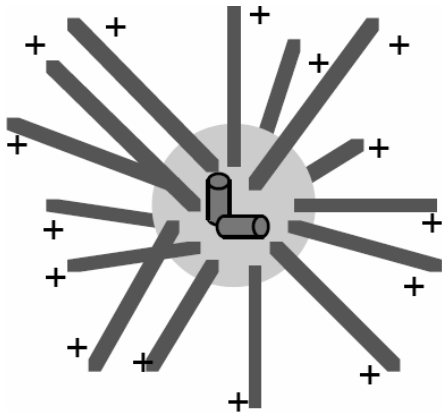
(hopeless for long transport along neurons)

At speeds up to $2\text{-}5 \text{ } \mu\text{m}/\text{s}$, a chemical cargo can be transported in 2-6 days from a production site in the brain to the end of a neuron a meter away.

Such cargo is transported on microtubules in a cell.

Motors on microtubules

- Hundreds of microtubules radiate from the **centrosome** of most animal cells.
- By probing the cell surface, the microtubules can push the nucleation region towards the center of the cell.
- Because they are relatively stiff, microtubules provide highways for transport in the cell.



dynein

+ → -

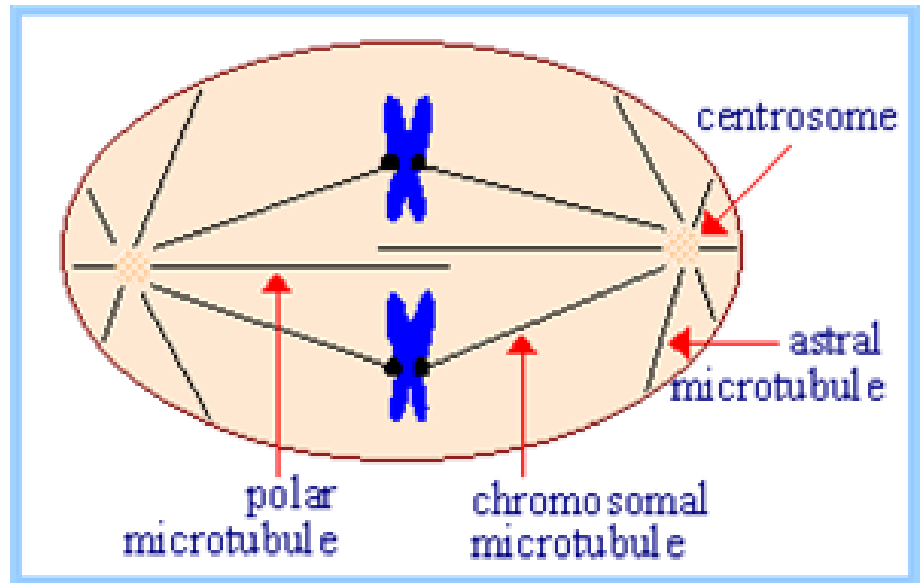
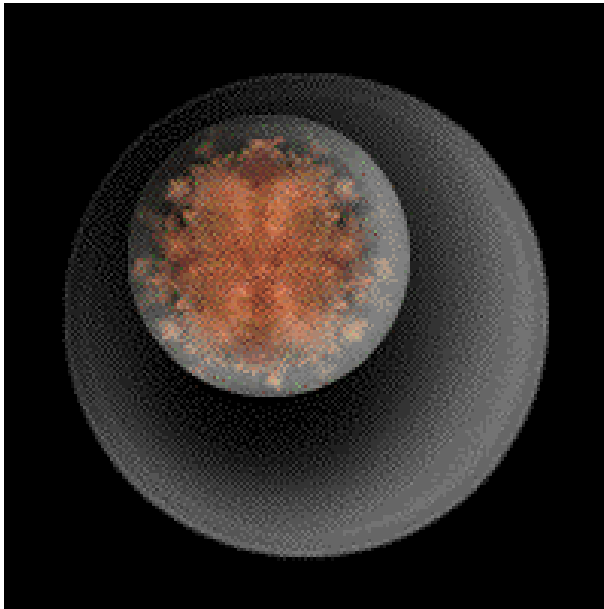
kinesin

- → +

Microtubules in mitosis

Different microtubules are involved in separating chromosomes during cell division, including

- (i) Polar microtubules,
- (ii) Chromosomal microtubules,
- (iii) astral microtubules.



Animation: http://www.virtualscience.com/gallery_animations.htm

1D filament formation/mechanics

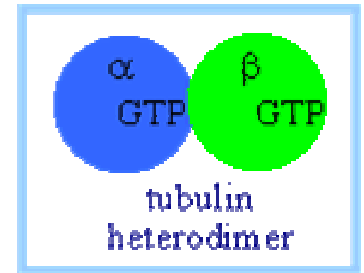
Structure of microtubules

Basic unit: α , β -tubulin (球形分子) heterodimer.

Each tubulin is a **G-actin** (globular actin).

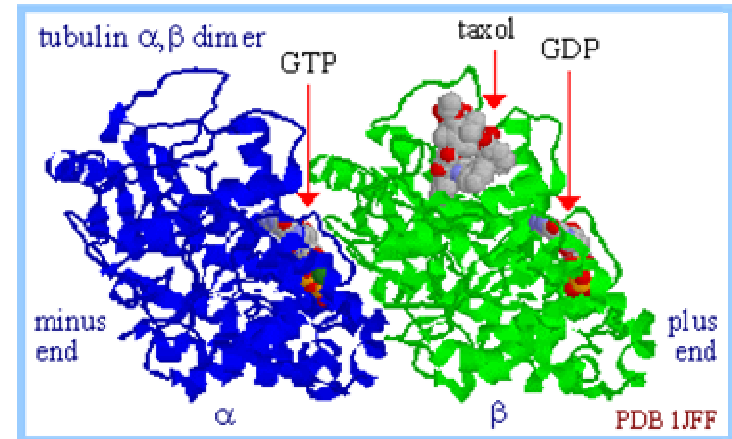
A chain of these tubulins is called a **F-actin** (fibrous actin).

The heterodimer does not come apart, once formed.



α -tubulin has a bound molecule of GTP, that **does not hydrolyze**.

β -tubulin may have bound GTP or GDP. Under certain conditions β -tubulin can hydrolyze its bound GTP to GDP plus Pi, release the Pi, and exchange the GDP for GTP.

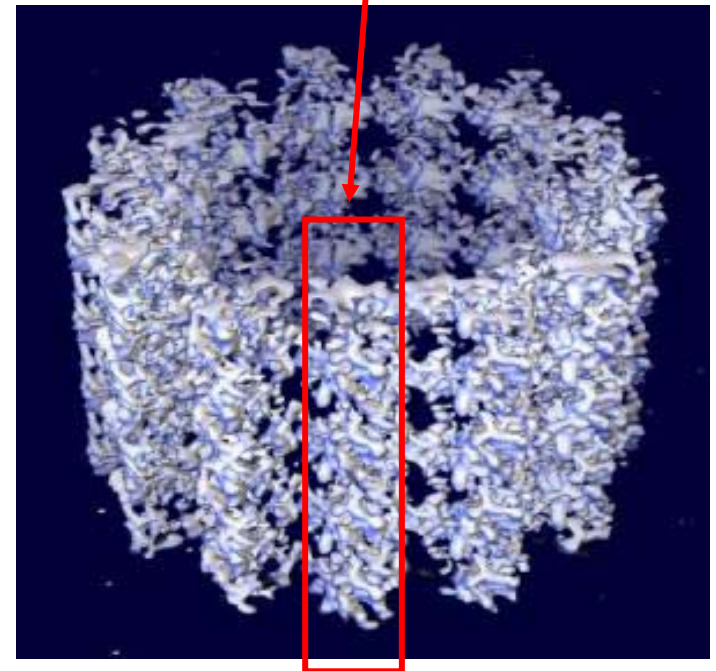
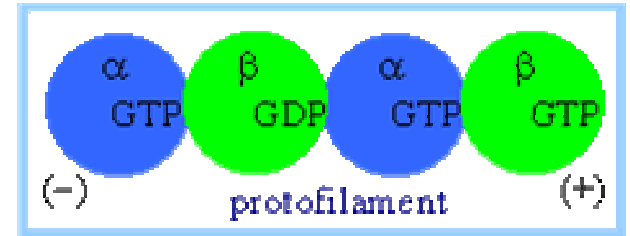


Microtubules

Protofilaments are formed by joining the α and β tubulins alternately.

A **microtubule** is a hollow cylinder about 25 nm in diameter, which is formed by putting 13 protofilaments to a helical cylinder wall.

GTP must be bound to both α and β subunits for a tubulin heterodimer to associate with other heterodimers to form a protofilament or microtubule.



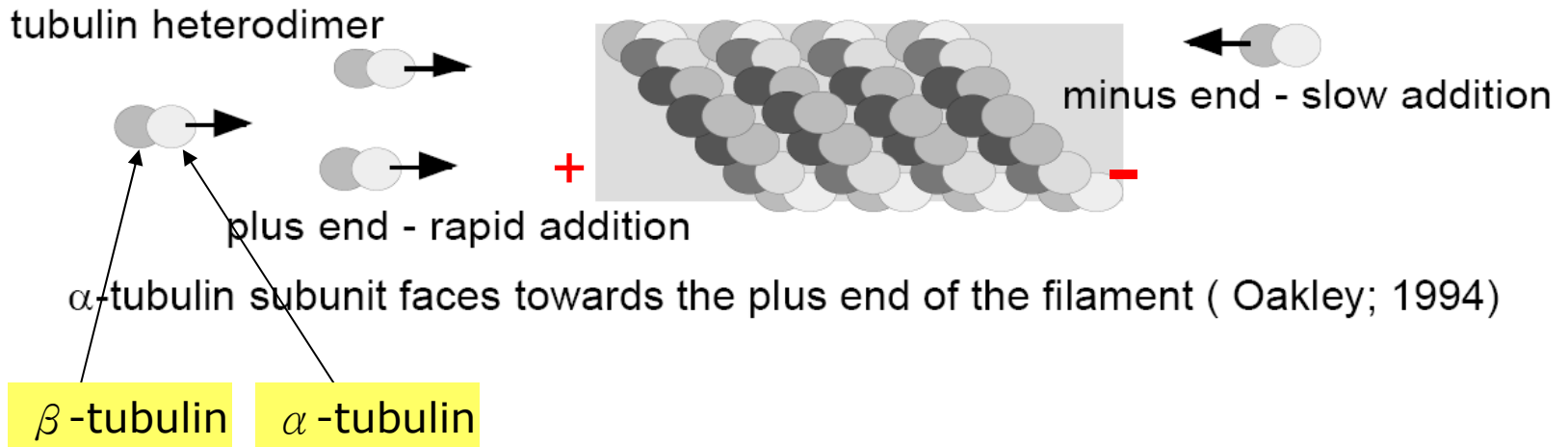
Polymerization of microtubules

Subunit addition brings β -tubulin that was exposed at the plus end into contact with α -tubulin.

This promotes hydrolysis of GTP bound to the now interior β -tubulin.

Pi dissociates, but β -tubulin within a microtubule cannot exchange its bound GDP for GTP.

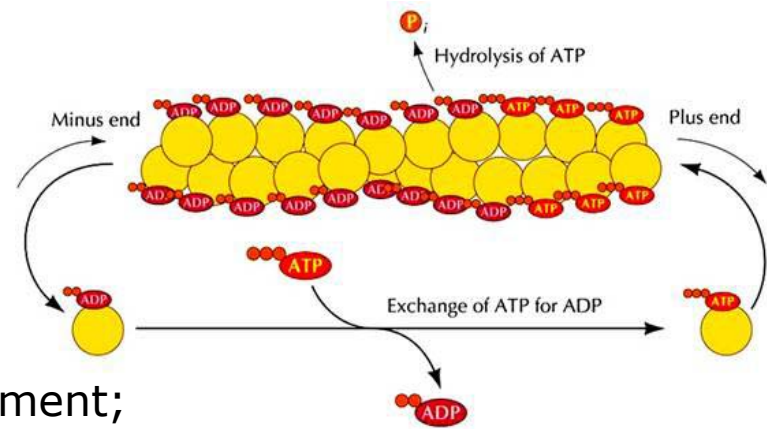
The GTP on α -tubulin does not hydrolyze.



Polymerization of actin and tubulin

(I) Equivalent filament ends:

n = number of monomers in a single filament;
 t = time;
 $[M]$ = concentration of free monomer in solution.



Capture rate of monomers by a single filament is proportional to the number of monomers available for capture

$$dn/dt = +k_{on} [M] \quad (\text{capture}) \quad (1)$$

k_{on} = capture rate constant, with units of $[\text{concentration} \cdot \text{time}]^{-1}$

Release rate **does not depend on $[M]$**

$$dn/dt = -k_{off} \quad (\text{release}) \quad (2)$$

k_{off} has units of $[\text{time}]^{-1}$

The net change of filament size is

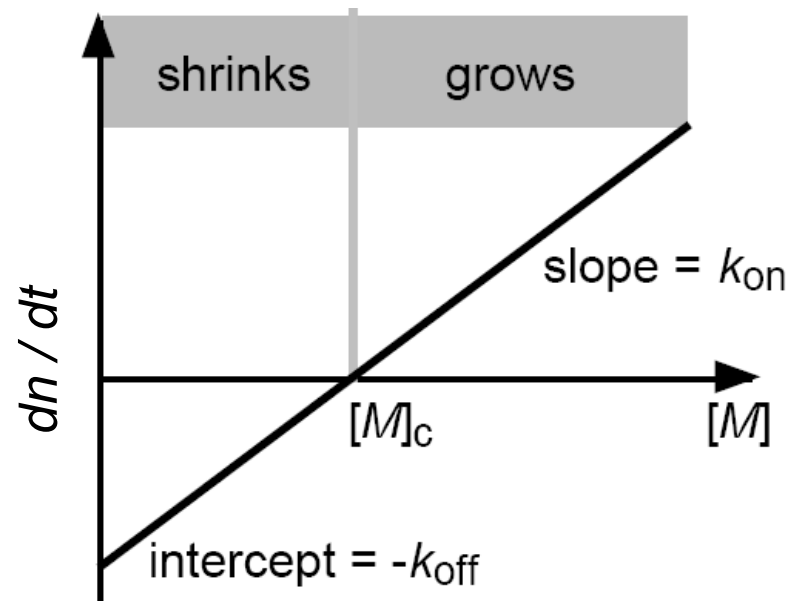
$$dn/dt = +k_{on} [M] - k_{off} \quad (3)$$

Obtain k_{on} and k_{off} from a plot of dn /dt against $[M]$

$dn /dt < 0 \rightarrow$ filament is shrinking

Minimum concentration for filament growth (often called the **critical concentration**) occurs at $dn /dt = 0$, where

$$[M]_c = k_{off} / k_{on}.$$



(II) Chemically inequivalent filament ends with different rate constants:

Treadmilling (Wegner, 1976)

Example: 2 x 2 inequivalent rate constants (+/- refer to the filament end):

$$dn^+/dt = k_{on}^+ [M] - k_{off}^+ \quad (5a)$$

$$dn^-/dt = k_{on}^- [M] - k_{off}^- \quad (5b)$$

Take $[M]$ to be the concentration of free triphosphate proteins.

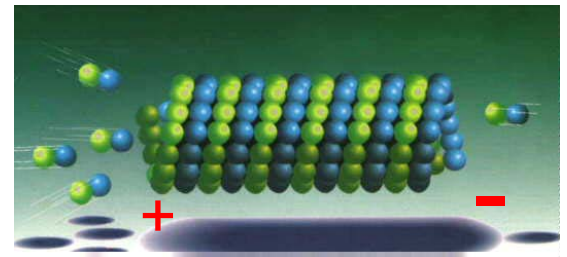
Probably different critical concentration at each end

$$\begin{aligned} [M]_c^+ &= k_{off}^+ / k_{on}^+ \\ [M]_c^- &= k_{off}^- / k_{on}^- \end{aligned} \quad (6)$$

Using data above for actin:

$$[M]_{ss} = 0.17 \mu M \quad \text{and} \quad dn^+/dt = 0.6$$

Direct measure of $[M]_{ss}$ under not dissimilar solution conditions yields $0.16 \mu M$ (Wegner, 1982).



The filament doesn't change length,

(I) if

$$[M]_c^+ = [M]_c^- \text{ (as in tubulin):}$$

both ends grow or both ends shrink simultaneously.

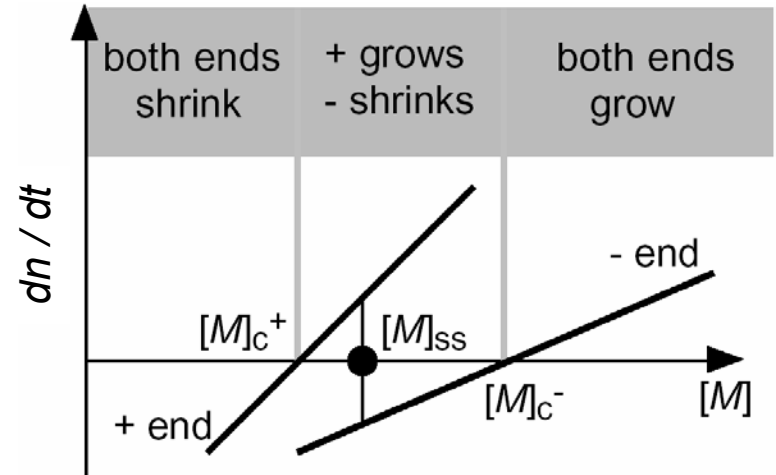
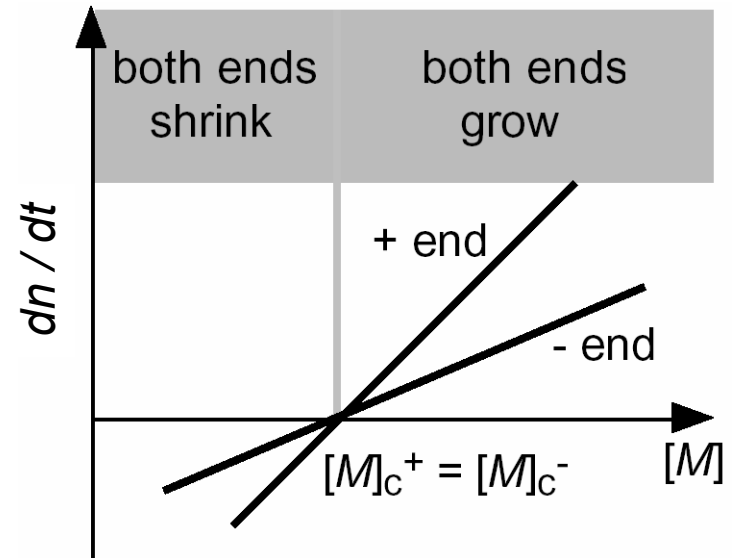
(II) if

$$dn^+ / dt = - dn^- / dt :$$

one end of the filament grows at the same rate as the other shrinks.

