
NNPDF1.0: parton distribution functions with faithful errors

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The NNPDF Collaboration

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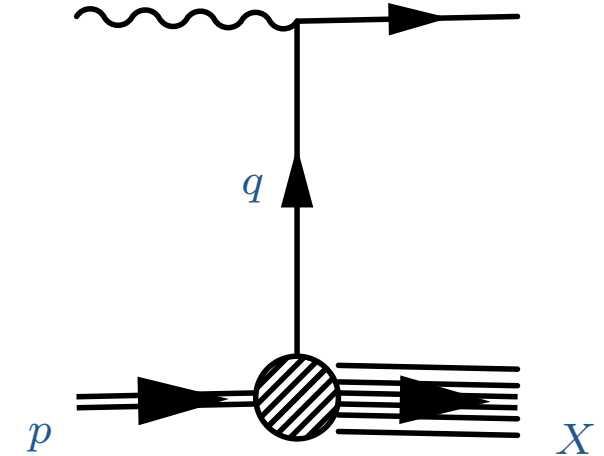
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DIS Parton distribution functions

Deep inelastic observables:

$$F_I(x, Q^2) = \sum_j C_{Ij}(x, \alpha_s(Q^2)) \otimes f_j(x, Q^2)$$

γ^*, W^*, Z^*



Scale-dependence of PDFs:

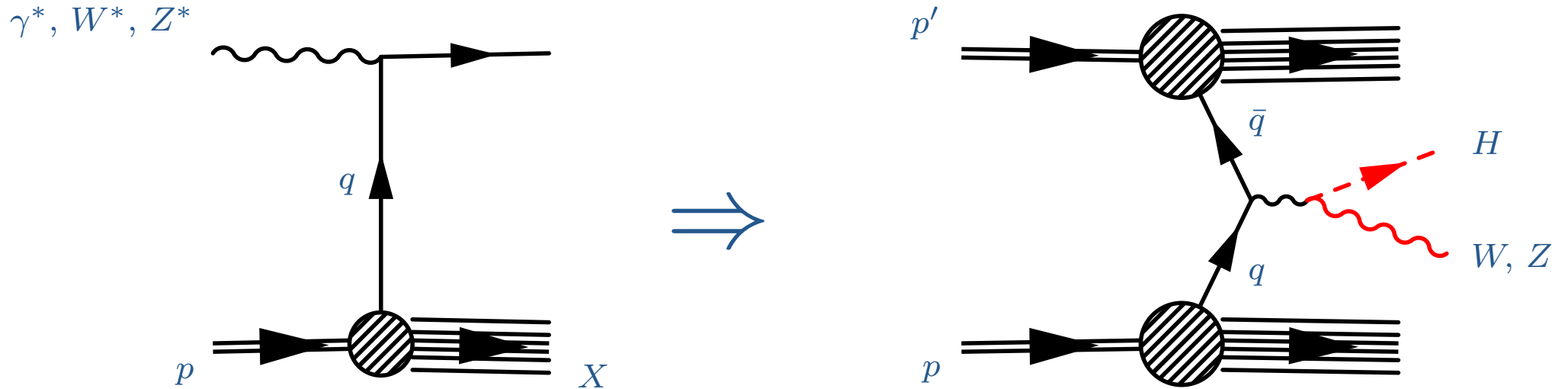
$$Q^2 \frac{\partial}{\partial Q^2} f_i(x, Q^2) = \sum_j P_{ij}(x, \alpha_s(Q^2)) \otimes f_j(x, Q^2)$$

$$f_i(x, Q^2) = \sum_j \Gamma_{ij}(x, \alpha_s, \alpha_s^0) \otimes f_j(x, Q_0^2)$$

Back to the observables:

$$\begin{aligned} F_I(x, Q^2) &= \sum_{jk} C_{Ij}(x, \alpha_s) \otimes \Gamma_{jk}(x, \alpha_s, \alpha_s^0) \otimes f_k(x, Q_0^2) \\ &= \sum_j K_{Ij}(x, \alpha_s, \alpha_s^0) \otimes f_j(x, Q_0^2) \end{aligned}$$

