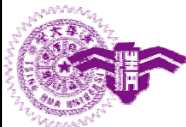

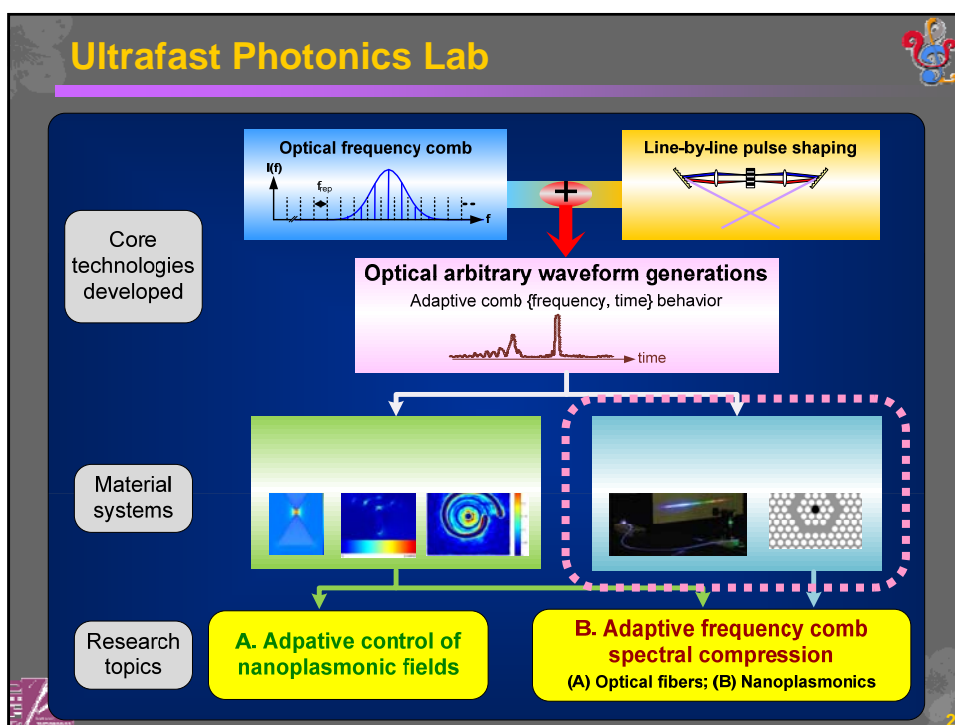


AMO Seminar, Department of Physics, National Tsing Hua University

Large-scale laser spectral compression and its applications

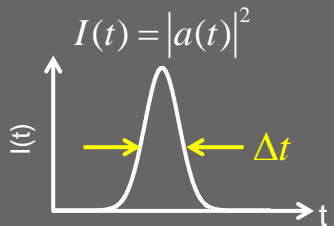
Chen-Bin Huang (黃承彬)

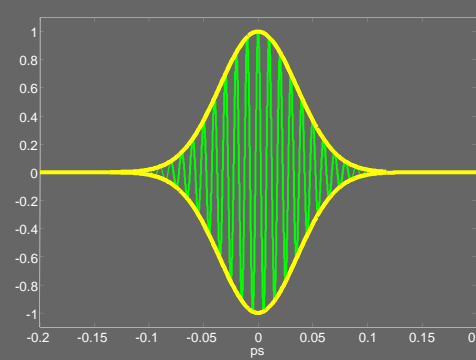
Institute of Photonics Technologies
National Tsing Hua University, Taiwan
March 26, 2012



Ultrashort optical pulse

- Electric field
 - Carrier
 - Envelope function $a(t)$

$$I(t) = |a(t)|^2$$


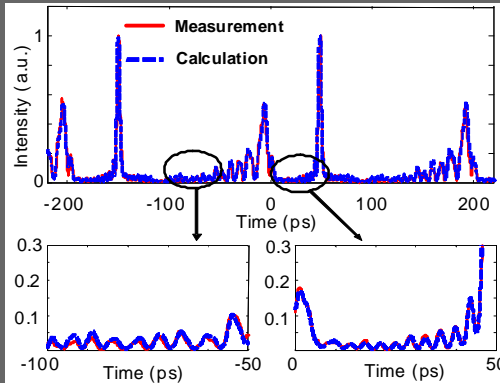
$$E(t) = \text{Re}\{a(t)\exp^{i\omega t}\}$$


- Ultrafast optics
 - Full-width half-maximum pulse duration $\Delta t < 1\text{ps}$

Optical pulse shaping

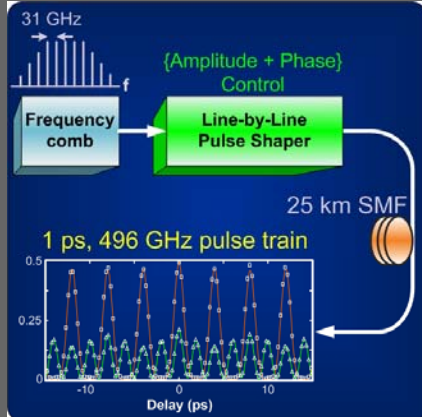
- We can shape the optical waveform however you like



Intensity (a.u.) vs Time (ps)

— Measurement (red)
— Calculation (blue)

Time (ps) ranges from -200 to 200.