This occurs at steady state value $[M]_{ss}$

$$[M]_{ss} = (k_{off}^+ + k_{off}^-) / (k_{on}^+ + k_{on}^-)$$

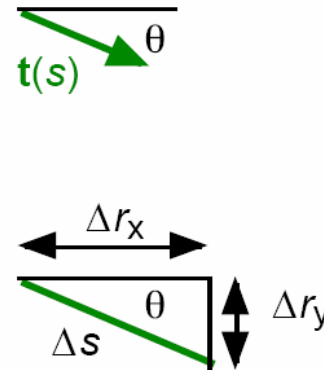
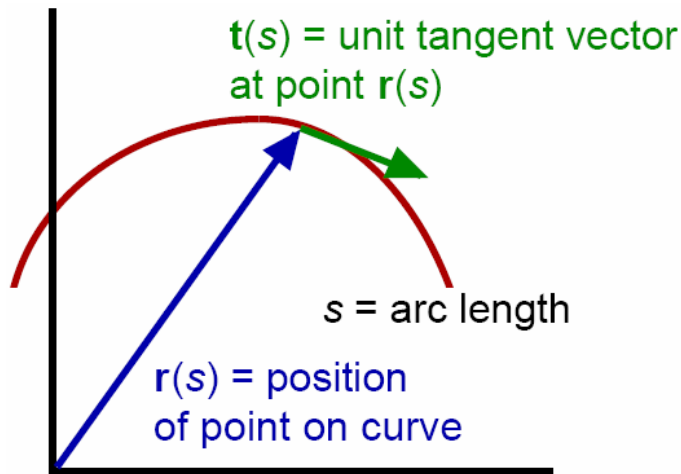


Bending energy of a thin polymer rod (Kratky-Porod model)

A straight rod of length L_c with uniform density and cross section, bent into an arc. The total energy becomes

$$E_{\text{bend}} = (k_f/2) \int_0^{L_c} (\partial \mathbf{t} / \partial s)^2 ds,$$

where k_f = **flexural rigidity**; units of [energy] · [length]



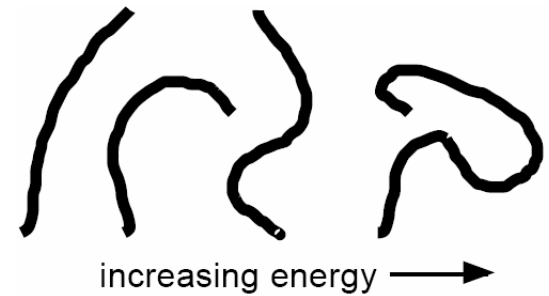
Thermal fluctuations and persistence length

At $T > 0$, shape of a filament can fluctuate:

The entropic force is the main driving in such a small filament.

Arc s of a circle with radius R_c : $\theta = s / R_c$
has an energy:

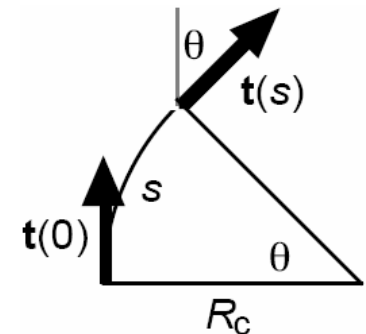
$$E_{\text{arc}} = k_f s / 2R_c^2 = k_f \theta^2 / 2s$$



Persistence length ξ_p decreases with increasing temperature

2 dim: $\langle E_{\text{arc}} \rangle = k_B T / 2$
 $\langle \theta^2 \rangle = s / \beta k_f$
 $\xi_p = 2 \beta k_f$

3 dim $\langle E_{\text{arc}} \rangle = k_B T$
 $\langle \theta^2 \rangle = 2s / \beta k_f$
 $\xi_p = \beta k_f$

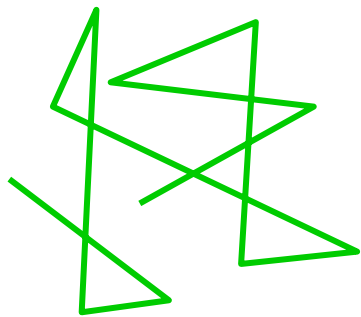


Sizes of flexible polymer chains

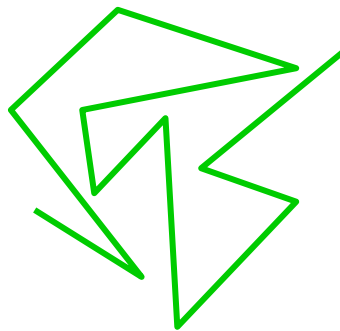
Radius of gyration of N particles is $R_g = \sqrt{\sum_{k=1}^N (r_k - r_m)^2 / N}$, where r_k is position of the k -th particle and r_m is the mean position of the particles.

Summary of scaling exponents for $\langle R_g^2 \rangle^{1/2} \sim N^\nu$

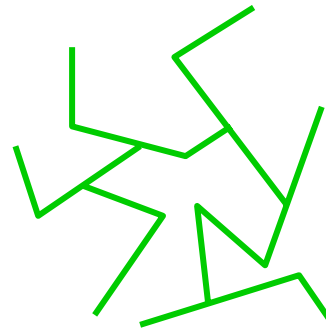
| <i>Configuration</i> | $d = 2$ | $d = 3$ | $d = 4$ |
|----------------------|---------|---------|---------|
| Ideal chains | 1/2 | 1/2 | 1/2 |
| Self-avoiding chains | 3/4 | 0.59 | 1/2 |
| Branched polymers | 0.64 | 1/2 | 1/2 |
| Collapsed chains | 1/2 | 1/3 | 1/4 |



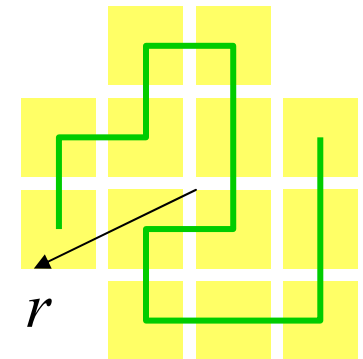
Ideal (random)



Self-avoiding



Branched

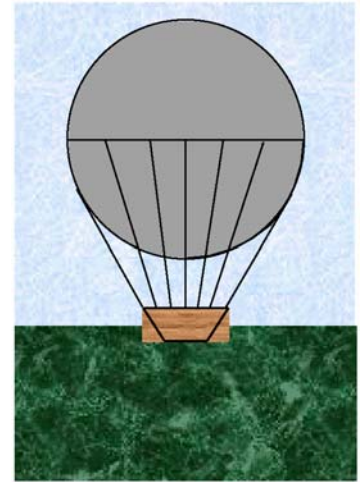


Collapsed

2D filament network

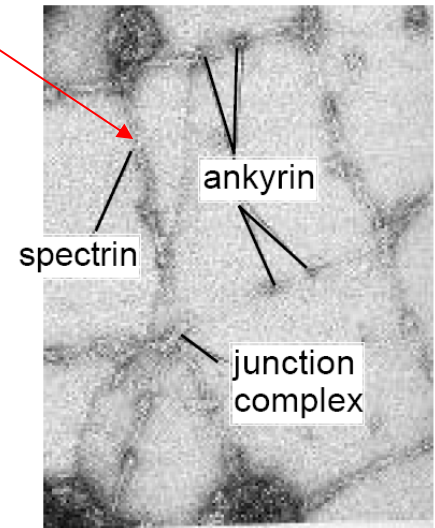
The explosion pressure on cells

(a) In a hot air balloon, a thin membrane confines the gas within the balloon, and an external **network** provides mechanical attachment points and may aid in maintaining the balloon's shape.



(b) A two-dimensional **network** of spectrin is attached to the red blood cell membrane to provide shear resistance (Liu, Derick and Palek, 1987).

The interior pressure of some cells, such as many varieties of bacteria, may be much higher than their surroundings. Thus, the engineering problem facing a bacterium is one of **explosion** like balloon rather than **collapse**.



(from Byers and Branton, 1985)

Deformation energy in 2D network

The strain tensor u_{ij} , related to the rate of change of \mathbf{u} with position \mathbf{x} by

$$u_{ij} = 1/2 [\partial u_i / \partial x_j + \partial u_j / \partial x_i + \sum_k (\partial u_k / \partial x_i) (\partial u_k / \partial x_j)],$$
$$\approx 1/2 [\partial u_i / \partial x_j + \partial u_j / \partial x_i]. \quad (\text{small deformations})$$

where i, j, k are Cartesian indices.

Just as the potential energy of a Hooke's law spring is quadratic in the square of the displacement, the change in the free energy density ΔF of a continuous object under deformation is quadratic in the strain tensor u_{ij} :

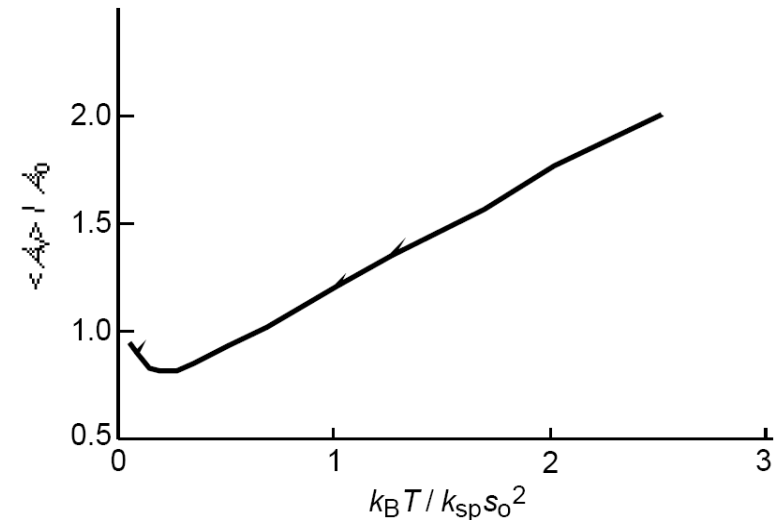
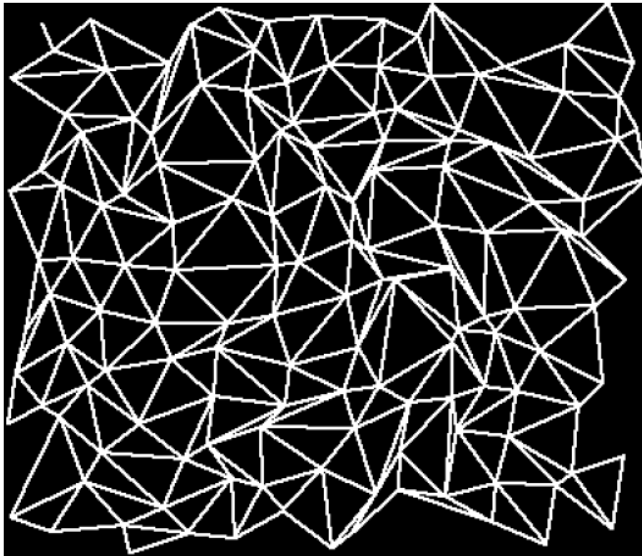
$$\Delta F = 1/2 \sum_{i,j,k,l} C_{ijkl} u_{ij} u_{kl}.$$

with the **elastic stiffness constants** or **elastic moduli** C_{ijkl} .

Network size at non-zero temperature

The network has a **negative** coefficient of thermal expansion at low temperatures (Lammert and Discher, 1998).

At higher temperatures, the area increases linearly with temperature with a slope close to 1/2.



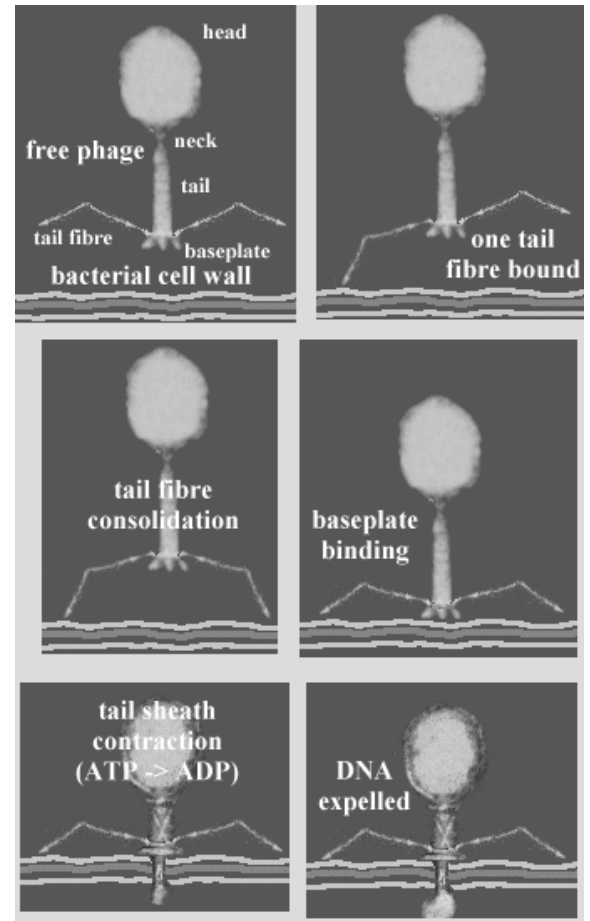
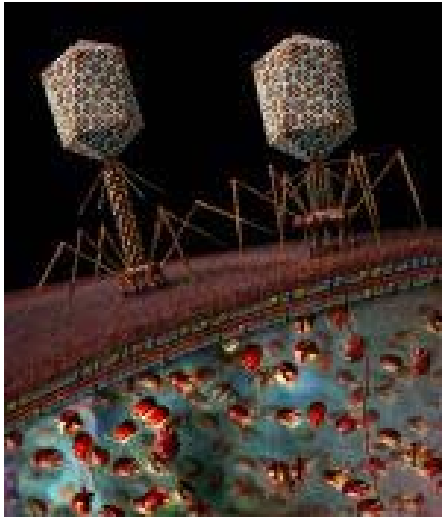
Snapshot of a triangular network of springs at $k_B T = k_{sp} s_0^2 / 4$.

4. Energy transduction (biological motors)



Biological motor I

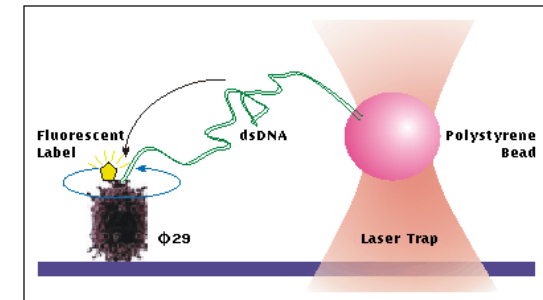
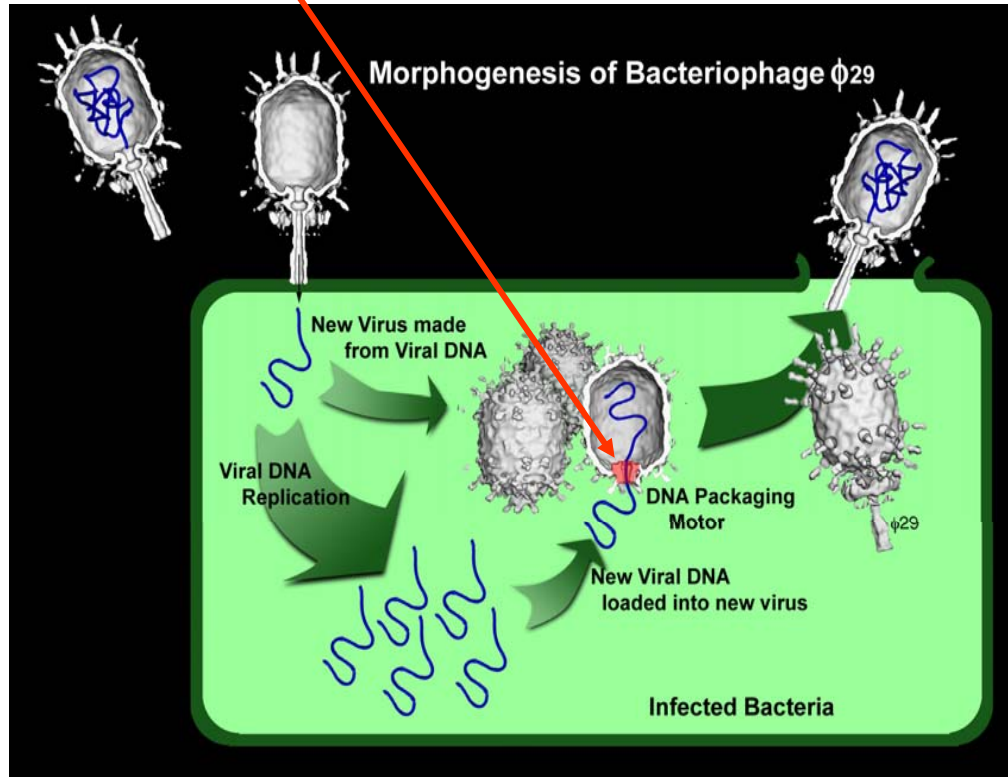
DNA packing motor in phage = bacteriophage (噬菌體)



Phage T4 - Enterobacteria phage T4, genus "T4-like Viruses", family *Myoviridae*, or viruses with 34-170 kbp dsDNA genomes, isometric heads and contractile tails - infects the gram-negative bacterium *E coli*.

<http://www.mcb.uct.ac.za/tutorial/virusentbacteria.htm>

Packing motor

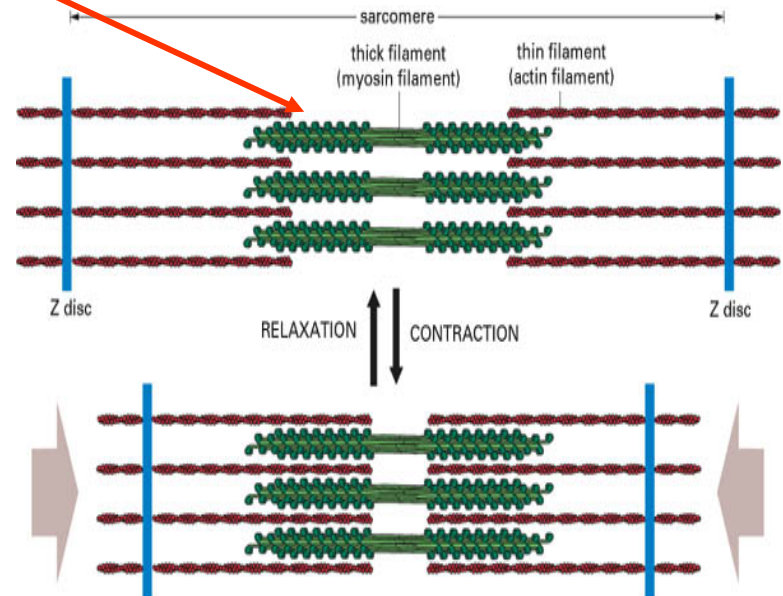
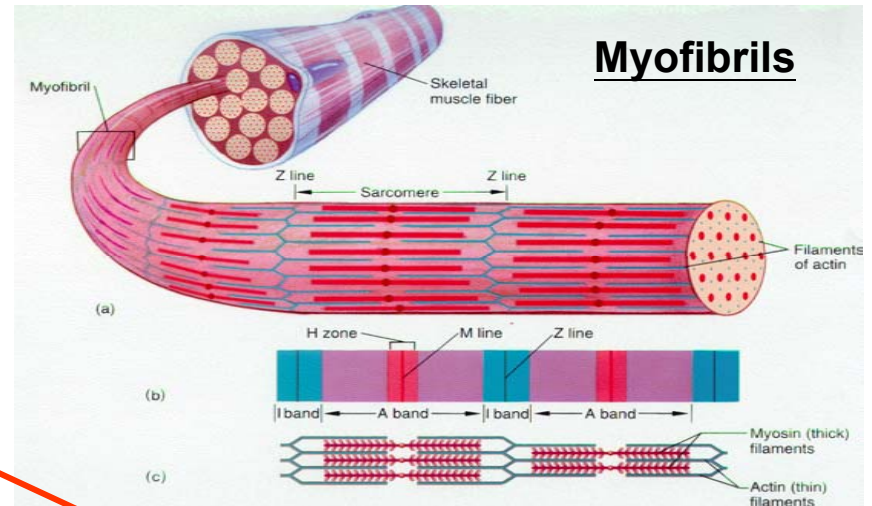
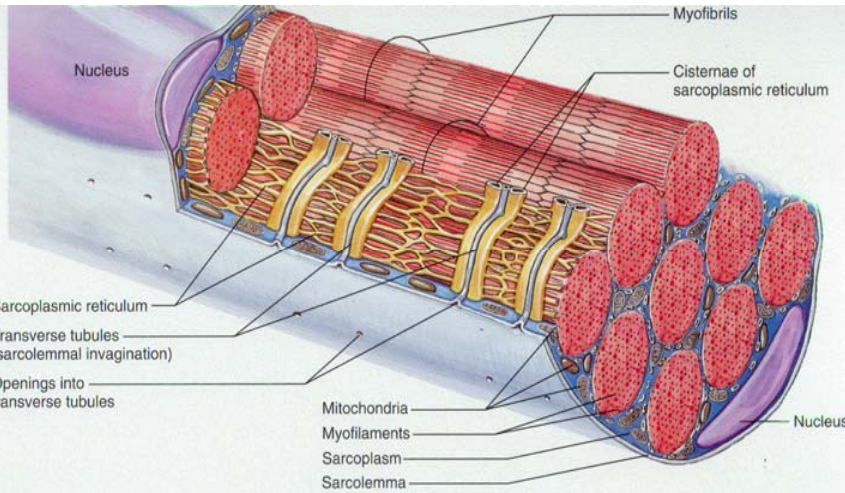


Bacteriophage phi29 (DNA packaging motor 60 pN) packages a DNA about 130 times longer than the viral shell in just 3 minutes (60 atmosphere inside the viral shell).