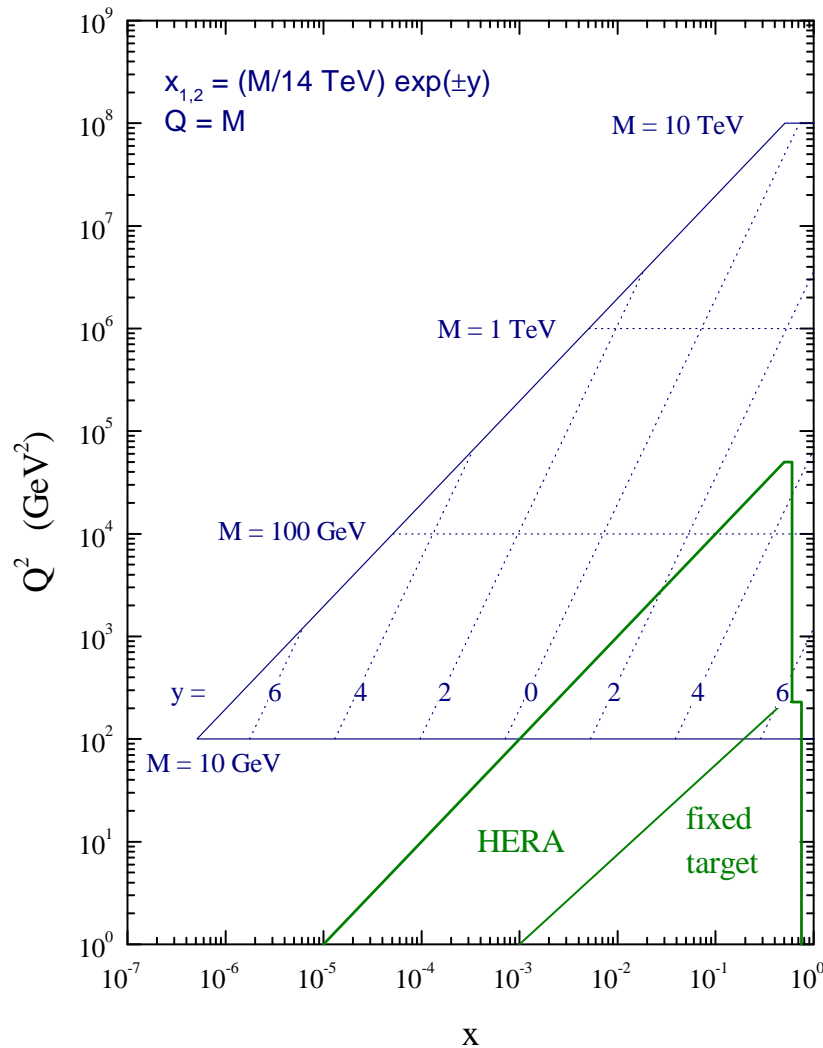
Parton distributions for LHC



- different kinematics
- nonperturbative nucleon structure described by the *same* PDFs
- evolved to the relevant scales