31 GHz

Frequency comb f

{Amplitude + Phase} Control

Line-by-Line Pulse Shaper



25 km SMF

1 ps, 496 GHz pulse train

Delay (ps)

Jiang*, Huang*, Leaird, Weiner, Nat. Photon. 1, 463 (2007)

Chuang and Huang*, Opt. Express 18, 24003 (2010)

Outline

- Ultrafast 3rd-order nonlinear optics
 - Self-phase modulation
 - Optical solitons
 - Raman scattering
 - Self-steepening
- Supercontinuum generation
- Spectral compression
- Summary



5

Nonlinear wave propagation equation

- The inclusion of nonlinear polarization

$$\nabla \times \nabla \times \vec{E} = \nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2}$$

$$\vec{D} = \varepsilon_0 \vec{E} + \vec{P}_{(1)} + \vec{P}_{NL} = \vec{D}_{(1)} + \vec{P}_{NL}$$

$$\nabla_T^2 \vec{E} + \frac{\partial^2 \vec{E}}{\partial z^2} - \mu_0 \varepsilon_{(1)} \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

- Scalar treatment

$$D \approx \varepsilon_0 (n_0^2 + 2n_0 \delta n_{(1)} + 2n_0 \delta n_{NL}) E \approx (\varepsilon_{(1)} + 2\varepsilon_0 n_0 \delta n_{NL}) E$$

$$P_{NL} \approx 2\varepsilon_0 n_0 \delta n_{NL} E$$



6

Nonlinear propagation equation in waveguide

- Basic formulation

$$\frac{\partial a}{\partial z} + \beta_1 \frac{\partial a}{\partial t} - \frac{j\beta_2}{2} \frac{\partial^2 a}{\partial t^2} + j\gamma |a|^2 a + \frac{\alpha}{2} a = 0 \quad \gamma = \frac{\omega_0^2 n_0 n_2}{\beta_0 c^2 A_{eff}}$$

optical Kerr term
- NonLinear Schrödinger Equation

$$\frac{\partial a}{\partial z'} - \frac{j\beta_2}{2T_0^2} \frac{\partial^2 a}{\partial \tau^2} + j\gamma |a|^2 a = 0$$
 - Lossless
- Generalized NLSE

$$\frac{\partial A}{\partial z} + \frac{\alpha}{2} A - \frac{j}{2} \beta_2 \frac{\partial^2 a}{\partial T^2} - \frac{1}{6} \beta_3 \frac{\partial^3 a}{\partial T^3} = -j\gamma \left[\underbrace{|a|^2 a}_{\text{dispersion}} - \underbrace{T_R a \frac{\partial(|a|^2)}{\partial T}}_{\text{SPM}} - \underbrace{\frac{j}{\omega_0} \frac{\partial(|a|^2 a)}{\partial T}}_{\text{Raman}} \right]$$

Self-steepening

Self-phase modulation (SPM)

- No dispersion

$$\frac{\partial a}{\partial z} = -j\gamma |a|^2 a$$
- Gives broadened spectrum

$$a(z, t) = a(0, t) \exp[-j\gamma |a(0, t)|^2 z]$$
- instantaneous frequency:

$$\Delta\omega(z, t) = \frac{\partial(\Delta\phi)}{\partial t} = -\frac{\partial}{\partial t} (\gamma |a|^2 z)$$

Optical solitons

- For anomalous dispersive materials: $\beta_2 < 0$
- Dimensionless NLSE
- Solution: solitons

$$j \frac{\partial a}{\partial \zeta} = \frac{1}{2} \frac{\partial^2 a}{\partial \tau^2} + |a|^2 a$$

$$a(z, t) = \sqrt{P_c} \operatorname{sech}\left(\frac{t}{T_0}\right) \exp\left[\frac{-j|\beta_2|z}{2T_0^2}\right]$$

- Complete balance between dispersion and SPM

- Pulse energy

$$U = \frac{2|\beta_2|}{\gamma T_0}$$

- Soliton order

$$N^2 = \frac{\gamma T_0 P_{peak}}{|\beta_2|}$$



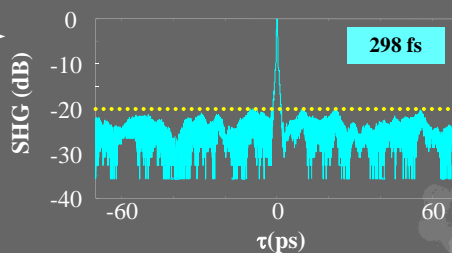
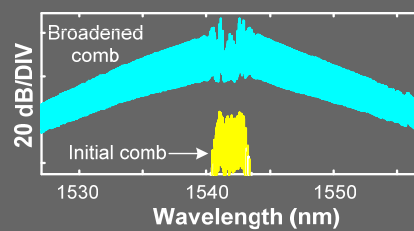
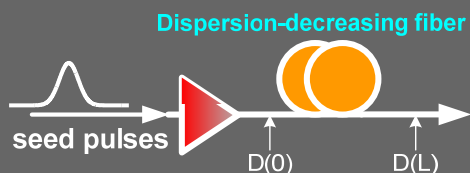
9

Adiabatic soliton temporal compression

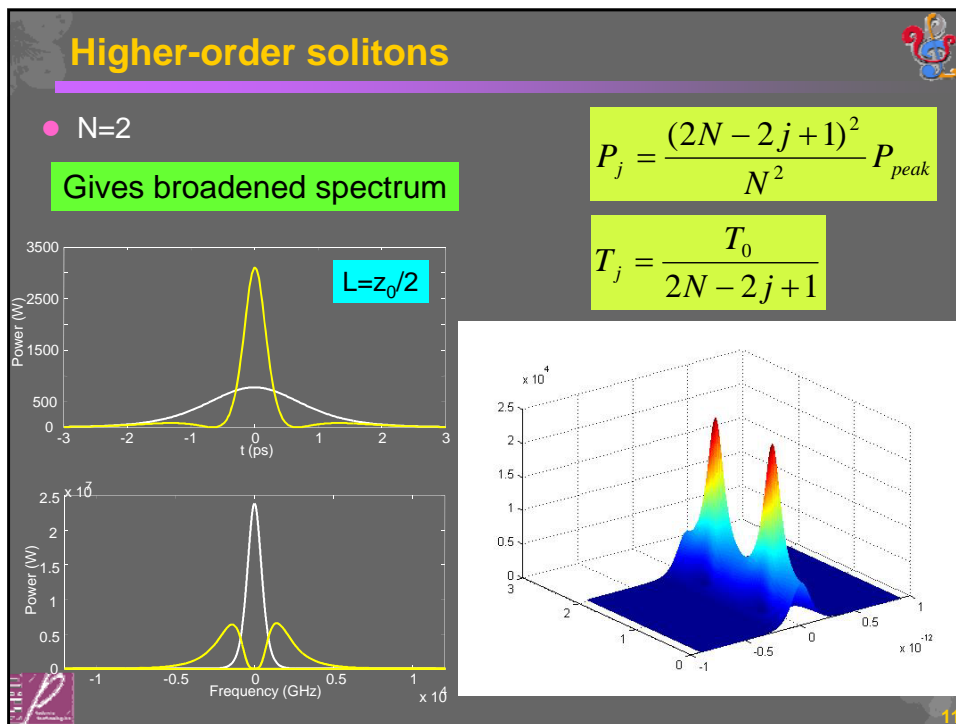
- Dispersion control
 - Energy conservation

$$U = \frac{2|\beta_2|}{\gamma T_0} \Rightarrow \frac{T_0(0)}{T_0(z)} = \frac{\beta_2(0)}{\beta_2(z)}$$