Biological motor II

Myosin: Muscle contraction (5pN)

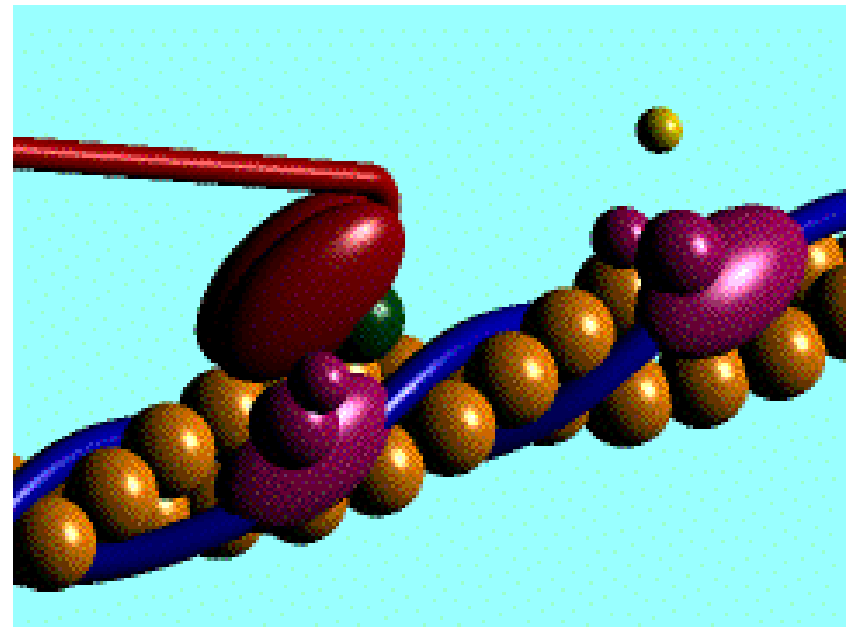
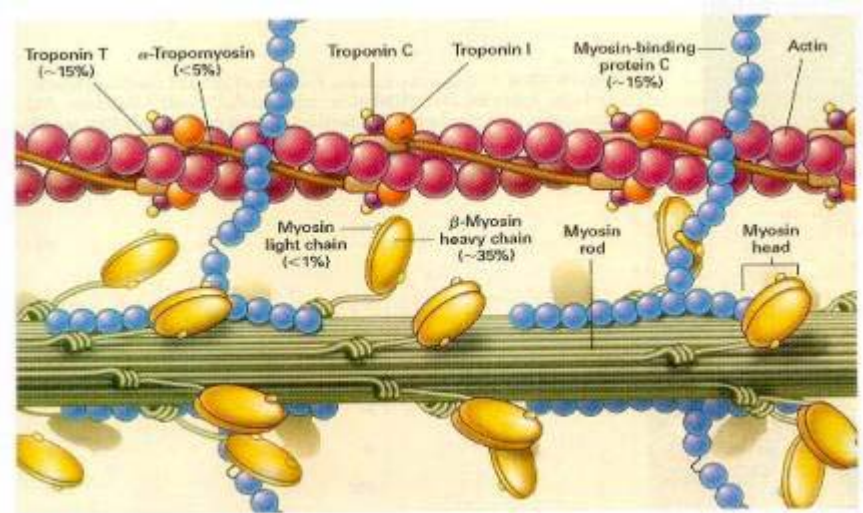
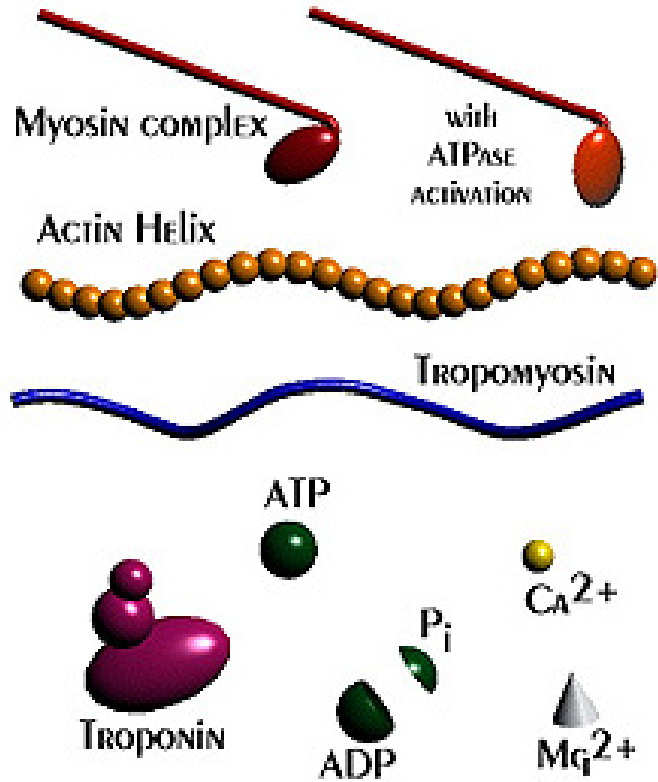
Cytology of a skeletal muscle



<http://www.gdn.edu/Faculty/pjen/Anatomy%20and%20physiology%202211K%20-%20lecture%201%20notes.htm>

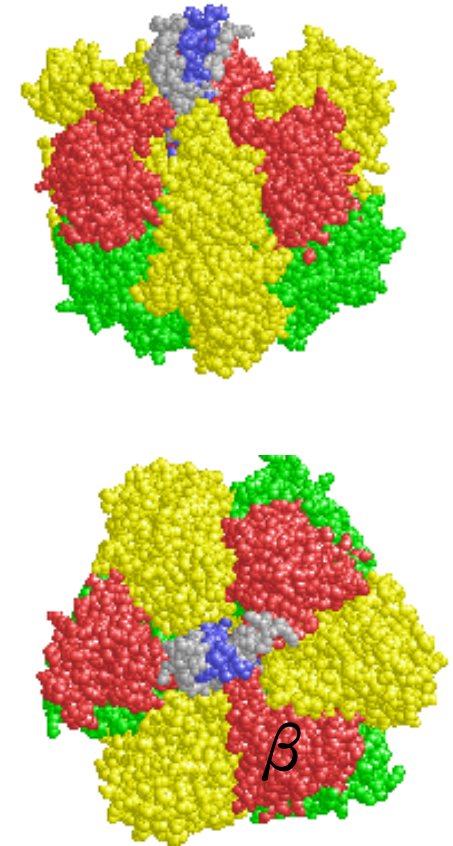
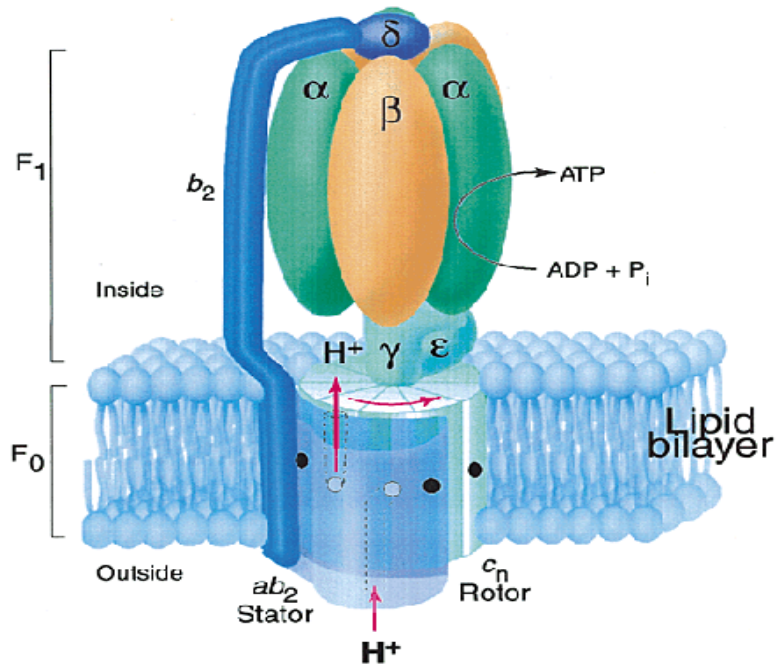
Animation: <http://entochem.tamu.edu/MuscleStrucContractswf/index.html>

Myosin head



Biological motor III

F₀F₁ ATPase: a machine to produce ATP



generates energy ATP (synthesis) $\text{ADP} + \text{P}_i \rightarrow \text{ATP}$
consumes energy ATP (hydrolysis) $\text{ATP} \rightarrow \text{ADP} + \text{P}_i$

Animation: <http://plantcell.lu.se/ltm/06/3ATP.html>
<http://vcell.ndsu.nodak.edu/animations/atpgradient/movie.htm>

Rotation Experiments

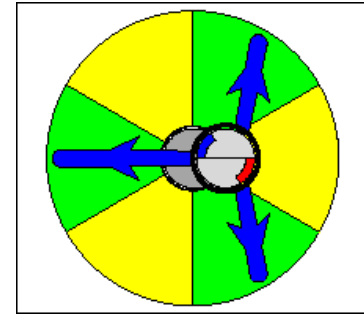
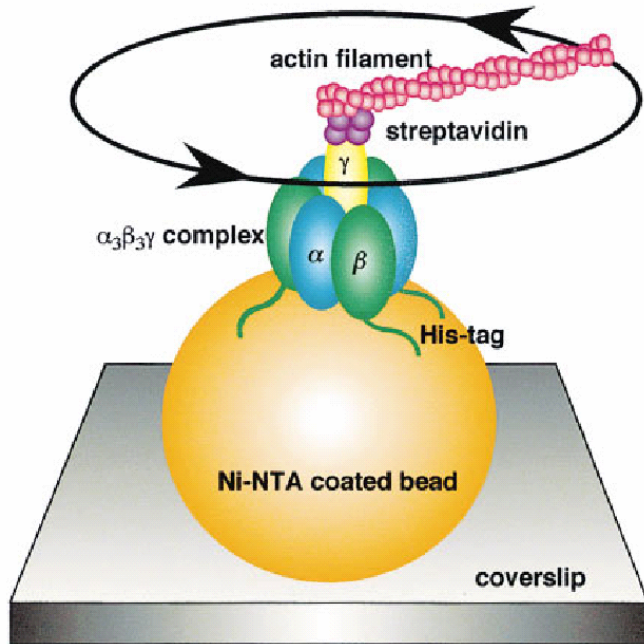
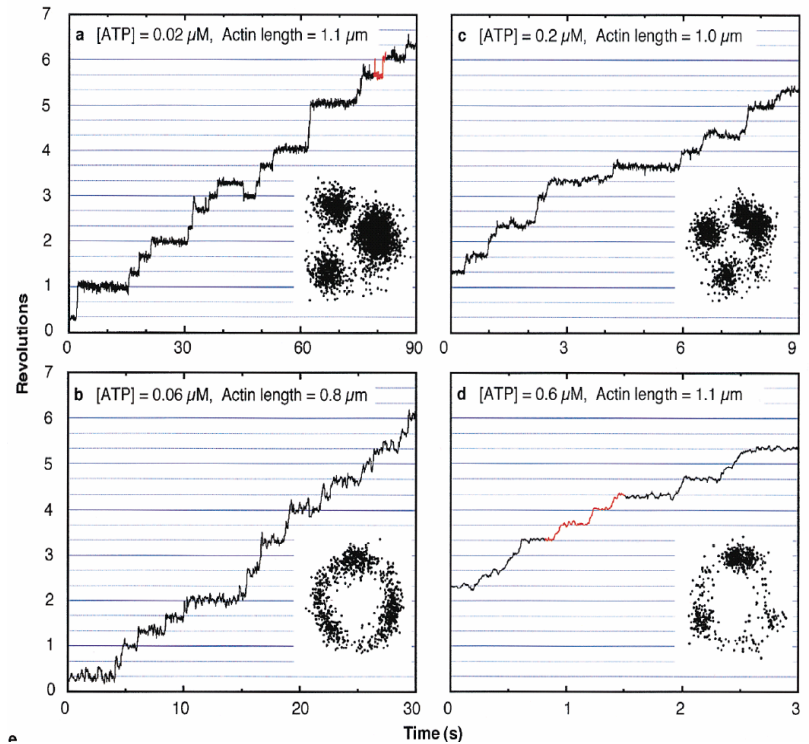


Figure 1. Experimental System (Not to Scale)

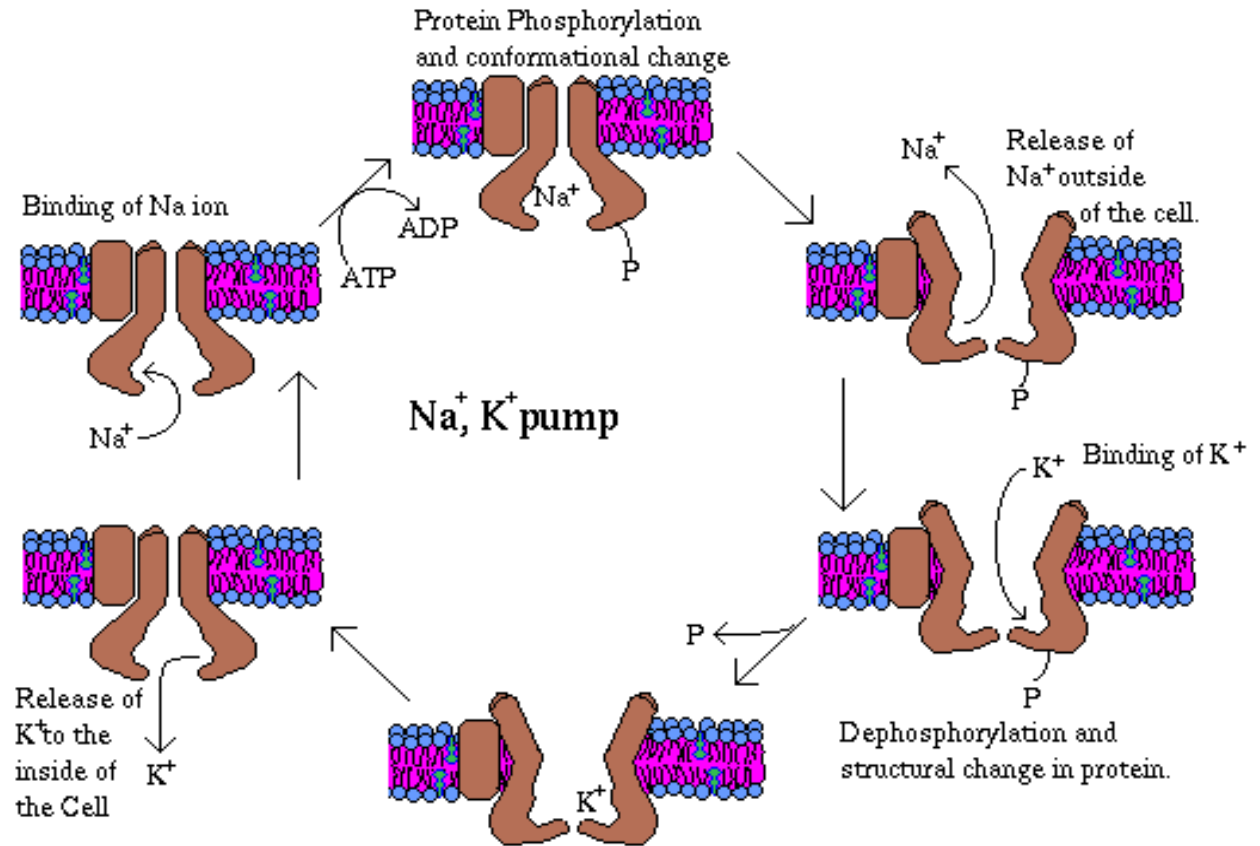
consume ATP



Biological motor IV

Ion pump

(Na,K-ATPase,
Ca-ATPase)



- The Top is the Outer membrane.
- The Bottom is the inner membrane (inside of the Cell)

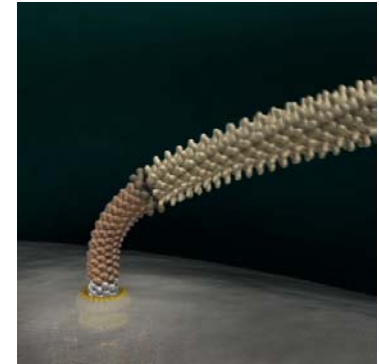
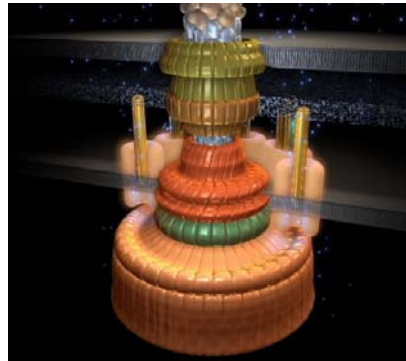
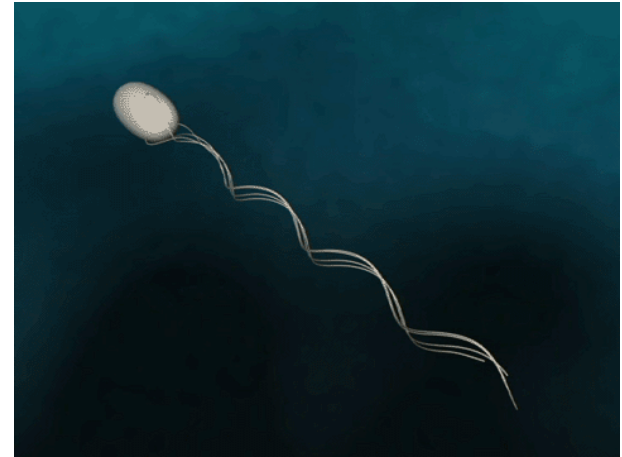
Animation:

<http://www.cm.utexas.edu/academic/courses/Spring1999/CH339K/Caras/Web339K/Animations/IonPump.html>

Biological motor V

Flagellar motor in *E. coli*

<http://www.npn.jst.go.jp/>



Biological motors

| Category | Motion | Function | Force (pico N) | Weight (kDa) |
|-----------------|-----------------|--------------------------------------------|----------------|--------------|
| Myosin | Muscle Filament | Muscle contraction Cell division | 5-6 | 460-500 |
| Kinesin | Microtubule | Chromosom deviation Cell unit transport | ~ 6 | 400 |
| Dynein | Microtubule | Fragellan | Unclear | 1000 |
| RNA Polymerase | DNA | Synthesis | 25-30 | 490 |
| DNA Polymerase | DNA | Synthesis | 34 | 94 |
| Helicase | DNA | DNA rewinding | Unclear | 100 |
| ATP synthase | None | ATP synthesis & hydrolysis | 80 | 500000 |
| Flagellar motor | None | Swimming | 200 | Rather large |

Transcription, Translation

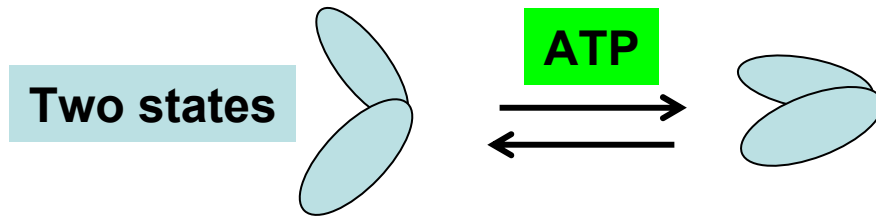
1Da = 1 u \approx 1.66×10^{-27} kg,

C atom has 12 Da.