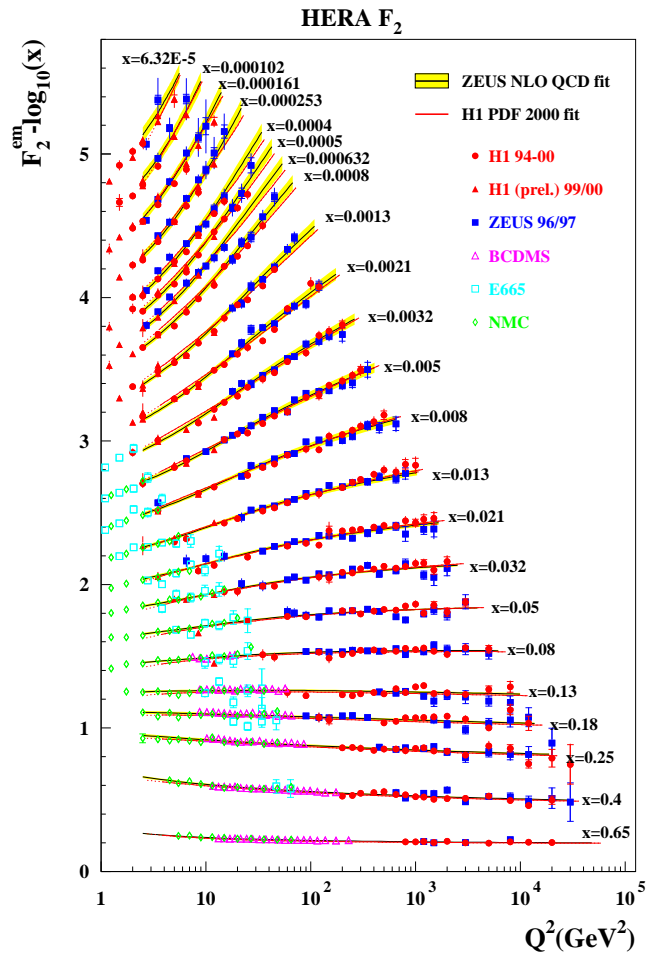
LHC kinematics

LHC parton kinematics



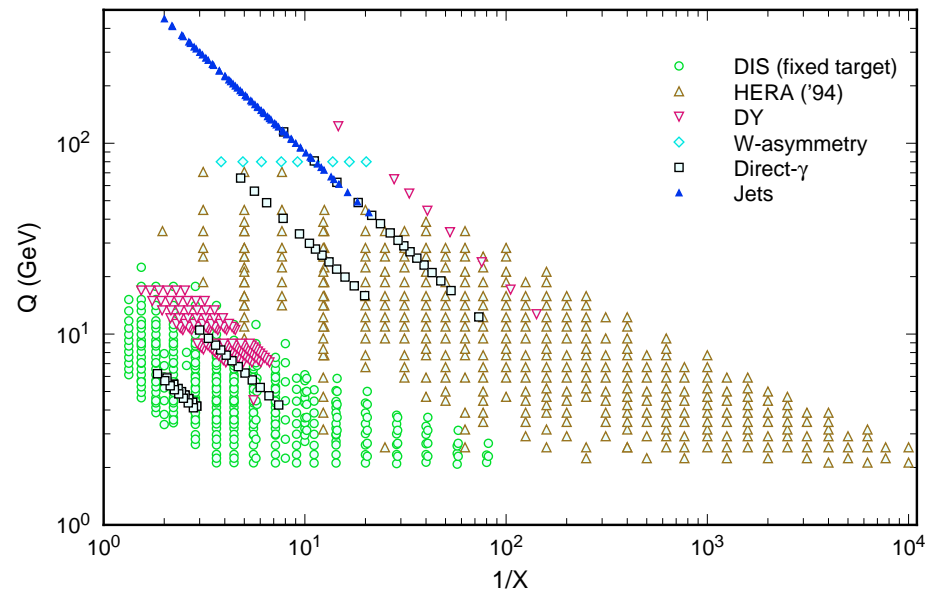
- PDFs need to be extrapolated
- uncertainty in the extrapolation region
- uncertainty propagates into LHC physical observables
- important for modelling the QCD background at LHC
- necessity to develop PDFs with faithful errors

Partons with errors



given a set of **data points**, determine a set of **functions with errors**

data included in CTEQ5 parton fit



What's the problem? [Kosower 99]

- for a single quantity, we quote 1 sigma errors: value \pm error
- for a pair of numbers, we quote a 1 sigma ellipse
- for a function, we need an “error bar” in a space of functions

we must determine the probability density (measure) $\mathcal{P}[f_i(x)]$ in the space of parton distribution functions $f_i(x)$ (i =quark, antiquark, gluon)

EXPECTATION VALUE OF $\mathcal{F}[f_i(x)] \Rightarrow$ **FUNCTIONAL INTEGRAL**

$$\langle \mathcal{F}[f_i(x)] \rangle = \int \mathcal{D}f_i \mathcal{F}[f_i(x)] \mathcal{P}[f_i],$$

we must extract from the data a description of the probability distribution \mathcal{P}