Gives broadened spectrum



10

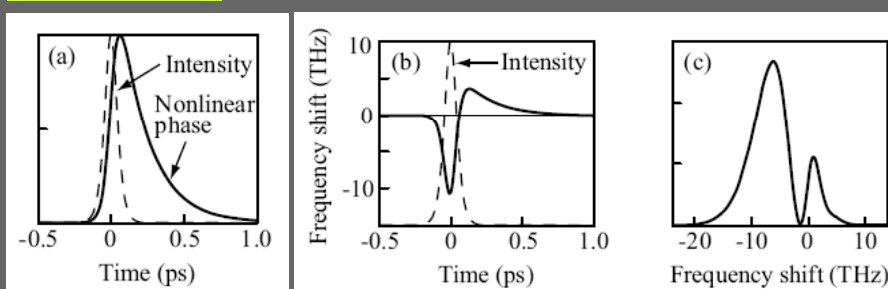


Raman response

- Delayed nonlinear phase → delayed instantaneous frequency
- Red-shift of the pulse power spectrum

Gives broadened spectrum

Time-domain view



A. M. Weiner, *Ultrafast Optics* (Wiley, 2009)

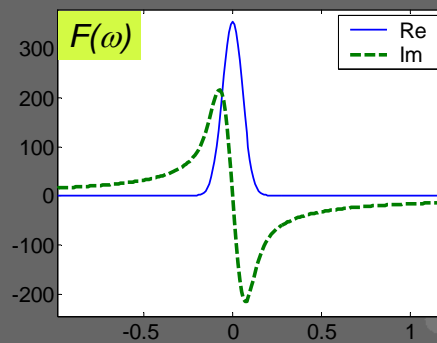
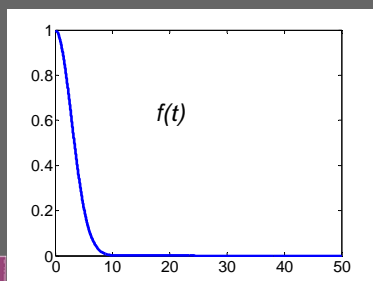
13

Fourier-transform properties

- Raman response real $\Leftrightarrow F(-\omega) = F^*(\omega)$
- A single-sided function
 - $\text{Im}\{F(\omega)\} < 0$ for $\omega > 0$
 - $\text{Im}\{F(\omega)\} > 0$ for $\omega < 0$

Frequency-domain view

$$\text{Im}\{F(\tilde{\omega}_s - \tilde{\omega}_p)\}$$



14

Self-steepening

- Gives rise to **intensity-dependent group velocity**
 - Dispersionless, only electronic response

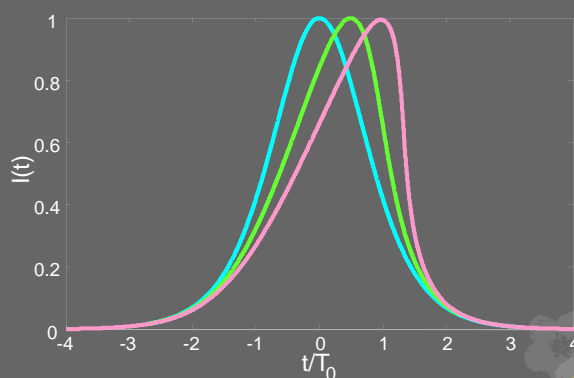
$$\frac{\partial a}{\partial z} + j\gamma \left(1 - \frac{j}{\omega_0} \frac{\partial}{\partial t'} \right) |a|^2 a = 0$$

Gives broadened spectrum

$$\frac{\partial |a|}{\partial z} + \frac{3\gamma}{\omega_0} |a|^2 \frac{\partial |a|}{\partial t'} = 0$$

$$|a| \propto f\left(t' - \frac{z}{v}\right)$$

$$v = \frac{\omega_0}{3\gamma|a|^2}$$

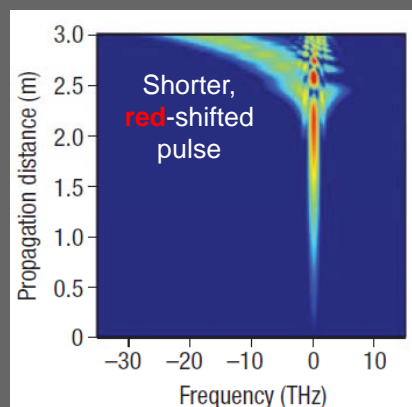
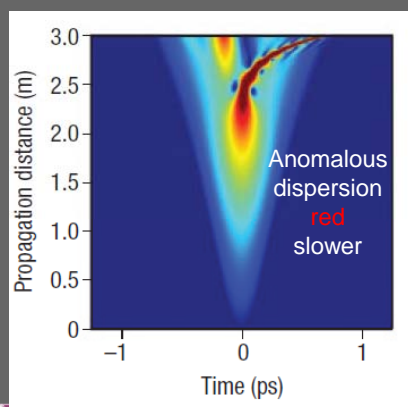


15

Soliton fission


- N^{th} -order solitons will break into N fundamental solitons
 - Raman: soliton self-frequency shift (**SSFS**)
 - Self-steepening

Gives broadened spectrum




Dudley et.al, Nat. Phys. 3, 597 (2007)

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- Ultrafast 3rd-order nonlinear optics
- Supercontinuum generation
- Spectral compression
- Summary

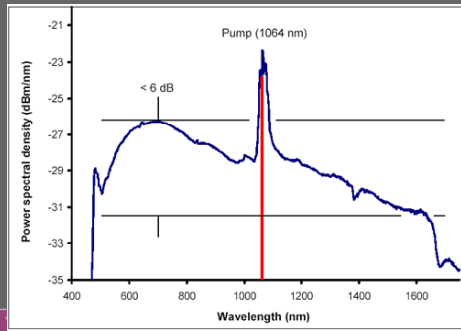


17

SC: Definition



- What is SC?
 - Coherent light source having large bandwidth
 - Broadened input spectrum by nonlinear optical processes
- Applications
 - Metrology, spectroscopy, sensing, ultrashort pulse.....