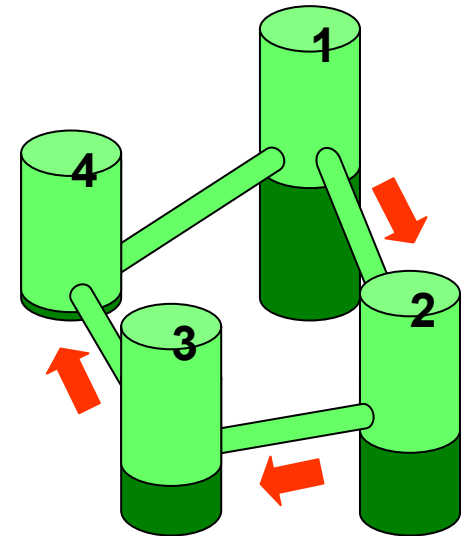
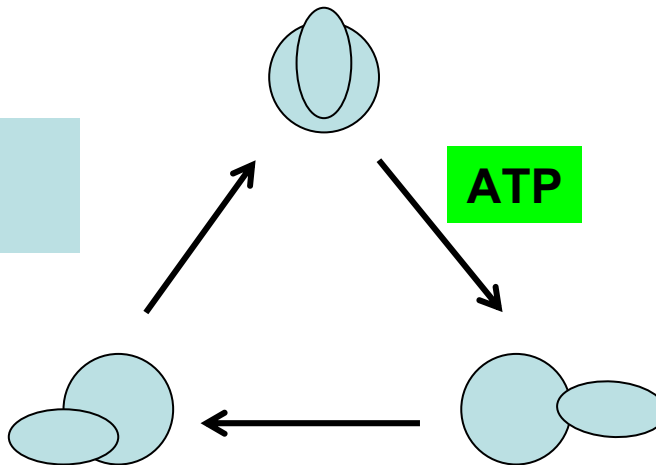
<http://vcell.ndsu.nodak.edu/animations>

Working principle of biological motors:

cyclic conformational change

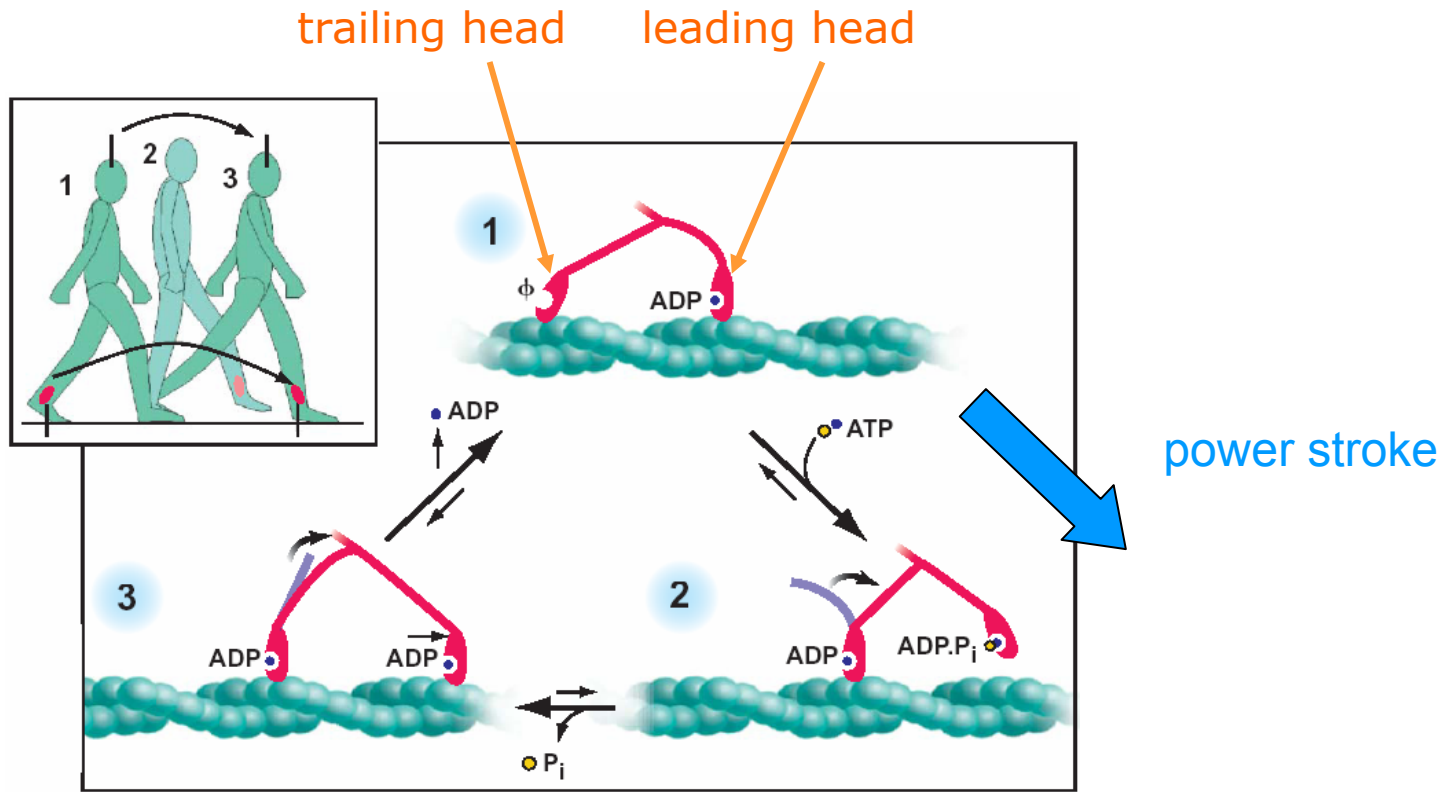


Multi-states



Free energy $1 > 2 > 3 > 4$

Example: Myosin three-state model



ATP = ADP (●) + inorganic phosphate (○).

The molecule dwells in state 1 at low ATP concentrations (because the empty site, ϕ , requires ATP to bind before it can be released from actin)

Force and speed of motors on a filament

Stokes' Law

For spherical objects of radius R at low speeds moving through a fluid of viscosity η , the drag force it encounters is

$$F_{\text{drag}} = c_1 v \quad \text{with } c_1 = 6 \pi \eta R.$$

If an object has an initial speed v_0 , it will come to rest in a distance $x = mv_0 / c_1$ with a linear dependence of F_{drag} on v .

Example:

A vesicle of radius 50 nm is carried along a filament at speed $0.5 \mu\text{m/s}$ in a cell with viscosity $10^{-1} \text{ kg / m} \cdot \text{s}$ which is 100 times more viscous than water. Thus

$$c_1 = 6 \pi \eta R = 6 \pi \cdot 10^{-1} \cdot 5 \times 10^{-8} = 30 \pi \cdot 10^{-9} = 9.4 \times 10^{-8} \text{ kg/s.}$$

$$F_{\text{drag}} = c_1 v = 9.4 \times 10^{-8} \cdot 5 \times 10^{-7} = 5 \times 10^{-14} \text{ N} = 0.05 \text{ pN.}$$

A typical molecular motor can generate 2-4 pN of force, easily enough to drive the vesicle.

Typical speed and viscosity

Observed speed of movement

| <i>motion</i> | <i>typical speed</i> ($\mu\text{m/s}$) | <i>example</i> |
|---------------------------|------------------------------------------|--------------------------------------------------|
| actin filament growth | $10^{-2} - 1$ | 0.3 $\mu\text{m/s}$ at $[M] = 10 \mu\text{M}$ |
| actin-based cell crawling | $10^{-2} - 1$ | fibroblasts move at $\sim 10^{-2} \mu\text{m/s}$ |
| myosin on actin | $10^{-2} - 1$ | 0.1 - 0.5 $\mu\text{m/s}$ common in muscles |
| microtubule growth | up to 0.3 | 0.03 $\mu\text{m/s}$ at $[M] = 10 \mu\text{M}$ |
| microtubule shrinkage | 0.4 - 0.6 | 0.5 $\mu\text{m/s}$ |
| fast axonal transport | 1-4 | |
| slow axonal transport | $10^{-3} - 10^{-1}$ | |

$[M]$ = monomer concentration. Fast axonal transport involves kinesin or dynein moving along microtubules.

Viscosity of different fluids

| Fluid | η (kg/m \cdot sec at 20 $^{\circ}\text{C}$) |
|-----------|-----------------------------------------------------|
| Water | 1.0×10^{-3} |
| Olive oil | 0.084 |
| Glycerine | 1.34 |
| Glucose | 10^{13} |

Force scale at the molecular level

Langevin force is the force bacteria encountered under Brownian motion, which is about

$$10^{-14} \text{ N.}$$

Biological motor perform a force of the order:

$$10^{-11} \text{ N.}$$

Cohesion force for (i) hydrophobic interactions and (ii) hydrogen bonding contributing to the stability of biomolecules is of order

$$10^{-10} \text{ N.}$$

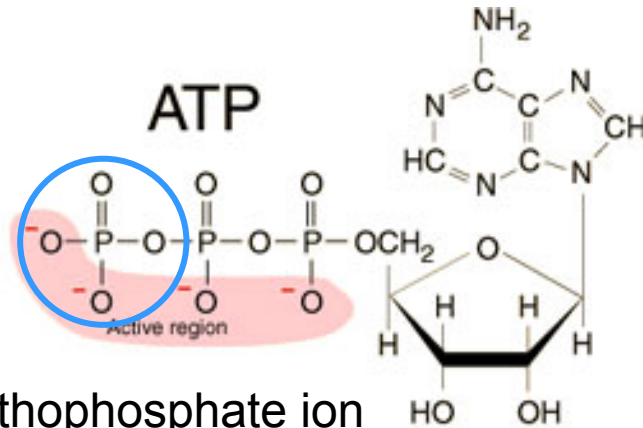
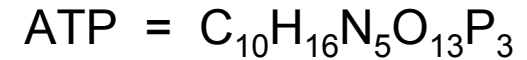
Covalent bond is the strongest force at the molecular level, which can denature a protein and is on the order of

$$10^{-9} \text{ N.}$$

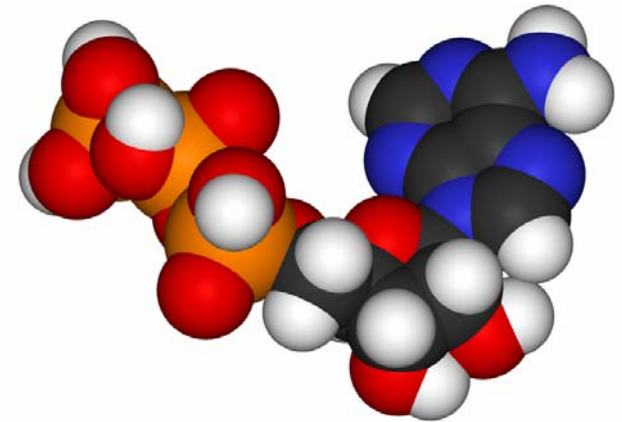
Energy currency ATP

Energy source of the motors

ATP (adenosine triphosphate) 三磷酸腺苷



The orthophosphate ion HPO_4^{2-} written as P_i is the inorganic phosphate.



A ATP contains energy 5×10^{-20} J, corresponding to -31 kJ/mol.

A typical biological motor consumes 10^2 - 10^3 ATP/sec ($\approx 10^{-16}$ - 10^{-17} W).

Each cell contains 10^9 ATP.

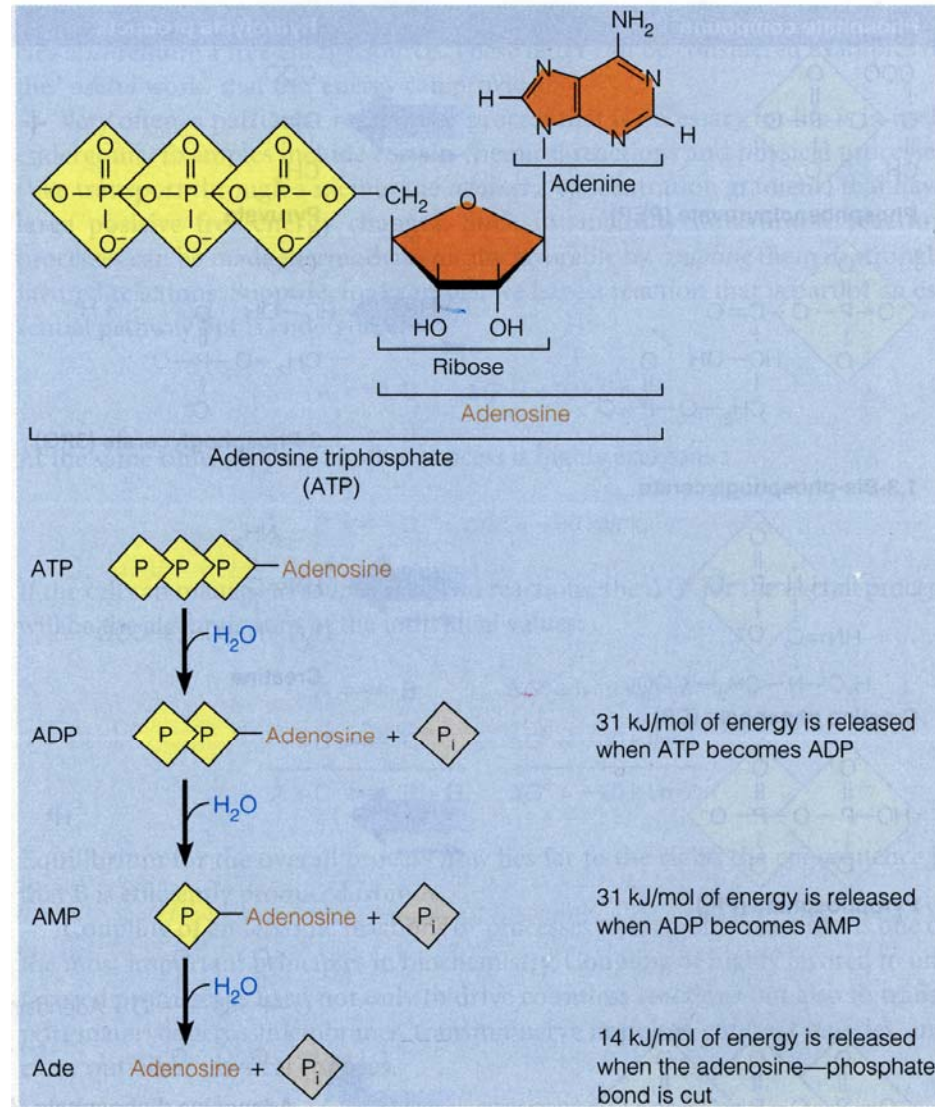
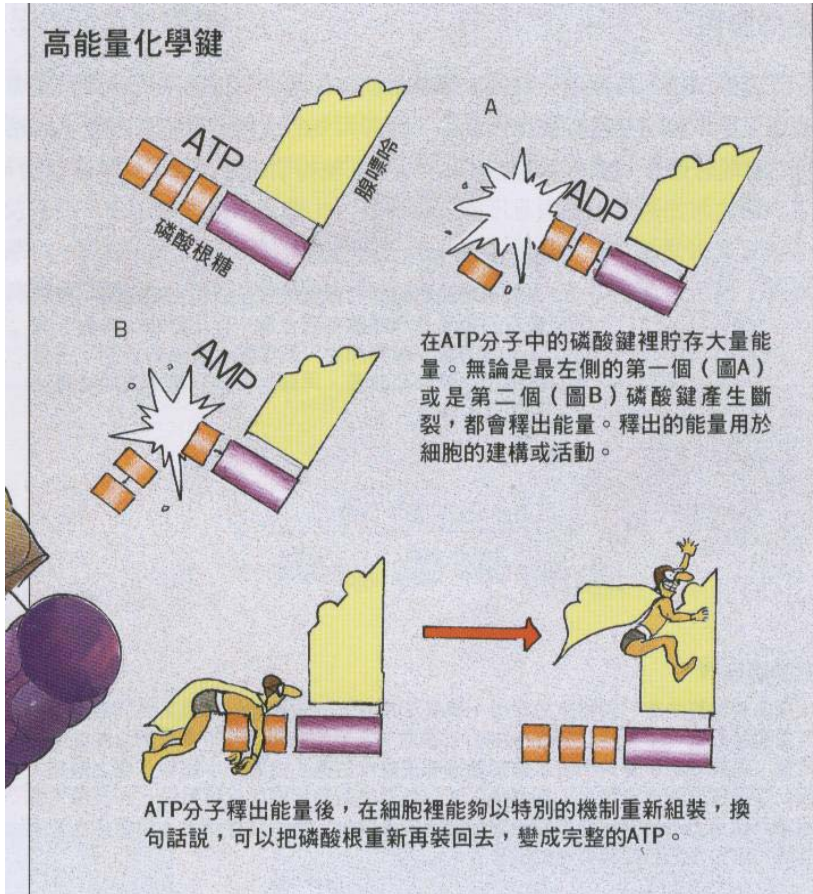
The total human body contains about 50 grams ATP.

Our average daily ATP consumption is 180 kg.