The standard solution

- choose a parameterization at a reference scale
- evolve to desired scale & compute physical observables
- determine best-fit values of parameters
- determine error by propagation of error on parameters ('hessian method') or by parameter scans ('lagrange multiplier method')

problem projected onto the finite-dimensional space of parameters

Comparing global fits

W production cross-section Tevatron

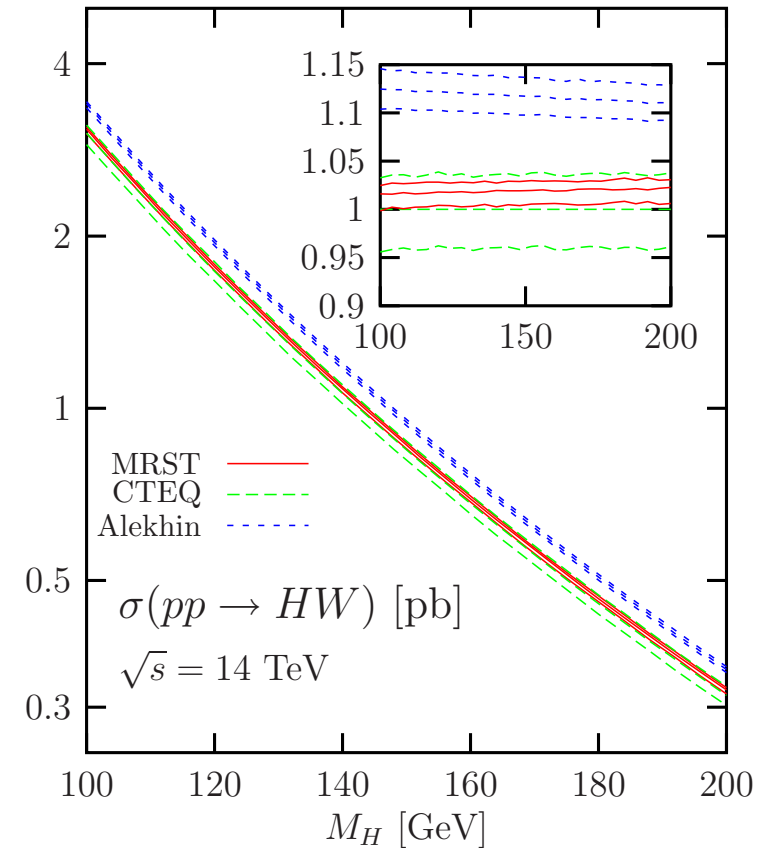
PDF set	xsec [nb]	PDF uncertainty
Alekhin	2.73	± 0.05
MRST2002	2.59	± 0.03
CTEQ6	2.54	± 0.10

[Thorne 03]

Alekhin vs. MRST/CTEQ \rightarrow W production xsect at tevatron do not agree within respective errors

Alekhin vs. MRST/CTEQ \rightarrow predictions for associate Higgs W production LHC do not agree within respective errors

Higgs production at LHC

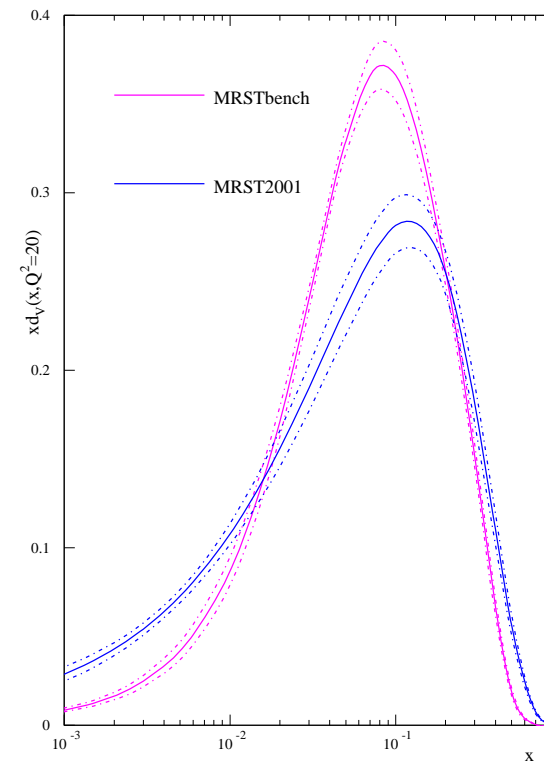
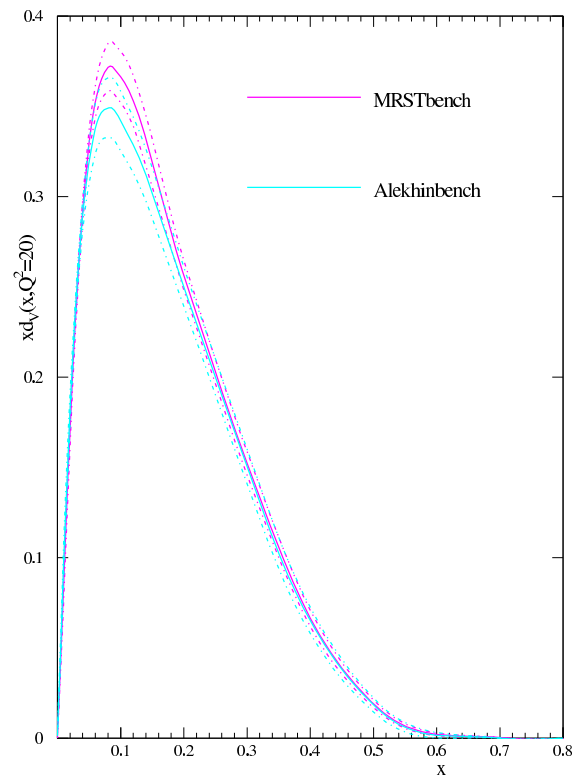