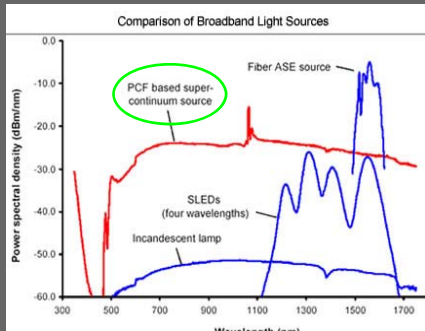


Pump (1064 nm)

< 6 dB

Power spectral density (dBm/nm)

Wavelength (nm)



Comparison of Broadband Light Sources

PCF based supercontinuum source


Fiber ASE source

SLEDs (four wavelengths)

Incandescent lamp

Power spectral density (dBm/nm)

Wavelength (nm)



18

SC: History

- Historical evolution

Year	Nonlinear medium	Laser type	Pulse width	Pulse intensity/ Peak power	SC -20 dB bandwidth
1970	Borosilicate	Nd:Glass	5 ps	1 GW/cm ²	300 nm
1977	Water	YAlG:Nd	30 ps	45 MW	600 nm
1983	Ethylene glycol	Rh6G	80 fs	3 GW	130 nm
1987	MMF	Nd:YAG	25 ps	1.5 GW/cm ²	60 nm
1995	DSF	Er ³⁺ fiber laser	1 ps	1.2 kW	300 nm
1999	MF	Ti:Sapphire	100 fs	8 kW	1200 nm



birth of optical frequency comb tightly linked

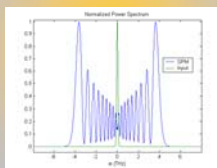
www.bath.ac.uk/physics/groups/ppmg/research_pcf_scg.html

19

SC: Major Nonlinear Processes

SPM

$$\Delta\omega_{inst} = -\gamma L \frac{\partial |A|^2}{\partial t}$$



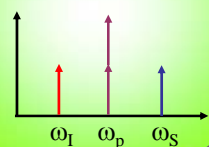
Elastic processes

FWM

$$\omega_1 + \omega_2 = \omega_3 + \omega_4$$

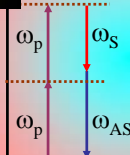
$$2\omega_p = \omega_s + \omega_i$$

$$\Delta k_M + \Delta k_{WG} + \Delta k_{NL} = 0$$

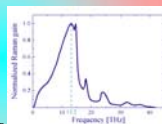


SRS

$$2\omega_p = \omega_s + \omega_{AS}$$



Inelastic Process
Along with SSFS + SS



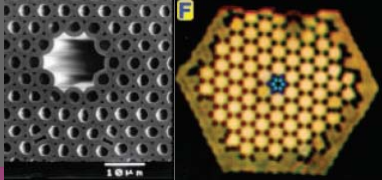
20

Microstructured Fibers

Photonic Crystal Fibers (PCF)

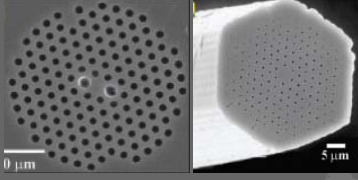
Photonic Band-gap Fibers (PBF)

Guided by photonic band-gap
Confined within air



Microstructured Fibers (MF)

Effective index
Confined within material



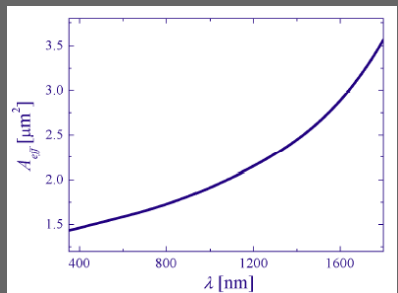
Russel, Science **299**, 358 (2003)

Variable MF Properties

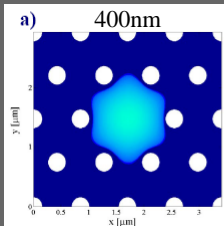
- Small effective modal area
 - Large nonlinearities!
- Endlessly single-mode
 - Cladding index engineering

Lattice pitch, not core diameter!

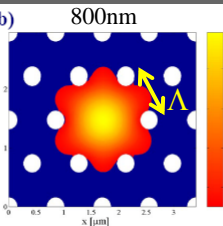
$$V_{eff} = \frac{2\pi\Lambda}{\lambda} \sqrt{n_{co}^2 - n_{cl}^2} < 2.405$$



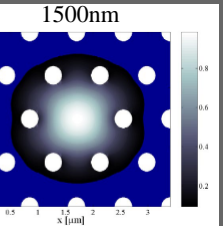
a) 400nm



b) 800nm



c) 1500nm



G. Genty, Ph.D. Dissertation, (Helsinki University of Technology,2004).

Variable MF Properties

- Zero-dispersion wavelength (λ_{ZD})
 - Core size
- Dispersion relation
 - Air-hole diameter
 - Lattice

Reeves et. al., Nature 424, 511 (2003) G. Genty, Ph.D. Dissertation 23

Comparison of MF Properties

Fiber/Properties	SMF	DSF	DCF	HNLF	MF
Attenuation (dB/km)	0.2	0.2	0.45	0.7	80-240
Modal area (μm^2)	85	50	19	12	3
γ (1/W/km)	1.8	2.7	5	15	50