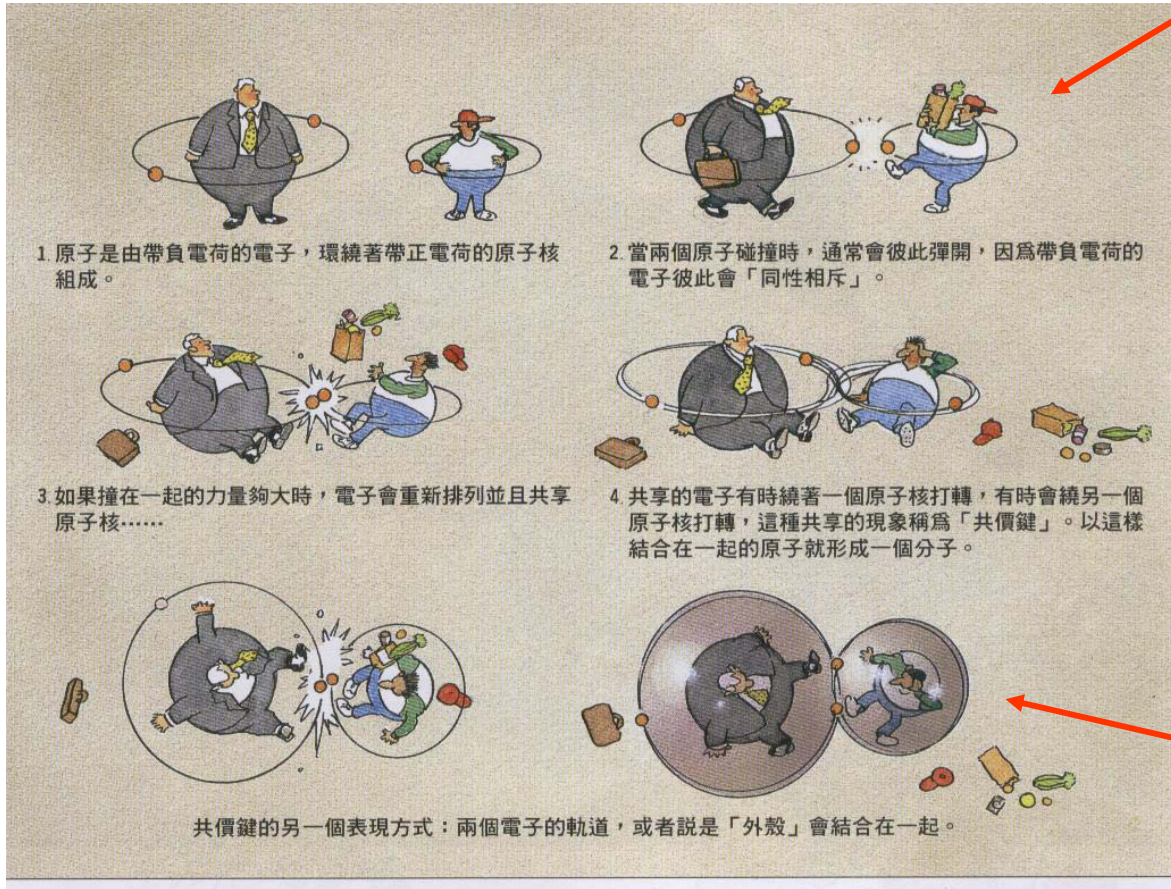
(Kornberg, Arthur. 1989. *For the love of enzymes*. Harvard Uni. Press. P.65)

Energy released by hydrolysis of ATP

高能量化學鍵



Physical picture of a high energy bond



Steric repulsion (exchange repulsion, hard core repulsion) comes from the overlap of electron clouds of atoms. This repulsion is a quantum mechanical effect and empirically behaves like

$$U \propto \frac{1}{r^n}$$

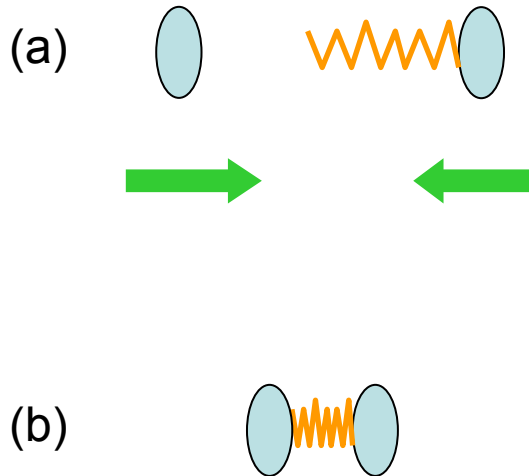
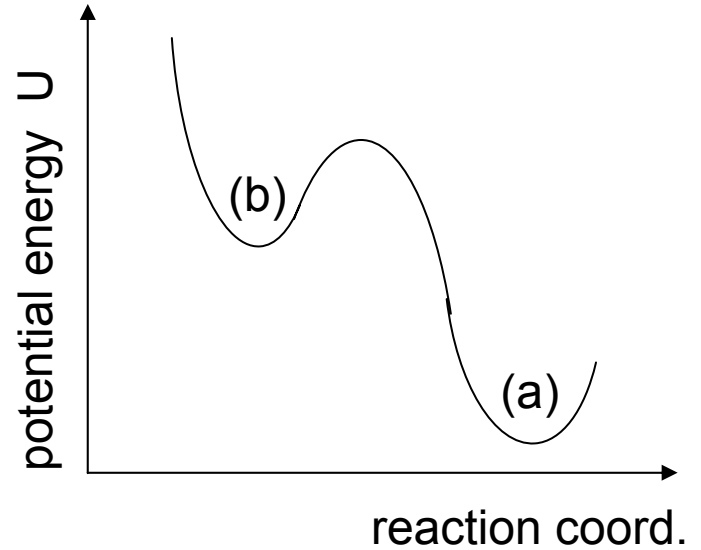
where n is an integer between 9 and 16.

After the collision the kinetic energy is largely reduced.

A bond has high energy because high kinetic energy is converted and stored in potential energy.

Energy exchange

A proper collision may convert kinetic energy to potential energy (a) \rightarrow (b).



有些時候，用力的撞擊會使原子結合成分子。

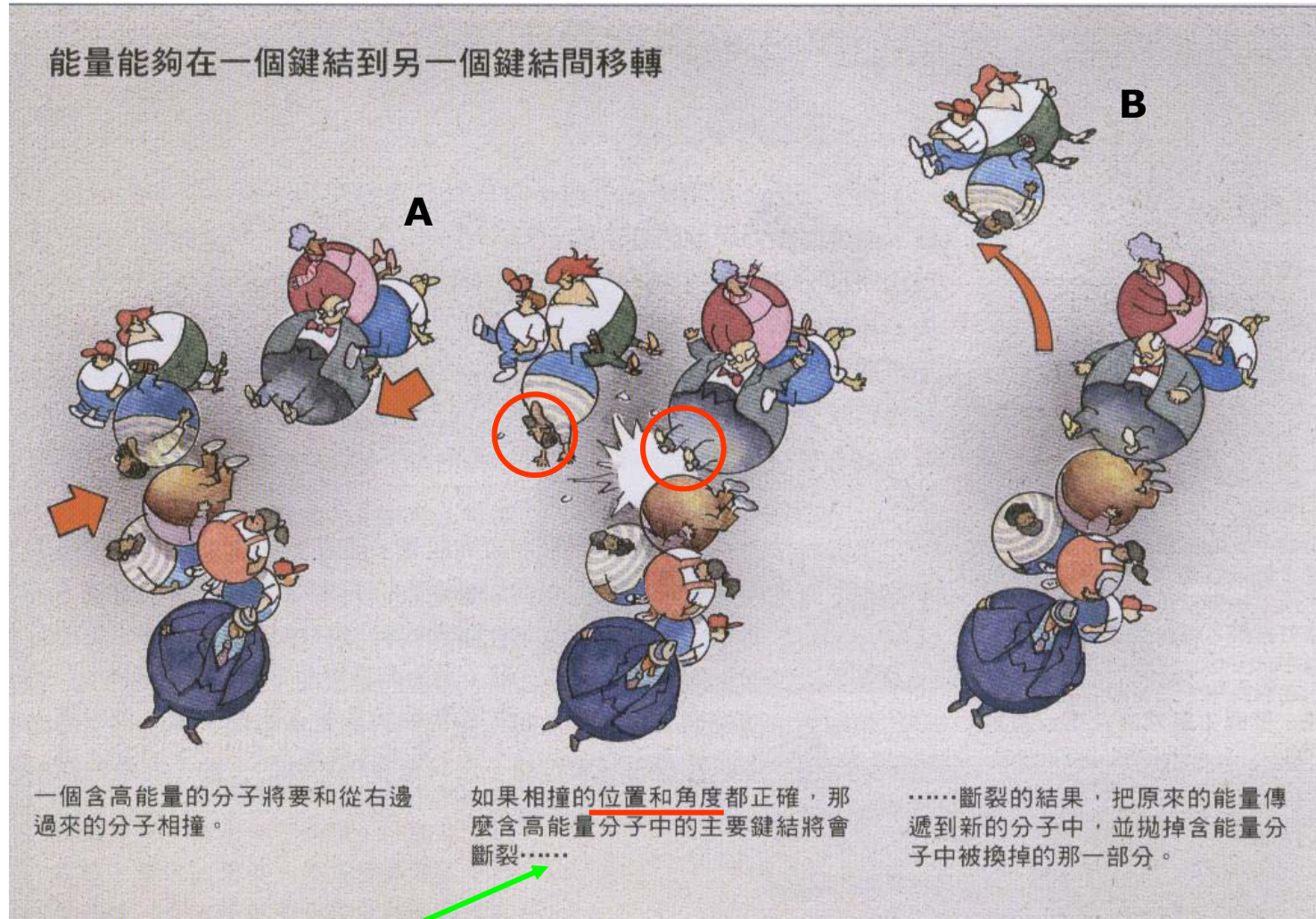


理論上，原子間成功的一連串撞擊應該能形成一串長鏈分子。



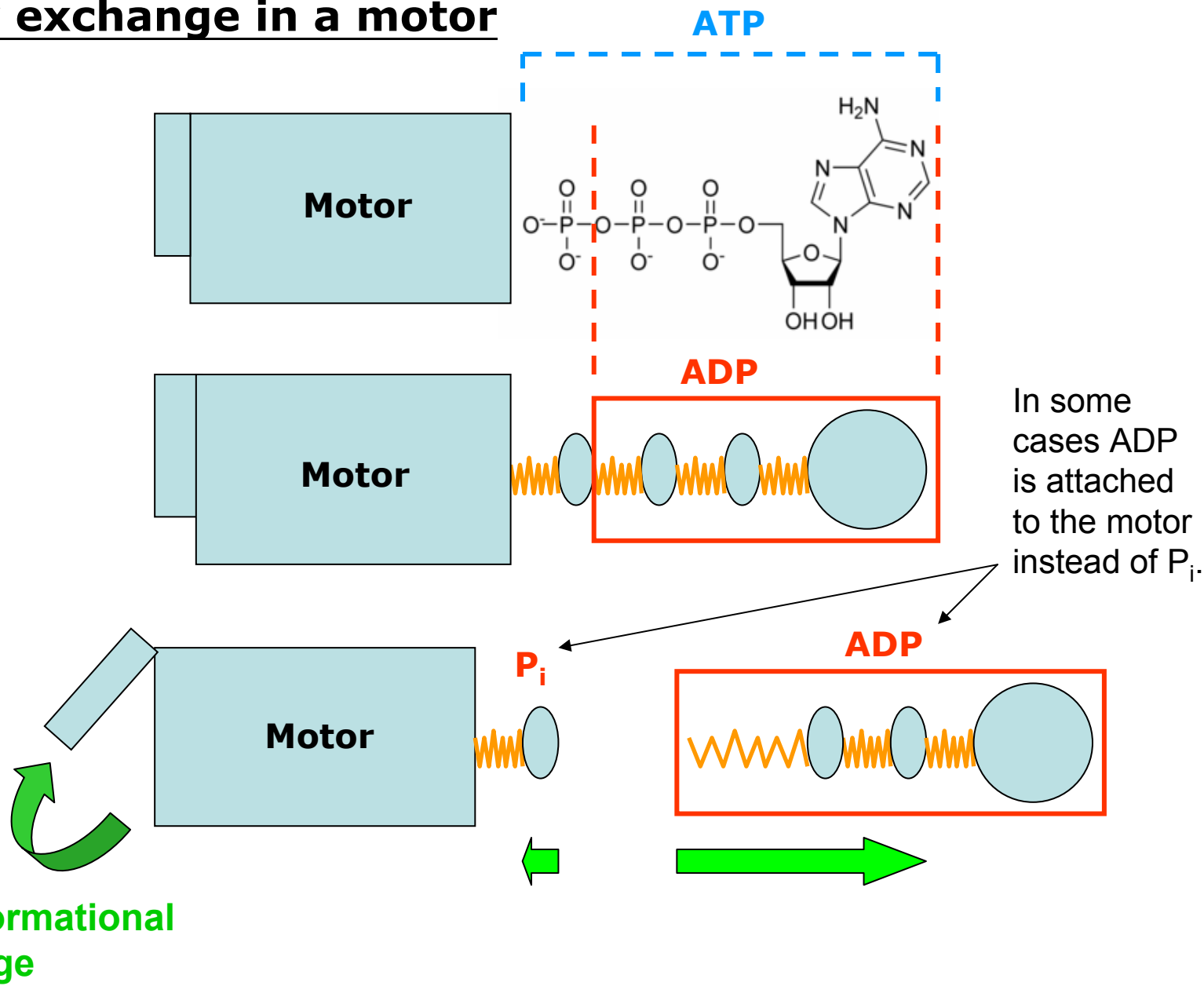
A proper collision also may release energy (b) \rightarrow (a)

Energy exchange between molecules



Even the kinetic energy of A is lower than B.

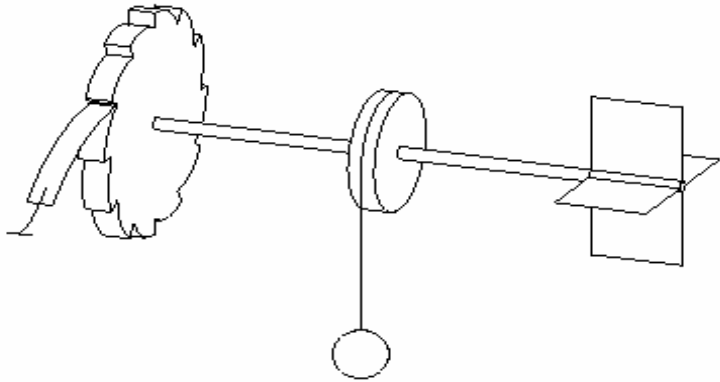
Energy exchange in a motor



Ratchet (directed motion)

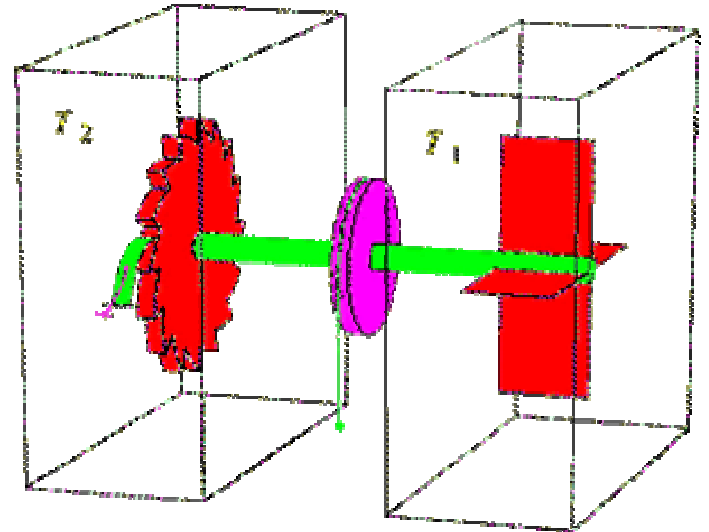
Ratchet effect

M.v.Smoluchowski
Physik. Zeitschrift, 13, 1069
(1912)



Thought experiment of a perpetuum mobile against the Second Law of Thermodynamics

R.P. Feynman
The Feynman Lectures on
Physics, Vol. 1, Chap.46,
(1963)



The Second Law of Thermodynamics is **not violated**.

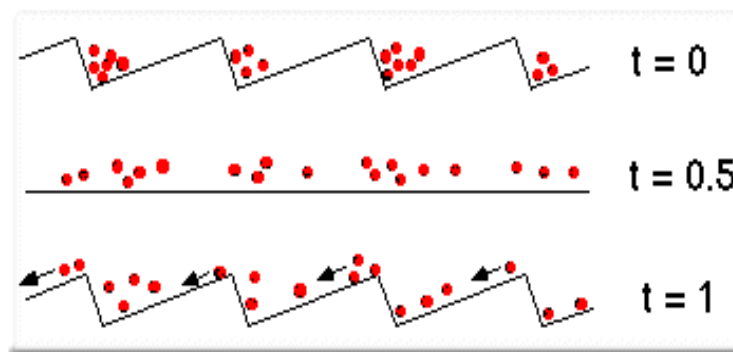
Non-equilibrium + proper asymmetry = directed motion

A simple mathematical Ratchet Model

- Ratchet ingredients:**
- Brownian particle (mass m)
 - periodic asymmetric potential $V(x,t)$
 - driving forces $f(t)$ of zero mean, i.e., $\langle f(t) \rangle = 0$

Ratchet model:
$$m\ddot{x} + \gamma\dot{x} + \frac{dV(x,t)}{dx} = f(t)$$
 damping coefficient γ

Interesting behavior: $\langle \dot{x} \rangle \neq 0$ even when $\langle f(t) \rangle = 0$



Animation: <http://monet.physik.unibas.ch/~elmer/bm/>

General classification of Ratchet Models

(1) A Langevin-based approach (Horsthemke 1994)

$$\xi \frac{dx}{dt} = -\partial_x W(x) + F(t) \quad \langle F(t) \rangle = 0$$

(2) An approach with fluctuating potential (Astumian & Bier 1994)

$$\xi \frac{dx}{dt} = -\partial_x W(x,t) + f(t) \quad \langle f(t) \rangle = 0, \quad \langle f(t)f(t') \rangle = 2\xi T \delta(t-t')$$

(3) A mode with several internal states of the particle described by the Langevin equation that depends on the state $i = 1, 2, \dots, N$

$$\xi_i \frac{dx}{dt} = -\partial_x W_i(x) + f_i(t) \quad \langle f_i(t) \rangle = 0, \quad \langle f_i(t)f_j(t') \rangle = 2\xi_i T \delta(t-t') \delta_{ij}$$

where W is the potential the particle experiences;

$F(t)$ is a fluctuating force of zero average (not white noise);

$f(t)$ is a Gaussian white noise.