[Djouadi and Ferrag 04]

Troubles with error bars

PDF4LHC workshop at CERN 08

- benchmark fits on reduced sets do not agree with global fits **within errors**
- incompatible experiments?
- lack of generality in the parametrization?
- tolerance criterion $\Delta\chi^2 > 1$?

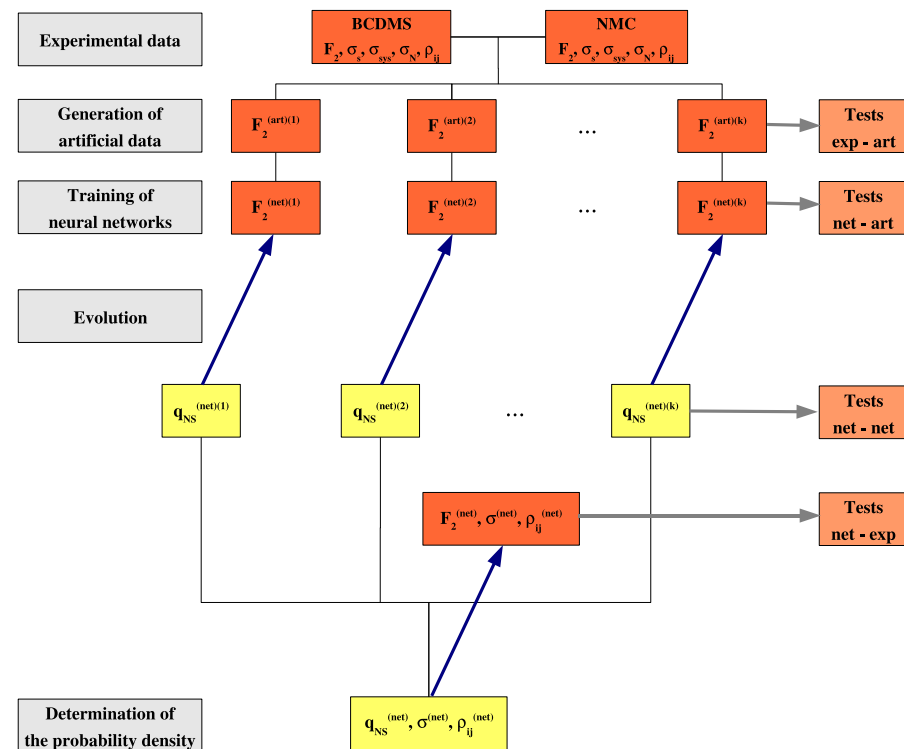


The Neural Monte Carlo

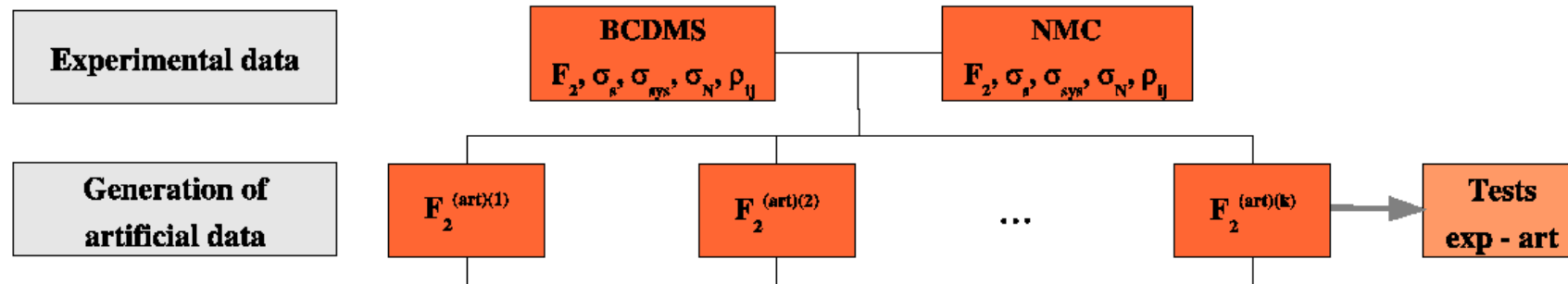
NNPDF collaboration (2004: Idd, Forte, Latorre, Piccione, Rojo; 2007: + Ball, Guffanti, Ubiali)

- Monte Carlo replicas $F_I^{(k)}(p_i)$ of the original dataset $F_I^{(\text{data})}(p_i)$
 \Rightarrow representation of $\mathcal{P}[F_I(p_i)]$ at discrete set of points p_i
- train a neural net for each pdf on each replica, \rightarrow neural representation of the pdfs $f_i^{(\text{net}), (k)}$
- The set of neural nets is a representation of the probability density:

$$\langle \mathcal{F}[f_i] \rangle = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \mathcal{F}[f_i^{(\text{net})}(k)]$$



MC sampling of exp data



$$F_{I,p}^{(art)(k)} = S_{p,N}^{(k)} F_{I,p}^{(exp)} \left(1 + \sum_{l=1}^{N_c} r_{p,l}^{(k)} \sigma_{p,l} + r_p^{(k)} \sigma_{p,s} \right), \quad k = 1, \dots, N_{rep},$$

where

$$S_{p,N}^{(k)} = \prod_{n=1}^{N_a} \left(1 + r_{p,n}^{(k)} \sigma_{p,n} \right) \prod_{n=1}^{N_r} \sqrt{1 + r_{p,n}^{(k)} \sigma_{p,n}}.$$