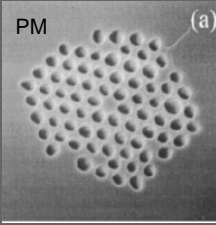
Ability to tune L and NL properties

Much lower power is required to generate SC

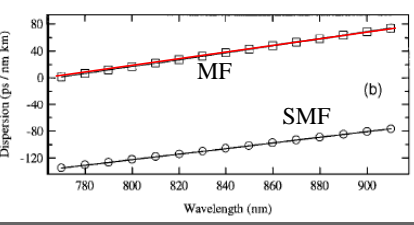
24

First SC Generated using MF

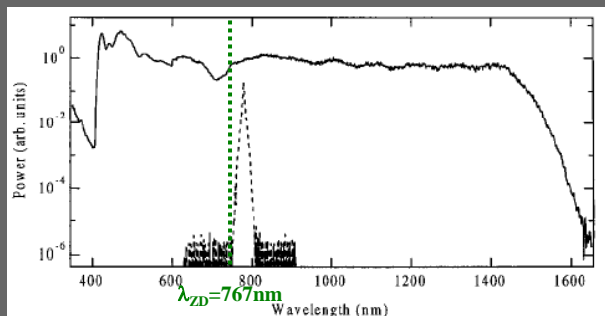
- Ranka, Windeler and Stentz (2000)



(a)



(b)



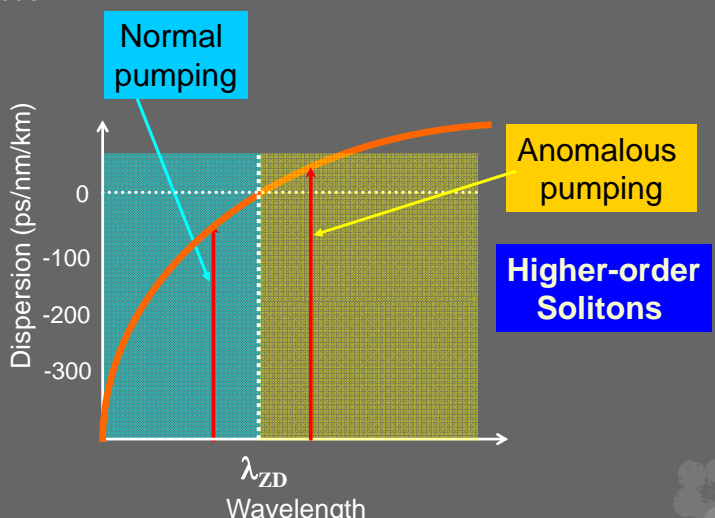
$\lambda_{ZD} = 767 \text{ nm}$

$\lambda_p = 790 \text{ nm}$
 $\lambda_{ZD} = 767 \text{ nm}$
 $\Delta t = 100 \text{ fs}$
 $P_p = 8 \text{ kW}$
 $P_{av} = 64 \text{ mW}$
 $L = 75 \text{ cm}$

Ranka et al., Opt. Lett. 25, 25 (2000)

SC Generation in MF

- Defining pumping modes
 - Anomalous
 - Normal



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Anomalous pumping flow chart

- Higher-order soliton formation
 - Soliton fission
 - Linear radiation wave (RW)
 - Stimulated Raman scattering (SSFS)
 - Later broadening
 - FWM + SRS for anomalous dispersion regime
 - SPM for normal dispersion regime

λ_p

λ_{ZD} wavelength

Higher-order soliton: anomalous pumping

- Effect of pump power
 - Wider, flatter

$$N = \sqrt{\frac{\gamma P_{peak}}{|\beta_2|} \frac{\Delta t}{1.665}}$$

$$T_j = \frac{T_0}{2N - 2j + 1}$$

$\lambda_p = 850 \text{ nm}$
 $\lambda_{ZD} = 806 \text{ nm}$
 $\Delta t = 200 \text{ fs}$
 $L = 6 \text{ m}$

$D [ps^2/(km \cdot nm)]$


$\lambda [\mu m]$

RW


SSFS

$\lambda \text{ (nm)}$

Ortigosa-Blanch et. al., JOSA-B 19, 2567 (2002)




- Ultrafast 3rd-order nonlinear optics
- Supercontinuum generation
- **Spectral compression**
 - Brief history
 - Our approach
 - Numerical and experimental results
- Summary

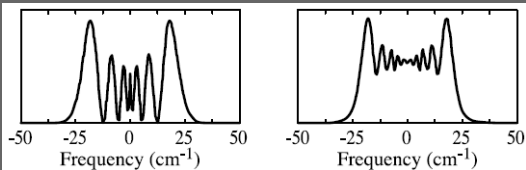


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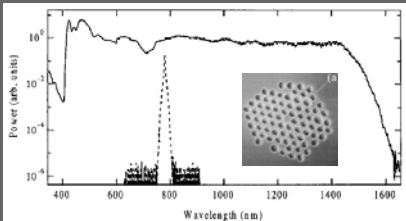
Nonlinear processes → spectral broadening




- Self-phase modulation



A. M. Weiner, *Ultrafast Optics* (Wiley, 2009)
- Supercontinuum generation
 - Higher-order solitons: SPM, dispersion, FWM, SSFS, soliton fission



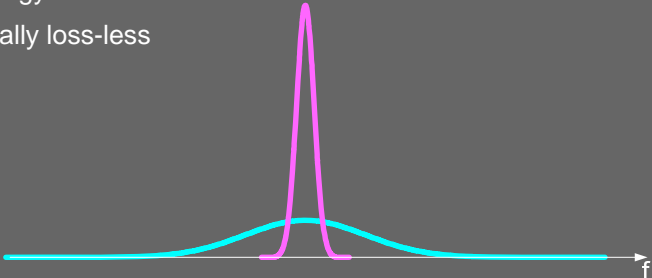
J. K. Ranka et al., *Opt. Lett.* **25**, 25 (2000)