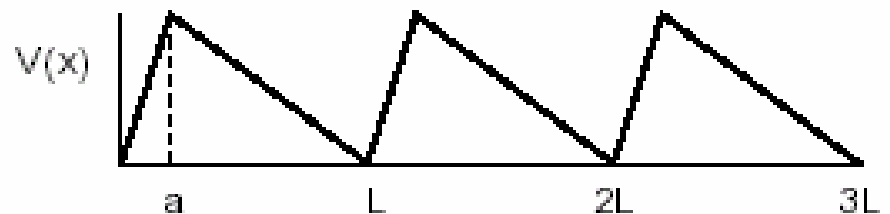
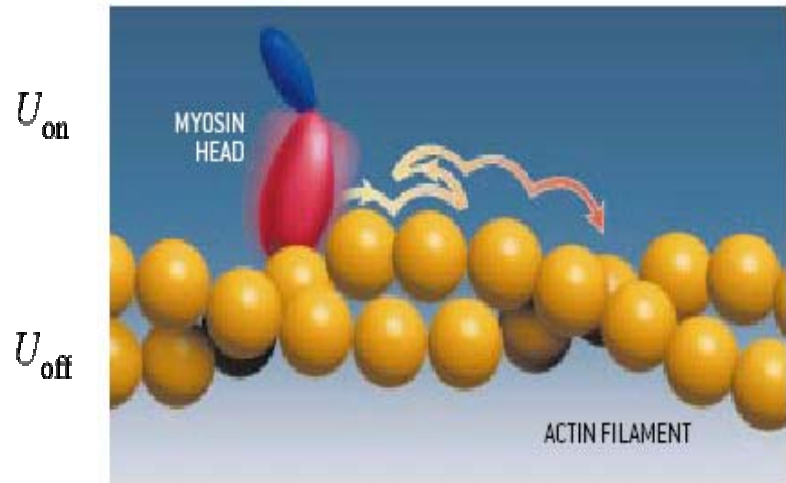
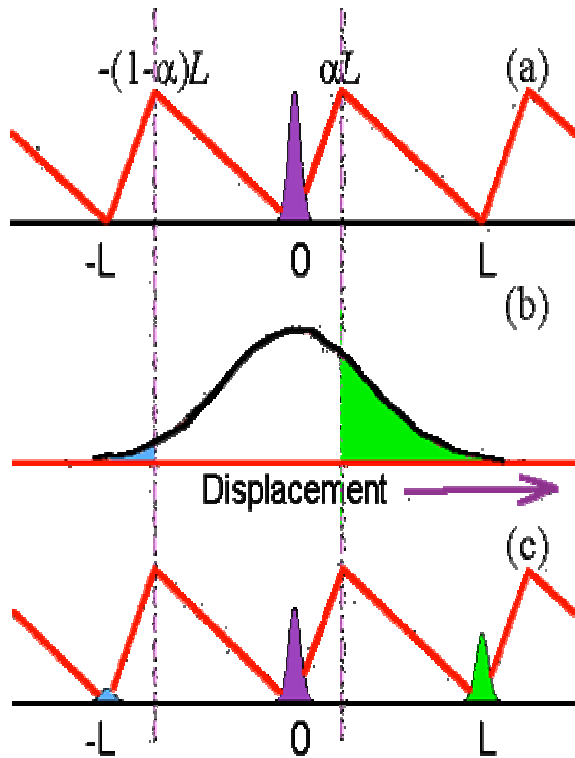
Frank Jülicher, et al Modeling molecular motors, Rev. Mod. Phys. 69, 1269 (1997)

Related problems:

Focker-Planck equation, Dissipation fluctuation relation, Einstein relation, ...

Flashing ratchet & Muscle contraction

Depending on whether ATP is on or away from the myosin head, the head is dissociated from or attached to the filament and experiences a flat respectively ratchet potential.



Note that without noise, the myosin head cannot move !!

Problem in energy transduction:

Chemical bond \leftrightarrow mechanical work

From an energetic point of view, however, it seems unlikely that the protein is able to extract, in an efficient manner, the free energy released during ATP catalysis, since the latter is an **ultrafast process** taking place on a **femtosecond** time scale. This was recently corroborated by computer simulations (Dittrich *et al.*, 2003, 2004) which revealed that **no net free energy is released during ATP catalysis itself**. In their paper (Dittrich *et al.*, 2003, 2004), the authors showed that the catalysis reaction energy profile changed from strongly endothermic in the β_{TP} catalytic site to approximately equienergetic in β_{DP} . This leaves either reactant binding (ATP), product unbinding (ADP, Pi), or both as possible candidates for force generation.

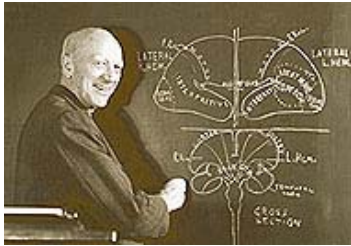
(Markus Dittrich and Klaus Schulten

Journal of Bioenergetics and Biomembranes, Vol. 37, No. 6, 441, 2005)

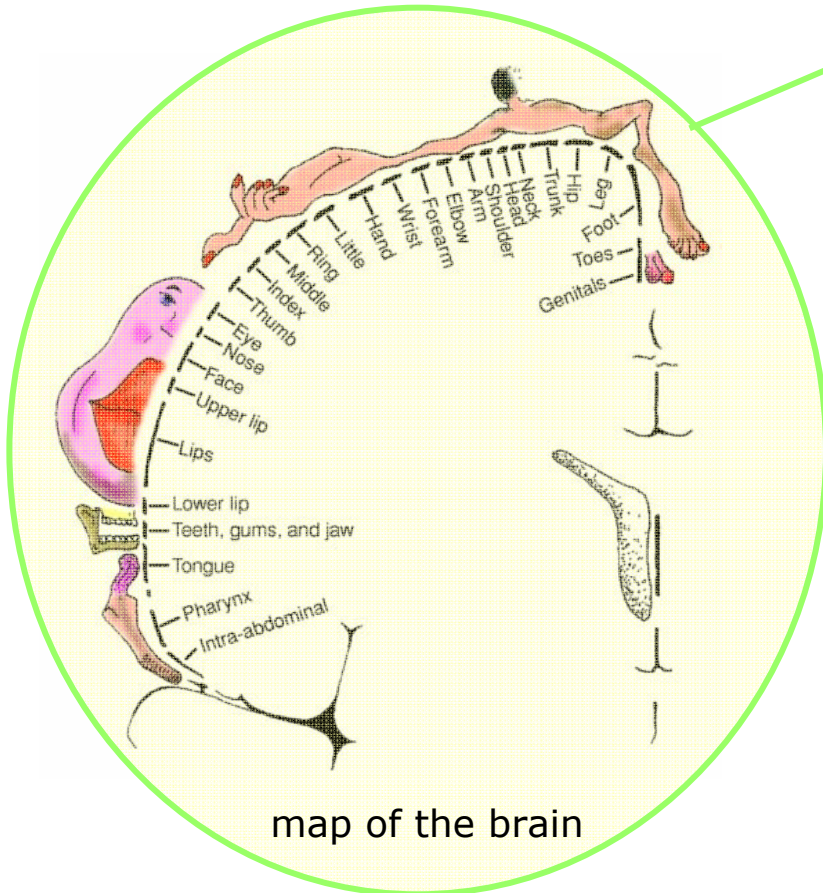
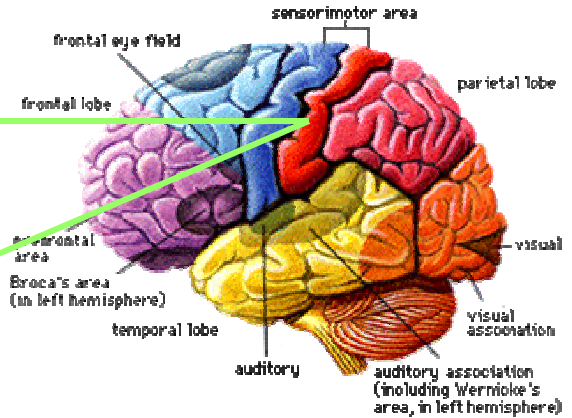
5. Electrostatics (neuron cells & networks)



Perception of electrical stimulation



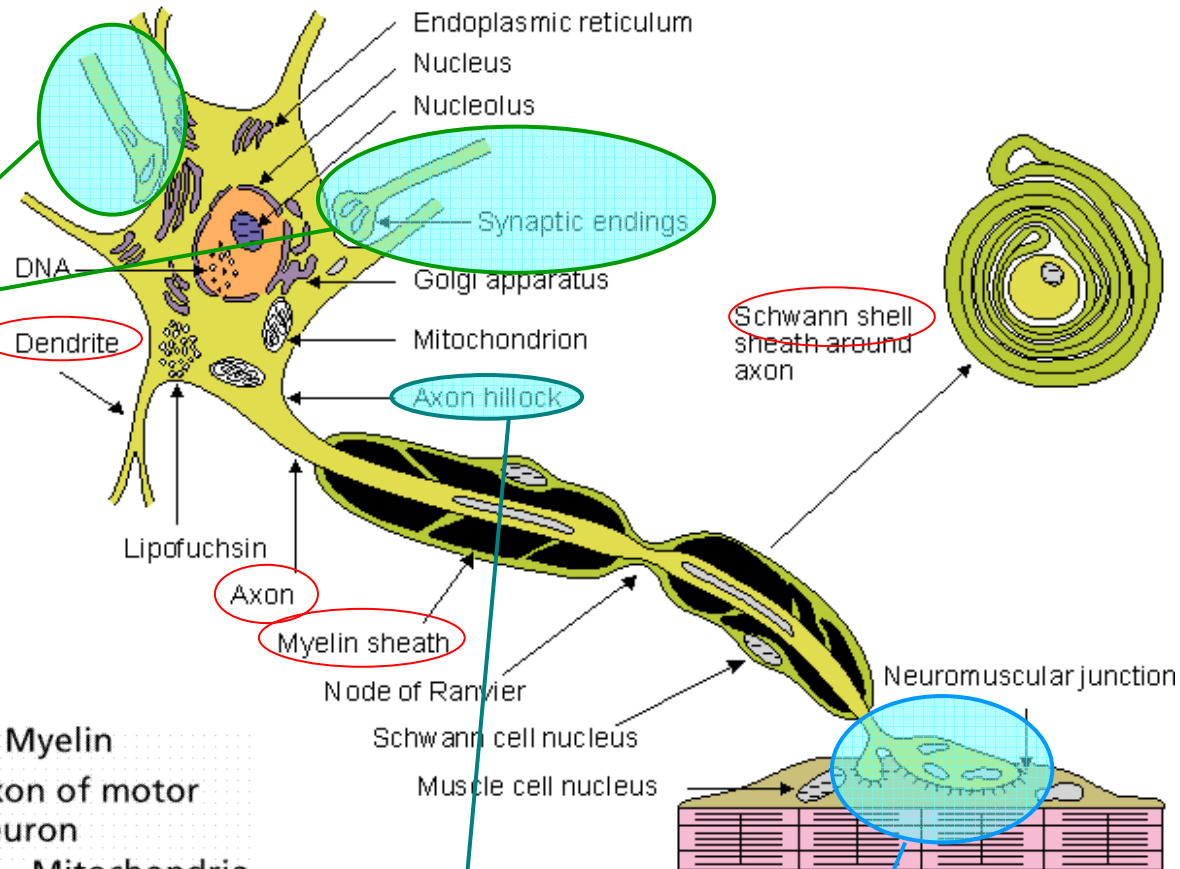
Dr. Wilder Penfield stimulated different areas of the **primary somatosensory cortex** of his patient and drew the map of the brain:



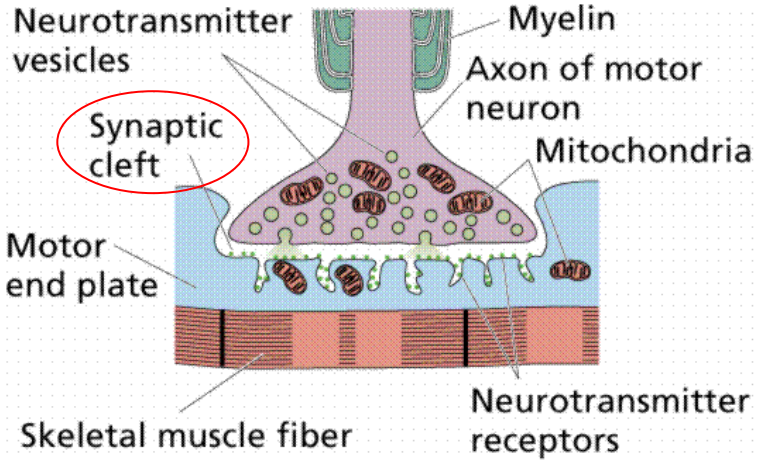
Homunculus (小矮人)



Neuron cell



Signal input:
Collect signals
from the synapse
of other cells.



Signal control:
Fire a signal,
if the total
strength of
the signal
exceeds
a threshold.

Signal output:
Release signal through
its own synapse to
other cells.

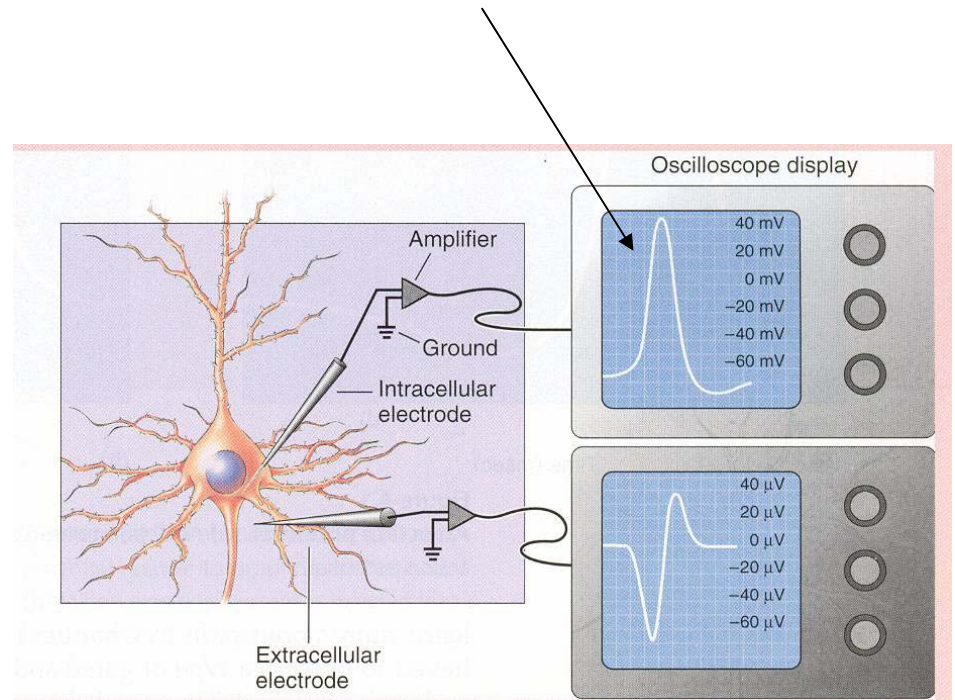
Measurement of action potentials

Action potentials are measured with the recording techniques. An oscilloscope recording the membrane potential from a single point on an axon shows each phase of the action potential as the wave passes.

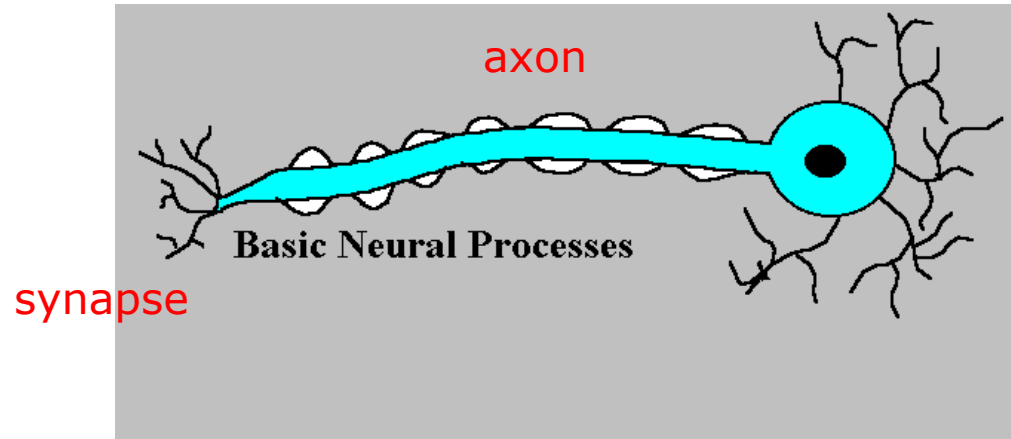
These phases trace an arc that resembles a **distorted sine wave**.

Its amplitude depends on whether the action potential wave has reached that point or passed it and how long ago.

The speed and simplicity of action potentials vary between different types of cells. However, the amplitudes of the voltage swings tend to be roughly the same. Within any one cell, consecutive action potentials typically are indistinguishable.



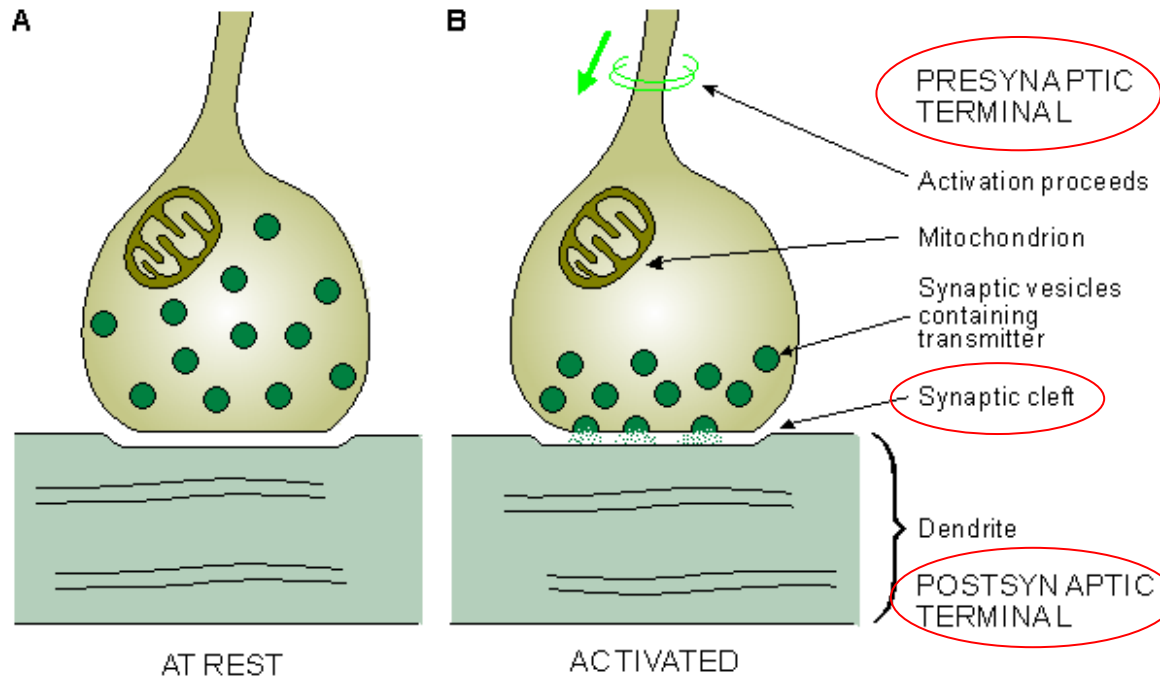
Internal neuron signal transmission (through axon)



Action potential

- is a wave of **electrical discharge** that travels along the membrane of a cell.
- carries fast internal messages between tissues and can be created by many types of body cells.
- is an essential carrier of the neural signal. Its properties may enable centralized control and coordination of organs and tissues.

External signal transmission (through synapse)



A The synaptic vesicles contain a chemical transmitter.

B When the activation reaches the presynaptic terminal the transmitter is released and it diffuses across the synaptic cleft to activate the postsynaptic membrane.

The formation of the action potential

When the membrane potential of an excitable cell is depolarized beyond a threshold, the cell will undergo an action potential (fire a spike).