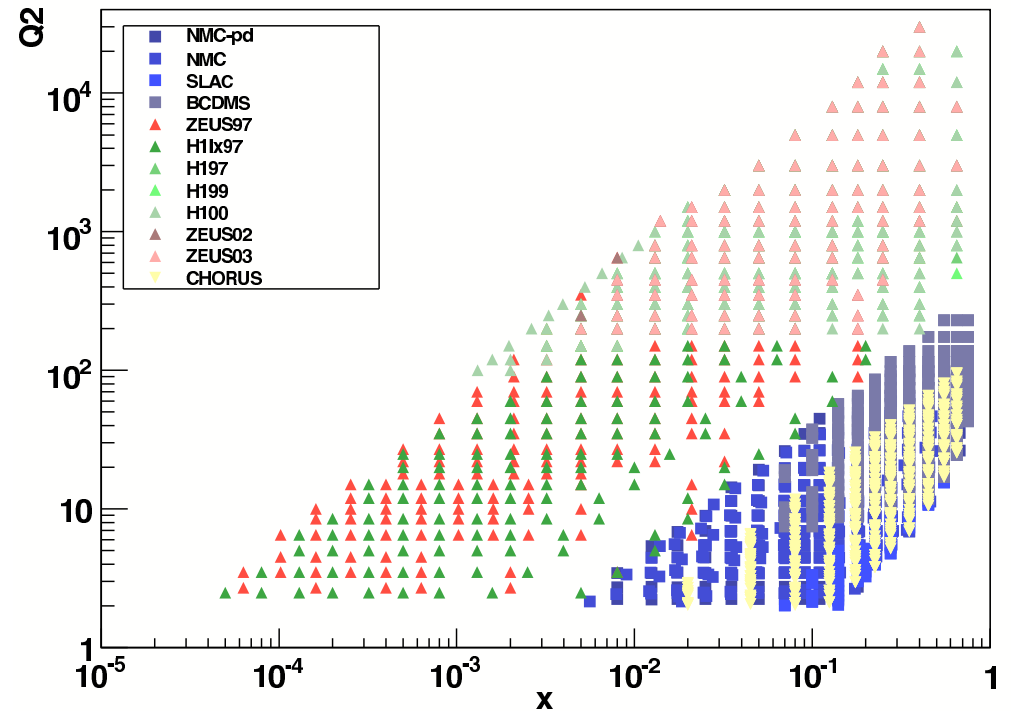
Monte Carlo errors

- for each replica (k) we fit one set of PDFs
- the ensemble of fitted replicas represents the probability distribution in the space of PDFs
- statistical properties of any function of the PDFs can be computed using standard methods:

$$\begin{aligned}\langle \mathcal{F}[f(x)] \rangle &= \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \mathcal{F}[f^{(k)(\text{net})}(x)] \\ \sigma_{\mathcal{F}[f(x)]} &= \sqrt{\langle \mathcal{F}[f(x)]^2 \rangle - \langle \mathcal{F}[f(x)] \rangle^2} \\ \rho[f_a(x_1, Q_1^2), f_b(x_2, Q_2^2)] &= \frac{\langle f_a(x_1, Q_1^2) f_b(x_2, Q_2^2) \rangle - \langle f_a(x_1, Q_1^2) \rangle \langle f_b(x_2, Q_2^2) \rangle}{\sigma_a(x_1, Q_1^2) \sigma_b(x_2, Q_2^2)}\end{aligned}$$

Experimental data

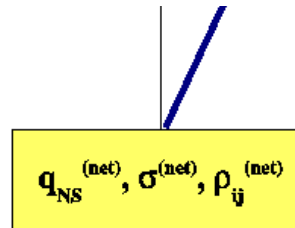
OBS	Data set	OBS	Data set
F_2^p	NMC	σ_{NC}^-	ZEUS
