Interlayer Exchange Coupling & Boundary Perturbation

張景皓 4/1 2008



Perturbation Wavefunctions:

$$\Phi(x,y) = \Phi_o(x,y) + \sum_{\vec{q}} a_{\vec{q}} e^{-iq_x x + iq_y y}$$
$$\Psi(x,y) = \Psi_o(x,y) + \sum_{\vec{q'}} b_{\vec{q'}} e^{iq'_x x + iq_y y}$$

Expanding order : $|k_x f(y)|$





極限: (震幅/ 波長) > 1/(2Pi)

$$\Phi(x,y) \approx \frac{1}{2} e^{ik_x \frac{d_o}{2}} \Phi_o(x - d_o/2, y) + \frac{1}{2} e^{-ik_x \frac{d_o}{2}} \Phi_o(x + d_o/2, y) + \sum_i a_{q_{y_i}}^{(1)} e^{-iq_{x_i}x + i(k_y \pm \omega_i)y}$$

$$\Psi(x,y) \approx \frac{1}{2} e^{ik_x \frac{d_o}{2}} \Psi_o(x - d_o/2, y) + \frac{1}{2} e^{-ik_x \frac{d_o}{2}} \Psi_o(x + d_o/2, y)$$

$$+\sum_{i} a_{q_{yi}}^{(1)} e^{iq'_{xi}x+i(k_y\pm\omega_i)y}$$



Fig 2: As a planewave be scattered by a seriously rough interface, the process finally could be simplified as a linear combination of two kinds scattering processes.

Interlayer exchange coupling

P. Bruno, PRB,52,00411 (1995)



Above result could be realized as:

$$k_F D \gg 1$$

$$\Delta E \approx \frac{1}{2\pi k_F D} \sum_{n=1}^{\infty} \frac{1}{n} \operatorname{Re}[(R_A(k_F) R_B(k_F))^n e^{i2k_F nD}]$$

n represent the "order" of coupling energy

n=1, similar to RKKY coupling





Magnetic films coupling defined as: $(\Delta E_F - \Delta E_{AF})$

Thickness Fluctuation

Thickness Fluctuations



$$J(D) = \sum_n P(n,D) J(n)$$

Interlayer Exchange Coupling

M. D. Stiles

October 31, 2002

There had been progress toward addressing two issues complicating the comparison between theory and expernent. One complication is the significant disorder that is present in real systems but absent in theoretical models. Fo one form of disorder, namely thickness fluctuations, averaging the interlayer exchange coupling over the growth fron 1 ad been proposed [30] as a solution. Unfortunately, the growth front had not been measured and this correctio 1 ad not been made quantitative. Progress on treating other types of disorder, like interdiffusion, was still to come



Apply our theory on interlayer coupling



$$\Delta E_{qi}^{r} = \frac{1}{4\pi^{3}} Im \int_{-\infty}^{E_{F}} dE \int_{IBZ} d^{2}\vec{k}_{\parallel} \{ -\sum_{n\leq 3} \frac{1}{n} r_{A}^{n} (1 - d_{A}^{2} k_{\perp}^{2})^{n} r_{B}^{n} (1 - d_{B}^{2} k_{\perp}^{2})^{n} e^{2ink_{\perp}D} + \sum_{q_{y}} b_{A}^{(1)} (k_{y} \to q_{y}) b_{B}^{(1)} (q_{y} \to k_{y}) e^{i[k_{\perp}(q_{y}) + k_{\perp}(k_{y})]D} + O[(n + \epsilon) > 3] \}$$

$$(22)$$



Fig 3: Coupling energy by our perturbative method(solid line) and static average(dashed line) . Ry is nature energy unit Rydberg, length scale is Bohr radius. Parameters: $E_F = 0.2Ry/2\pi^3, V_A = V_B = -0.1Ry/2\pi^3, f = 1/4a_o$.







Summary

- We use a new method to handle interface roughness which has physical meaning.
- Interface coupling can be corrected significant.
- Combine lattice structure would be another interesting problem.

References

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- [2] J. Kudrnovský, V. Drchal, et al., Phys. Rev. B. 53, 5125(1996)
- [3] M. D. Stiles, J. Magn. Magn. Mater. 200, 322(1999)