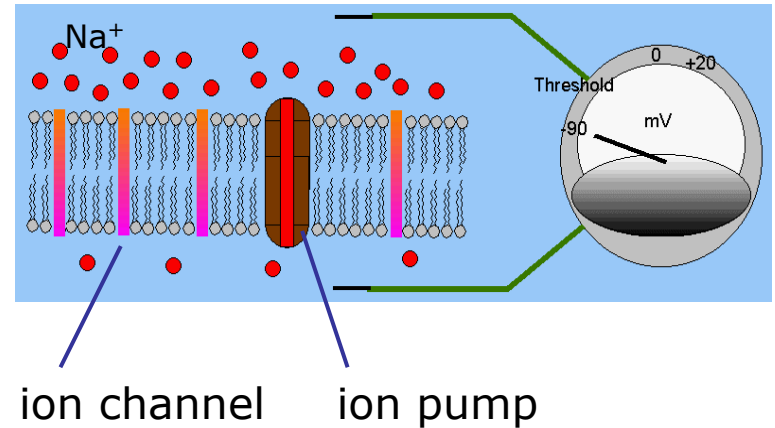
Four activities:

- (1) The outside and inside of the membrane is at a resting potential.
- (2) Sodium (Na^+) moves inside the cell causing an action potential, the influx of positive sodium ions makes the inside of the membrane more positive than the outside.
- (3) Potassium (K^+) ions flow out of the cell, restoring the resting potential net charges.
- (4) Na^+ are pumped out of the cell and K^+ are pumped into the cell, restoring the original distribution of ions.

Three phases:

These four activities will lead to four phases in a cycle:

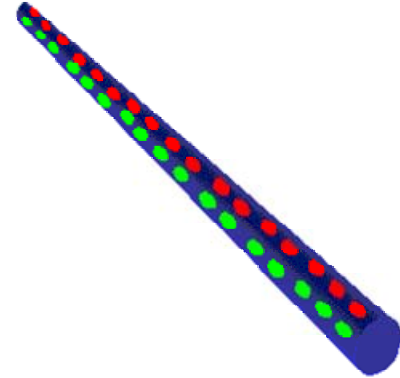
- (A) resting phase
- (B) rising phase
- (C) falling phase.



Propagating direction of the action potential

Usually action potentials travel unidirectionally along the axon.

Reason: where membrane has undergone an action potential, a **refractory period** follows. This period arises primarily because of the voltage-dependent inactivation of sodium channels (Hodgkin and Huxley in 1952).



(i) **Absolute refractory period:** immediately after an action potential, virtually all sodium channels are inactivated and thus it is impossible to fire another action potential in that segment of membrane.

(ii) **Relative refractory period:** with time, sodium channels are reactivated in a stochastic manner and as they become available, it becomes possible to fire an action potential, though one with a much higher threshold.

These two refractory periods together last approximately 5×10^{-3} seconds.

Propagating velocity of the action potential

For a unmyelinated axon,

$$v = \sqrt{\frac{i_{Na \max}}{r_i c_m^2 V_{th}}}$$

where

v = velocity of the nerve impulse [m/s]

$i_{Na \max}$ = maximum sodium current per unit length [A/m]

V_{th} = threshold voltage [V]

r_i = axial resistance per unit length [W/m]

c_m = membrane capacitance per unit length [F/m]

Example: for unmyelinated sensory neurons is 5 - 25 m / s.

For a myelinated axon:

$$v = 6d$$

where

v = velocity of the nerve impulse [m/s]

d = axon diameter \sim 2-20 [μ m]

Example: for mammalian motor neurons is 10 - 120 m / s.

Question: Can a baseball hitter hit a ball of 90 m/s?

The distance between the pitcher and the home base is 18.44 m.

The time before the ball arrives at the hitter is $18/90 = 0.2$ s.

The propagating distance of the action potential is $10 \times 0.2 = 2$ m.

Hodgkin-Huxley model

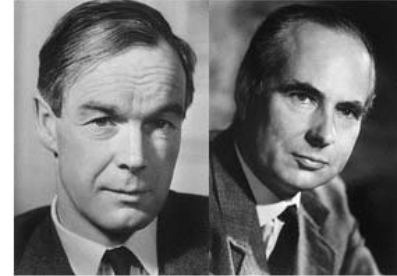
Goal: understanding the propagation of action potentials
(originally on the giant axon of the squid).

The semipermeable cell membrane separates the interior of the cell from the extracellular liquid and acts as a capacitor.

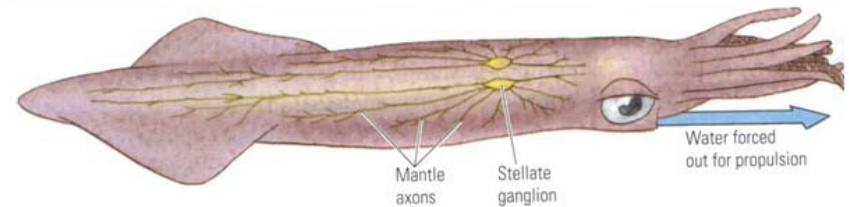
If an input current is injected into the cell, it may add further charge on the capacitor, or leak through the channels in the cell membrane.

Because of active ion transport through the cell membrane, the ion concentration inside the cell is different from that in the extracellular liquid.

The Nernst potential generated by the difference in ion concentration is represented by a battery.



Hodgkins and Huxley
Nobel Prize, 1963



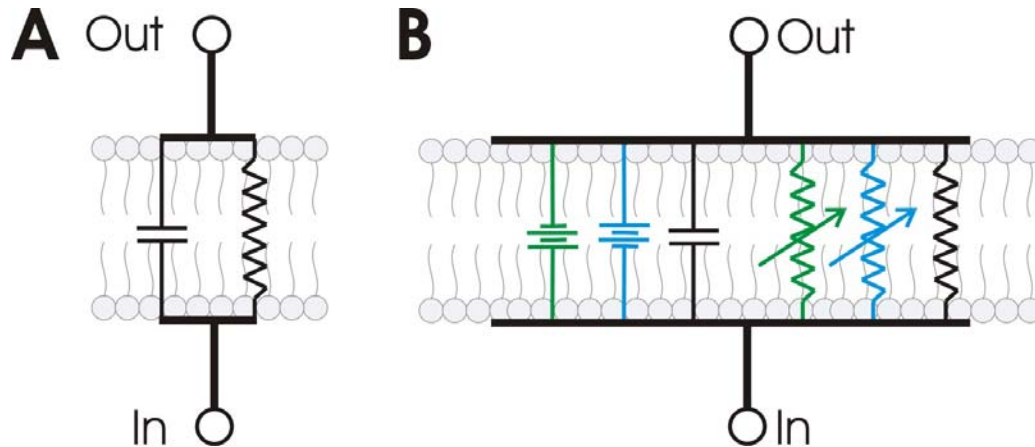
RC Circuit model

Membrane = insulator with **capacitor** and **high fixed resistor**

Voltage-gated channel = conductor with **low various resistor**

Leak channel = conductor with **low fixed resistor**

Voltage gradient of Na^+ and K^+ across the membrane = **battery**



A. Membrane without ion channels without transmembrane ion gradient.

B. Membrane with Na^+ (blue) and K^+ (green) ion channels and gradient.

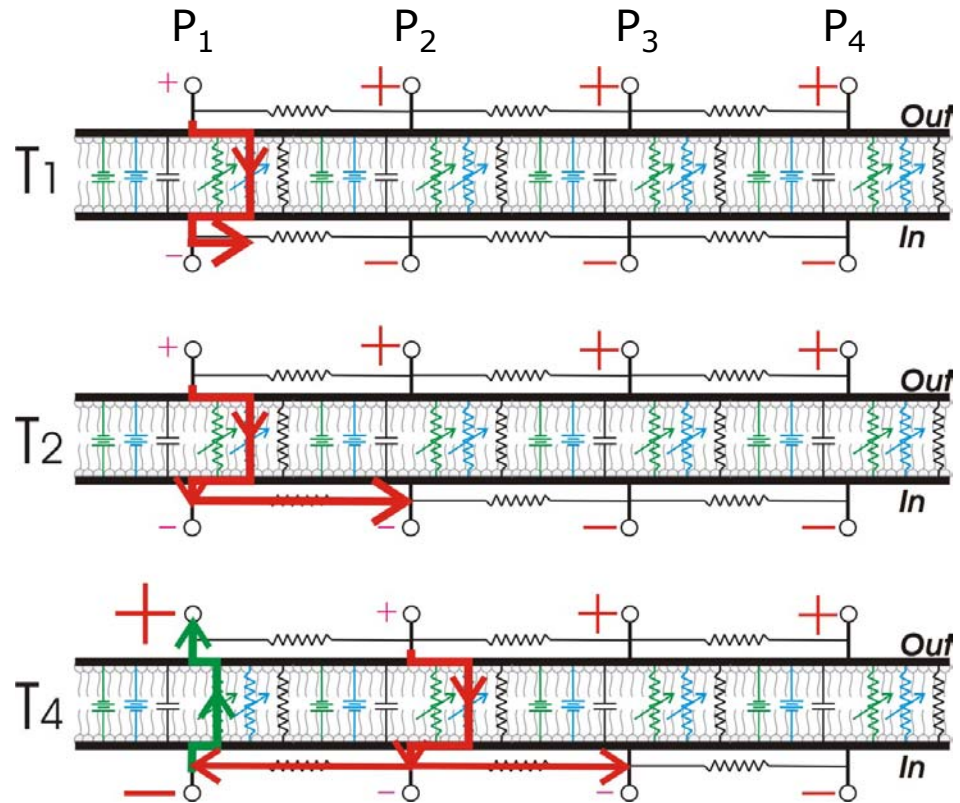
Propagating depolarization in the circuit model

(T1) A local depolarization opens local Na^+ channels at P_1 .

(T2) Na^+ current spreads to the right adjacent membrane at P_2 and causes a new depolarization there.

(T4) This depolarization opens the Na^+ channels at P_2 and new Na^+ current to cause new depolarization spreading further down the membrane to P_3 and P_4 .

Meanwhile, the delayed K^+ current flows in the membrane patch at P_1 , causing the falling phase of the action potential.



Note the size change of '+' and '-'.

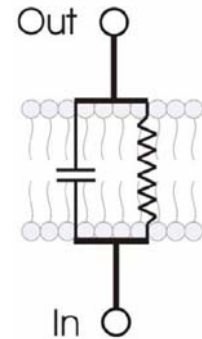
Mathematical Model

An input current $I(t)$ onto the membrane may be split in a capacitive current I_C which charges the capacitor C and further components I_k which pass through the ion channels. Thus

$$I(t) = I_C(t) + \sum_k I_k(t),$$

where

- $I_C = C \, du/dt$ with the membrane capacity $C = Q/u = 1 \text{ F/cm}^2$, in which Q is a charge and u is the voltage across the capacitor (membrane).
- the sum of $I_k(t)$ runs over the following 3 ion channels:
 - (i) an unspecific leakage channel with a constant resistance R and a constant conductance $g_L = 1/R$,
 - (ii) two voltage-gated channels (Na channel & K channel) with a time-dependent resistance and conductance, Suppose their maximum conductances are g_{Na} and g_K .



The time-dependent conductance is determined by the open probability m , n , and h of the channels with

$$\begin{aligned}
 dm/dt &= \alpha_m(u)(1-m) - \beta_m(u)m && \text{(fast Na}^+ \text{ gate)} \\
 dn/dt &= \alpha_n(u)(1-n) - \beta_n(u)n && \text{(slow Na}^+ \text{ gate)} \\
 dh/dt &= \alpha_h(u)(1-h) - \beta_h(u)h && \text{(K}^+ \text{ gate)}
 \end{aligned} \tag{1}$$

where α_i and β_i are some functions of the transmembrane voltage u .

Hodgkin and Huxley formulated the three current components as

$$\Sigma_k I_k = g_{Na} m^3 h (u - E_{Na}) + g_K n^4 (u - E_K) + g_L (u - E_L),$$

with the reversal potentials E_{Na} , E_K , and E_L which together with g_{Na} , g_K , and g_L as well as α_i and β_i are empirical (Hodgkin and Huxley):

| x | E_x | g_x | x | $\alpha_x(u / \text{mV})$ | $\beta_x(u / \text{mV})$ |
|-----|---------|------------------------|-----|-------------------------------------------|-----------------------------|
| Na | 115 mV | 120 mS/cm ² | n | $(0.1 - 0.01 u) / [\exp(1 - 0.1 u) - 1]$ | $0.125 \exp(-u / 80)$ |
| K | -12 mV | 36 mS/cm ² | m | $(2.5 - 0.1 u) / [\exp(2.5 - 0.1 u) - 1]$ | $4 \exp(-u / 18)$ |
| L | 10.6 mV | 0.3 mS/cm ² | h | $0.07 \exp(-u / 20)$ | $1 / [\exp(3 - 0.1 u) + 1]$ |

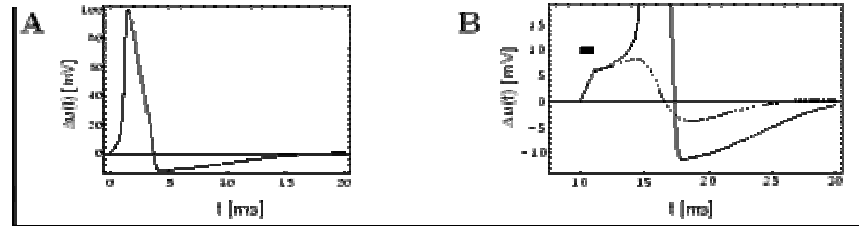
Dynamics

Using these parameters in the Hodgkin-Huxley equation:

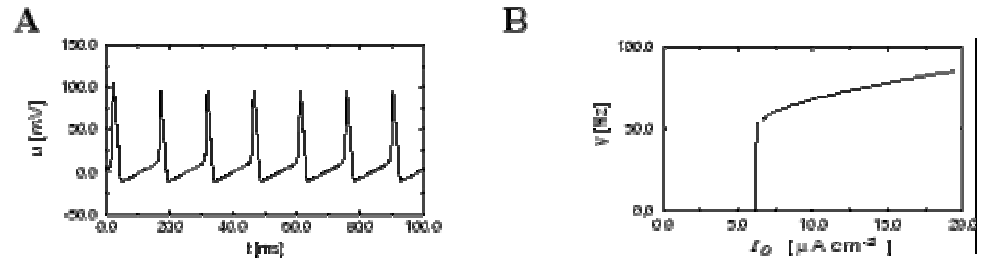
$$C \frac{du}{dt} = - \sum_k I_k(t) + I(t).$$

many important features of action potential can be realized:

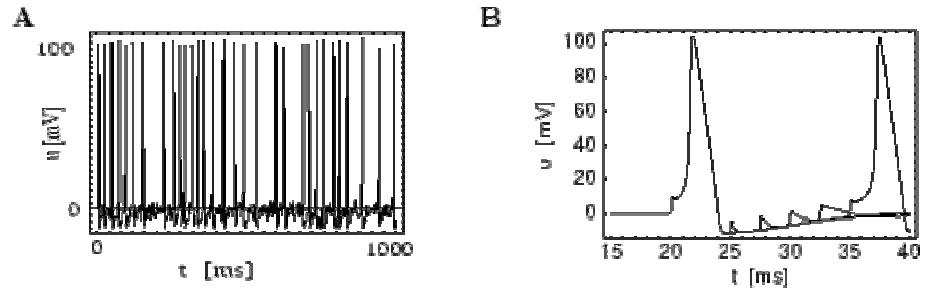
1. Spike generation



2. Mean firing rate



3. Stimulation by time-dependent input



Models outside the neuron cell

Rate model

TABLE I. Summary of many frequently used neuronal models.

| Model | Example | Variables | Remarks | References |
|----------------------------|--------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|-----------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|--------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|-----------------------------------------------------|
| Integrate-and-fire neurons | $\frac{dv(t)}{dt} = \begin{cases} -\frac{v(t)}{\tau} + I_{\text{ext}} + I_{\text{syn}}(t), & 0 < v(t) < \theta \\ v'(t_0) = 0, & v(t_0) = \theta \end{cases}$ $I_{\text{syn}}(t) = g \sum_{\text{spikes}} f(t - t_{\text{spike}})$ and $f(t) = A[\exp(-t/\tau_1) - \exp(-t/\tau_2)]$ | $v(t)$ is the neuron membrane potential; θ is the threshold for spike generation. I_{ext} is an external stimulus current; I_{syn} is the sum of the synaptic currents; and τ_1 and τ_2 are time constants characterizing the synaptic currents. | A spike occurs when the neuron reaches the threshold θ in $v(t)$ after which the cell is reset to the resting state. | Lapicque, 1907 |
| Rate models | $\dot{a}_i(t) = F_i(a_i(t)) [G_i(a_i(t)) - \sum_j \rho_{ij} Q_j(a_j(t))]$ | $a_i(t) > 0$ is the spiking rate of the i th neuron or cluster; ρ_{ij} is the connection matrix; and F, G, Q are polynomial functions. | This is a generalization of the Lotka-Volterra model [see Eq. (9)]. | Fukai and Tanaka, 1997; Lotka, 1925; Volterra, 1931 |
| McCulloch and Pitts | $x_i(n+1) = \Theta(\sum_j g_{ij} x_j(n) - \theta)$ $\Theta(x) = \begin{cases} 1, & x > 0 \\ 0, & x \leq 0 \end{cases}$ | θ is the firing threshold; $x_i(n)$ are synaptic inputs at the discrete "time" n ; $x_i(n+1)$ is the output. Inputs and outputs are binary (one or zero); the synaptic connections g_{ij} are 1, -1, or 0. | The first computational model for an artificial neuron; it is also known as a linear threshold device model. This model neglects the relative timing of neural spikes. | McCulloch and Pitts, 1943 |
| Hodgkin-Huxley | $Cv'(t) = g_L[v_L - v(t)] + g_{Na}m(t)^3h(t)[v_{Na} - v(t)] + g_{Kd}n(t)^4(v_K - v(t)) + I$ $m'(t) = \frac{m_{\infty}(v(t)) - m(t)}{\tau_m(v(t))}$ $h'(t) = \frac{h_{\infty}(v(t)) - h(t)}{\tau_h(v(t))}$ $n'(t) = \frac{n_{\infty}(v(t)) - n(t)}{\tau_n(v(t))}$ | $v(t)$ is the membrane potential, $m(t)$, and $h(t)$, and $n(t)$ represent empirical variables describing the activation and inactivation of the ionic conductances; I is an external current. The steady-state values of the conductance variables $m_{\infty}, h_{\infty}, n_{\infty}$ have a nonlinear voltage dependence, typically through sigmoidal or exponential functions. | These ODEs represent point neurons. There is a large list of models derived from this one, and it has become the principal tool in computational neuroscience. Other ionic currents can be added to the right-hand side of the voltage equation to better reproduce the dynamics and bifurcations observed in the experiments. | Hodgkin and Huxley, 1952 |
| FitzHugh-Nagumo | $\dot{x} = \mu x - cx^3 - y + I, \quad \dot{y} = x + by - a$ | $x(t)$ is the membrane potential, and $y(t)$ describes the dynamics of fast currents; I is an external current. The parameter values a, b , and c are constants chosen to allow spiking. | A reduced model describing oscillatory spiking neural dynamics including bistability. | FitzHugh, 1961; Nagumo et al., 1962 |