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Applications

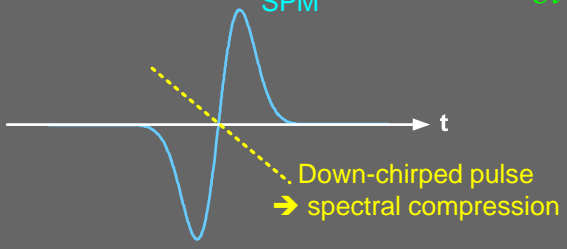
- Highly spectrally bright coherent sources
 - Optical metrology, linear/nonlinear spectroscopy
- ps pulse trains
 - Narrow-band, CARS
- Different compared to spectral filtering
 - Energy re-distribution
 - Ideally loss-less



31

Spectral narrowing/compression

- Earliest experimental observation
 - R. H. Stolen and C. Lin, Phys. Rev. A **17**, 1448 (1978)
- Physical insight explained
 - S. A. Planas et.al, Opt. Lett. **18**, 699 (1993).
 - M. Oberthaler and R. A. Höpfel, Appl. Phys. Lett. **63**, 1017 (1993).
- Instantaneous frequency
 - Pulse chirp, SPM and dispersion

$$\omega_{inst} = \frac{\partial \phi(t)}{\partial t}$$


32

Prior works

- Various fiber types and input chirp conditions
- Compression ratio

Publications	Fiber type (sign of β_2)	Input chirp	Compression ratio
Shen et.al, JLT 17 , 452 (1999)	SMF (-)	+	2
Washburn et.al, OL 25 , 445 (2000)	SMF (+)	-	3.5
Limpert et.al, APB 74 , 191 (2002)	Yb-doped (+)	-	16
Andresen et.al, OL 30 , 2025 (2005)	PCF (-0)	-	21
Sidorov-Biryukov et.al., OE 16 , 2502 (2008)	HNPCF (-)	X	7
Fedotov et.al., OL 34 , 662 (2009)	HNPCF (-)	X	6.5
Nishizawa et.al., OE 18 , 11700 (2010)	CPF (-)	X	25.9
Andresen et.al., OL 36 , 707 (2011)	HNPCF (+)	-	5

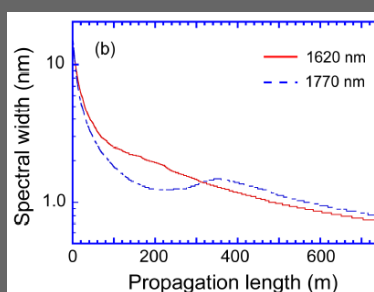
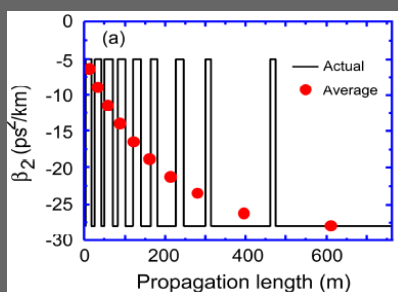
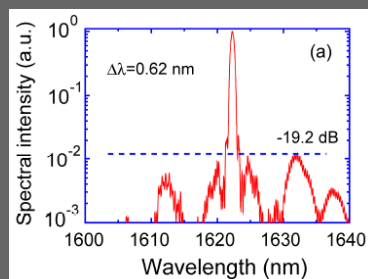
chirp-free, below soliton P_{peak}

solitons

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Soliton effect for spectral compression

- Comb-profile fiber
 - 19 segments of SMF and DSF
 - Large compression ratio of **25.9**
 - Excellent side-lobe suppression



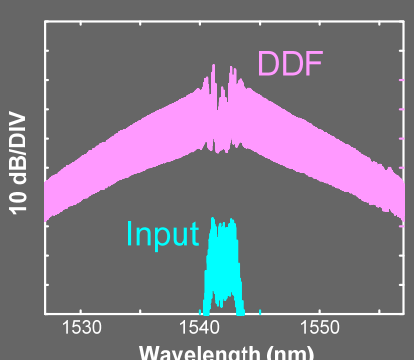
N. Nishizawa et.al., Opt. Express **18**, 11700 (2010)

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Adiabatic soliton temporal compression

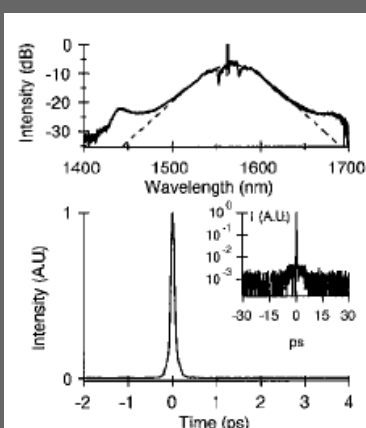
- Dispersion-decreasing fiber
 - Fundamental soliton
 - 60 times pulse compression

$$\Delta\tau_{out} = \frac{D_{out}}{D_{in}} \Delta\tau_{in}$$



10 dB/DIV

Wavelength (nm)




Intensity (dB)


Wavelength (nm)

Intensity (A.U.)

Time (ps)



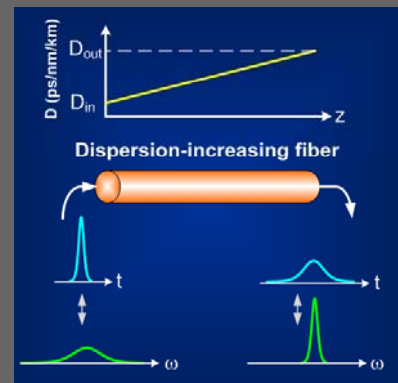
C.-B. Huang et al., Opt. Express 16, 2520 (2008)



K. R. Tamura et al., Opt. Lett. 26, 762 (2001)

Our idea

- Simply run the adiabatic soliton temporal compression in reverse!
- Adiabatic soliton spectral compression
 - Simple idea, never realized!
 - DDF → dispersion-increasing fiber (DIF)
 - Short pulse (wide spectrum) in, broad pulse (narrow spectrum) out



D (ps/nm/km)