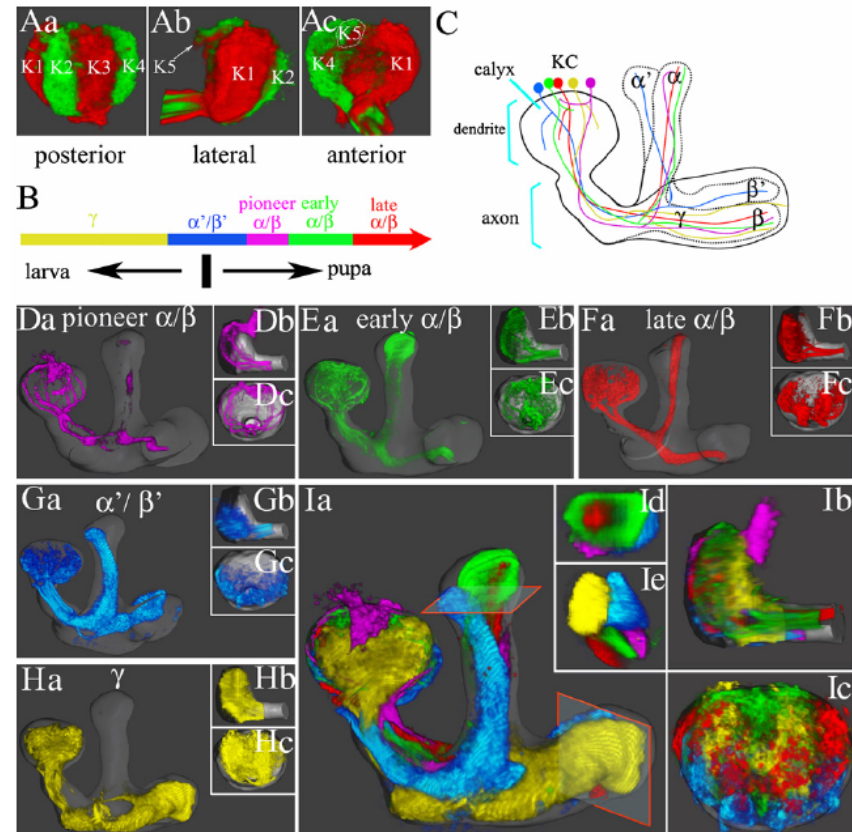
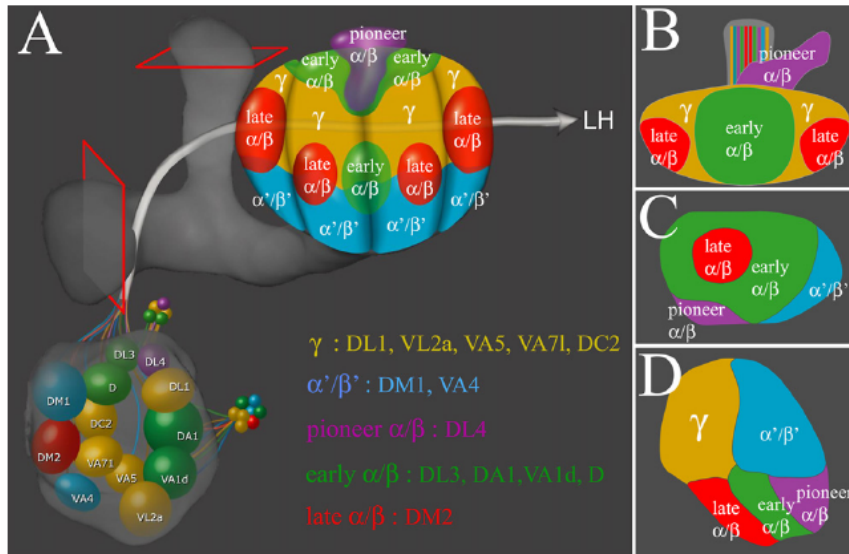
TABLE I. (Continued.)

| Model | Example | Variables | Remarks | References |
|-------------------------|------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|---------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|-----------------------------------------------------------------|
| Wilson-Cowan | $\mu \frac{\partial E(x,t)}{\partial t} = -E(x,t) + [1 - rE(x,t)] \times \mathcal{L} [E(x,t) \otimes w_{ee}(x) - I(x,t) \otimes w_{ei}(x) + I_e(x,t)]$ $\mu \frac{\partial I(x,t)}{\partial t} = -I(x,t) + [1 - rI(x,t)] \times \mathcal{L} [E(x,t) \otimes w_{ie}(x) - I(x,t) \otimes w_{ii}(x) + I_i(x,t)]$ | $\{E(x,t), I(x,t)\}$ are the number density of active excitatory and inhibitory neurons at location x of the continuous neural media. ($w_{ee}(x), w_{ie}(x), w_{ei}(x), w_{ii}(x)$) are connectivity distributions among the populations of cells. $\{\mathcal{L}_e, \mathcal{L}_i\}$ are nonlinear responses reflecting different populations of thresholds. The operator \otimes is a convolution involving the connectivity distributions. | The first "mean-field" model. It is an attempt to describe a cluster of neurons, to avoid the inherent noisy dynamical behavior of individual neurons; by averaging to a distribution noise is reduced. | Wilson and Cowan, 1973 |
| Morris-Lecar | $v'(t) = g_L[v_L - v(t)] + n(t)g_n \times [v_n - v(t)] + g_{Ca}m_{\infty}(v(t))[v_m - v(t)] + I$ $n'(t) = \lambda(n(v(t))[n_{\infty}(v(t)) - n(t)]$ $m_{\infty}(v) = \frac{1}{2} \left(1 + \tanh \frac{v - v_m}{v^0 - v_m} \right)$ $n_{\infty}(v) = \frac{1}{2} \left(1 + \tanh \frac{v - v_n}{v^0 - v_n} \right)$ $\lambda(v) = \phi_n \cosh \frac{v - v_n}{2v^0 - v_n}$ | $v(t)$ is the membrane potential; $n(t)$ describes the recovery activity of a calcium current; I is an external current. | Simplified model that reduces the number of dynamical variables of the HH model. It displays action potential generation when changing I leads to a saddle-node bifurcation to a limit cycle. | Morris and Lecar, 1981 |
| Hindmarsh-Rose | $x'(t) = y(t) + ax(t)^2 - bx(t)^3 - z(t) + I$ $y'(t) = C - xx(t)^2 - y(t)$ $z'(t) = r[s(x(t) - x_0) - z(t)]$ | $x(t)$ is the membrane potential; $y(t)$ describes fast currents; $z(t)$ describes slow currents; and I is an external current. | Simplified model that uses a polynomial approximation to the right-hand side of a Hodgkin-Huxley model. This model fails to describe the hyperpolarized periods after spiking of biological neurons. | Hindmarsh and Rose, 1984 |
| Phase oscillator models | $\frac{d\theta_i(t)}{dt} = \omega + \sum_j H_{ij} \theta_j(t) - \theta_i(t)$ | $\theta(t)$ is the phase of the i th neuron with approximately periodic behavior; and H_{ij} is the connectivity function determining how neuron i and j interact. | First introduced for chemical oscillators; good for describing strongly dissipative oscillating systems in which the neurons are intrinsic periodic oscillators. | Cohen et al., 1982; Ermentrout and Kopell, 1984; Kuramoto, 1984 |
| Map models | $x_{i+1}(t) = \frac{\alpha}{1 + x_i(t)^2} + y_i(t) + \frac{\epsilon}{N} \sum_j x_j(t)$ $y_{i+1}(t) = y_i(t) - \alpha x_i(t) - \beta$ | x_i represents the spiking activity and y_i represents a slow variable. A discrete time map. | One of a class of simple phenomenological models for spiking, bursting neurons. This kind of model can be computationally very fast, but has little biophysical foundation. | Cazelles et al., 2001; Rulkov, 2002 |

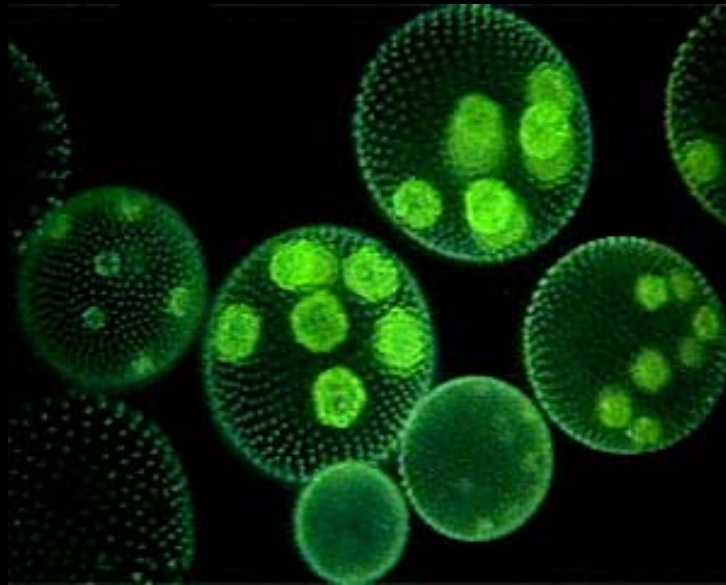
Neuron network

A Schematic Map of Olfactory Representations in the Adult Brain

For the 13 PNs analyzed, each PN sends dendrites to a single antennal glomerulus and an axon giving one to four major branches in the MB calyx. KC dendrites are segregated into 17 complementary domains defined by two orthogonal factors (four vertical K1–K4 clonal divisions by four horizontal sequential birth domains plus one 4-fold K5 domain).



6. Thermodynamics (non-equilibrium biosystems)



Useful work and free energy

The 2. law tells us

$$dS_o + dS \geq 0$$

Since

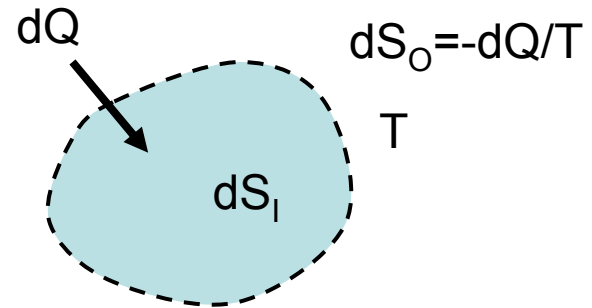
$$dS_o = -dQ/T$$

and

$$dU = dQ + dW$$

we have

$$dU - dW - T dS \leq 0.$$



Since the work W includes the hydrostatic work $-PdV$ performed on the system and other non-hydrostatic work W' , it yields

$$dU + P dV - d W' - T dS \leq 0.$$

At constant P and T ,

$$dW' \geq d (U - TS + PV) = dG. \quad (\text{Gibbs free energy})$$

For system in which P and V are irrelevant, $W = W'$ and

$$dW \geq d (U-TS) = dF. \quad (\text{Helmholtz free energy})$$

That is in a thermodynamic process from a state (macro) to another state, the work performed on the system is larger than the difference of the free energies between these two states:

$$W \geq \Delta G, \quad (\text{where } W' \text{ has been simplified as } W).$$

The inequality is ascribed to the inequality in the 2. law:

$$\Delta S \geq 0.$$

These two relations become equalities when the process is **reversible**, for which the process has to be infinitely slow so that

$$\Delta S = \Delta Q/T = (\Delta U - W) / T = 0 \quad \Rightarrow \quad \Delta U = W.$$

That is the work from the environment can be perfectly transformed to and completely stored in the internal energy.

In such a reversible process,

$$\Delta G = W_{\text{rev}}.$$

Jarzynski's Equality

Therefore, the external work W performed on the system can be used to estimate the maximum possible value of the free energy difference:

$$\Delta G \leq \langle W \rangle$$

(Gibbs free energy $G = U + PV - TS$)

where the $\langle \rangle$ denotes an average over several experimental measurements under identical conditions.

- Since that work given to the system is bound to dissipate away partially, this inequality is quite understandable.
- Determining free energy difference from measuring the work done on the system will require infinitely slow processes and therefore is impossible in practice.
- Nevertheless, recently Jarzynski showed that the identity for any process:

$$\exp\left(-\frac{\Delta G}{k_B T}\right) = \left\langle \exp\left(-\frac{W}{k_B T}\right) \right\rangle$$

- This gives the important implication that the equilibrium quantity ΔF or ΔG can be extracted from non-equilibrium (real) measurements.

RNA Experiment on Jarzynski equality

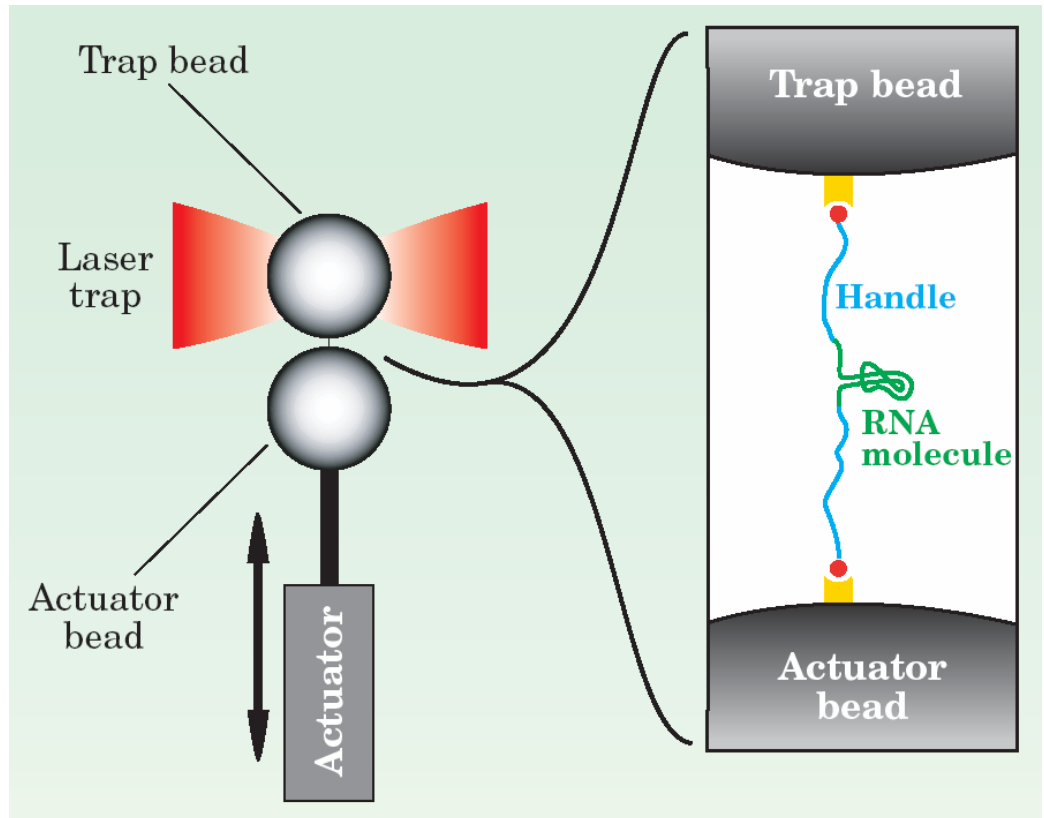
A molecule of RNA is attached to two beads and subjected to reversible and irreversible cycles of folding and unfolding.

A piezoelectric actuator controls the position of the bottom bead.

An optical trap captures the top bead and determines the force exerted on the RNA.

The position difference of the two beads gives the end-to-end length of the molecule.

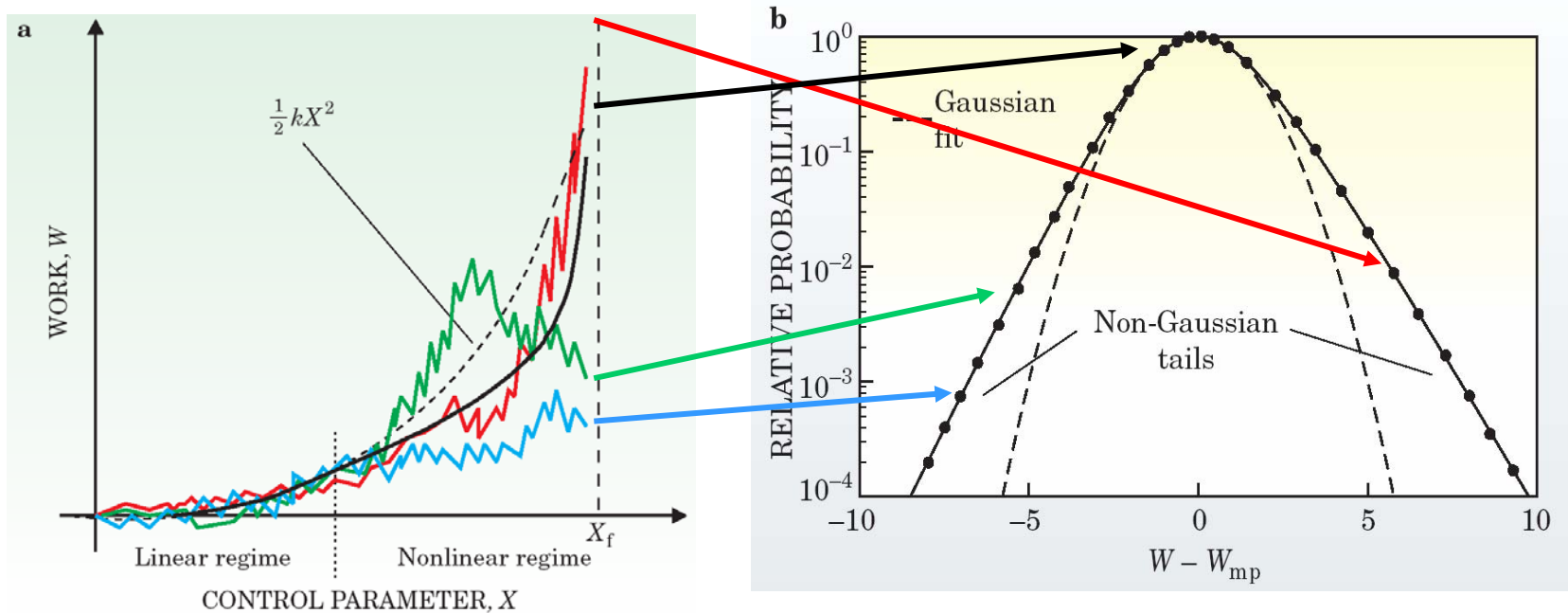
The diameter of the beads is around 3000 nm, much greater than the 20-nm length of the RNA.



The work to stretch a polymer

a. The work required to stretch a short polymer is a fluctuating function, As shown in the three different nonequilibrium trajectories obtained as the control parameter X varies from 0 to X_f . The continuous black line is the work averaged over all trajectories.

b. In the nonlinear regime, the work probability distribution $P(W)$ has a Gaussian component plus long non-Gaussian tails describing rare processes.



Crooks fluctuation theorem

Jarzynski's equality can be proved from the Crooks theorem.

Consider the forward (F) and reverse paths (R) with:

$$x_F(t) \quad \text{and} \quad x_R(t) = x_F(t_f - t).$$

At the beginning of the F and R paths the system is in equilibrium with the bath at temperature T.

Crooks fluctuation theorem:

$$\frac{P_F(W)}{P_R(-W)} = \exp\left(\frac{W - \Delta G}{k_B T}\right).$$

Rewriting this relation as

$$P_F(W) \exp(\Delta G/k_B T) = P_R(-W) \exp(W/k_B T),$$

and integrating both sides along paths, one gets the Jarzynski's equality after some steps of calculations.

Physical problems in Bio system

Molecular level: Quantum mechanics

Cell level: Mechanics, Electrodynamics, Thermodynamics

Tissue level: Nonlinear science, statistical mechanics