D_{out}


D_{in}


z

Dispersion-increasing fiber

t

ω







Split-step Fourier method

- Input pulse envelope $a(0, t)$
- Spectral envelope $A(0, \tilde{\omega})$
- Dispersed spectral envelope $A_D(\Delta z, \tilde{\omega}) = A(0, \tilde{\omega}) \exp[-j \frac{\beta_2 \tilde{\omega}^2 \Delta z}{2}]$
- Dispersed pulse envelope $a_D(\Delta z, t)$
- Dispersed, SPM pulse envelope $a_{D,SPM}(\Delta z, t) = a_D(\Delta z, t) \exp[-j\gamma |a_D(\Delta z, t)|^2 \Delta z]$
- *Now do iterations all the way to the output*

$$\frac{\partial a}{\partial z} = \frac{j\beta_2}{2} \frac{\partial^2 a}{\partial t^2} - j\gamma |a|^2 a$$

$$\downarrow \text{F.T.}$$


$$\downarrow \text{I.F.T.}$$

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Adiabatic soliton spectral compression

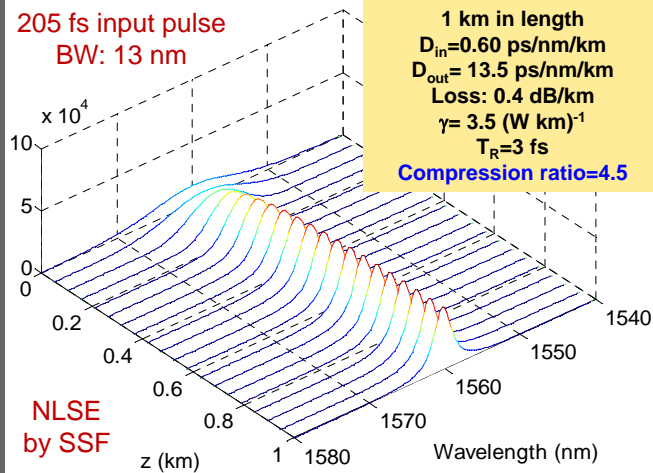
- Dispersion-increasing fiber
 - Ideal compression ratio = $D_{out}/D_{in}=22.5$





Dispersion-increasing fiber

205 fs input pulse
BW: 13 nm

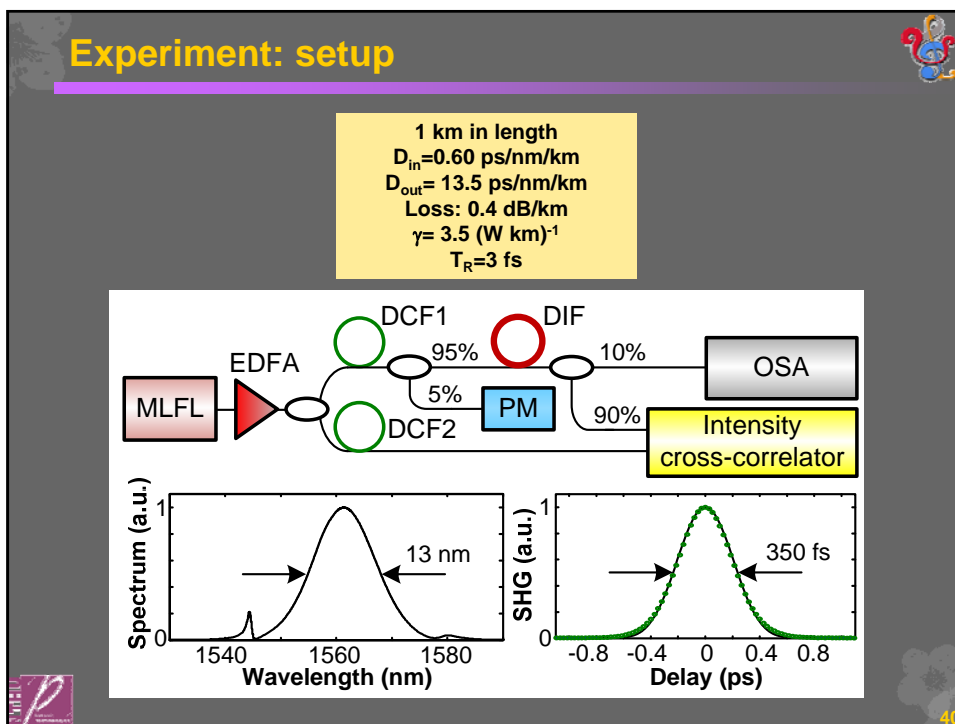
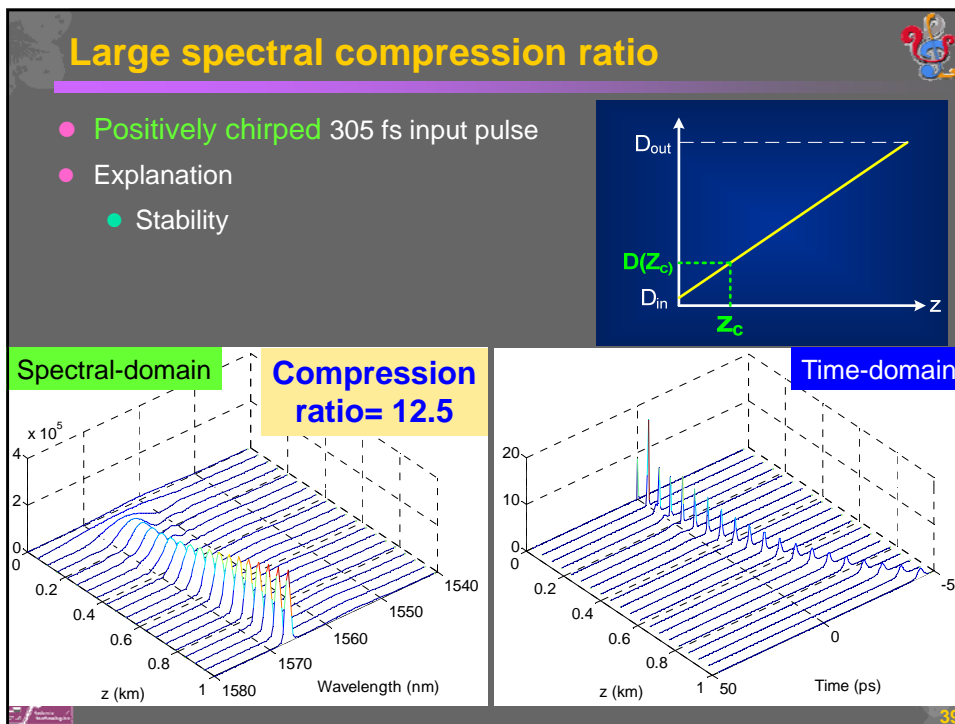
1 km in length
 $D_{in}=0.60 \text{ ps/nm/km}$
 $D_{out}= 13.5 \text{ ps/nm/km}$
Loss: 0.4 dB/km
 $\gamma= 3.5 \text{ (W km)}^{-1}$
 $T_R=3 \text{ fs}$
Compression ratio=4.5

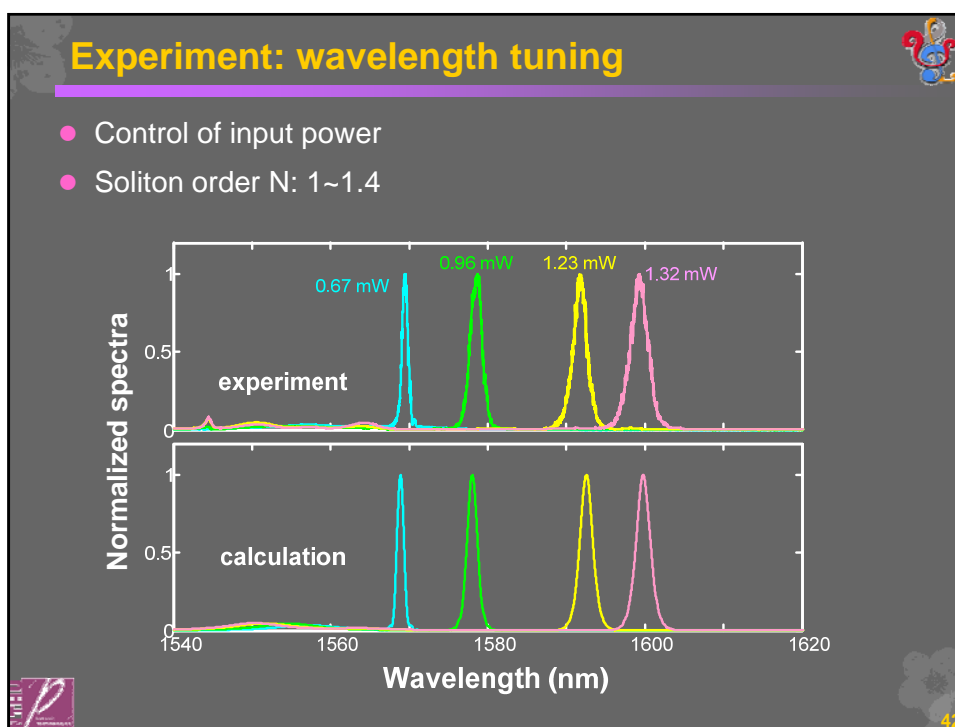
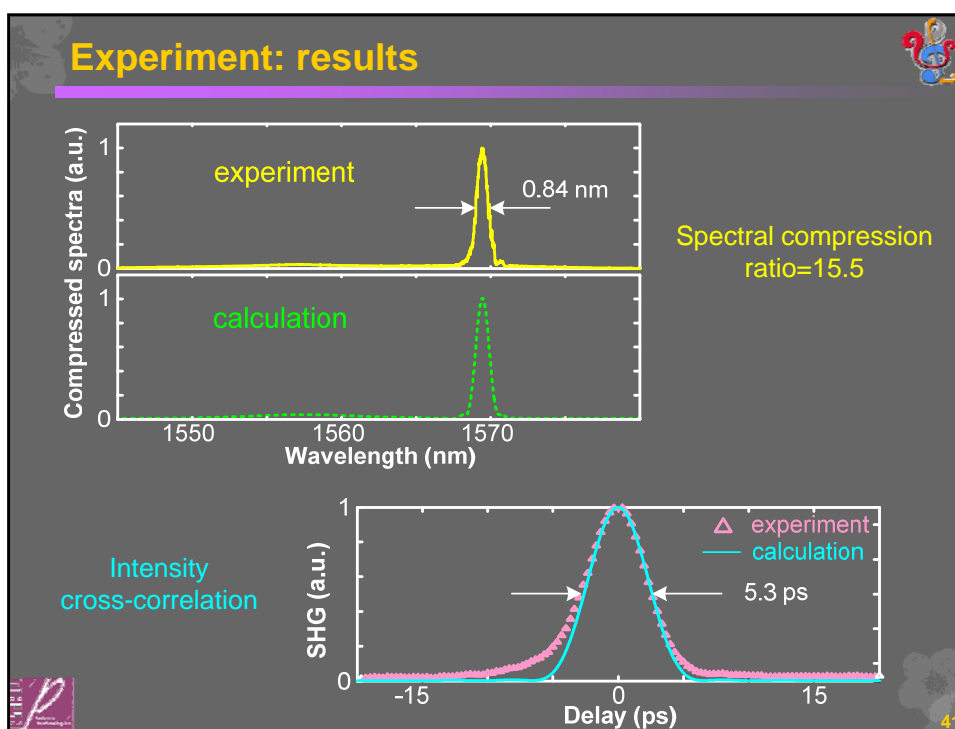


NLSE by SSF

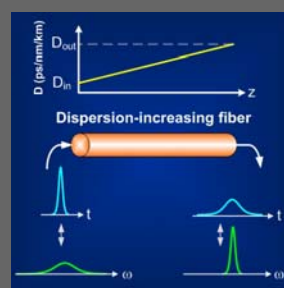
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Summary

- Adiabatic soliton propagation in a **dispersion-increasing fiber** is a simple means in achieving spectral compression
- **Positively chirped** pulse provides better spectral compression ratio if dispersion ramp is too large
- Experiments for **the first time**: 350 fs chirped pulse in a 1-km DIF
 - Spectral compression ratio of 15.5
 - Temporal broadening
 - 30 nm of wavelength tuning ability



H.-P. Chuang and C.-B. Huang, Opt. Lett. **36**, 2848 (2011)

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