	SLAC		H1
	BCDMS	σ_{CC}^+	ZEUS
F_2^d	SLAC		H1
	BCDMS	σ_{CC}^-	ZEUS
σ_{NC}^+	ZEUS		H1
	H1	$\sigma_\nu, \sigma_{\bar{\nu}}$	CHORUS
F_2^d / F_2^p	NMC-pd	F_L	H1



- Kinematical cuts:
 $Q^2 > 2 \text{ GeV}^2$
 $W^2 = Q^2(1 - x)/x > 12.5 \text{ GeV}^2$
- ~ 3000 points.

Neural Networks?

**Determination of
the probability density**



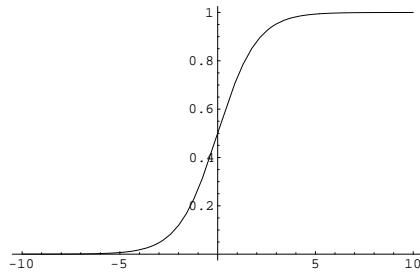
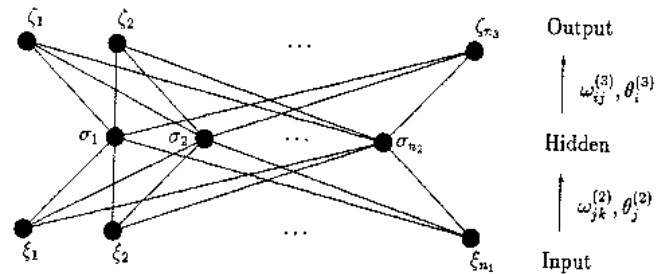
each PDF at the reference scale is parametrised by one NN:

$$f_i(x, Q_0^2) = N_i(x)$$

NN is a non-linear mapping:

$$N_i : \mathbb{R}^n \rightarrow \mathbb{R}^m$$

Some details about Neural Networks



Multilayer feed-forward networks

- Each neuron receives input from neurons in preceding layer and feeds output to neurons in subsequent layer
- Activation determined by weights and thresholds

$$\xi_i = g \left(\sum_j \omega_{ij} \xi_j - \theta_i \right)$$

- Sigmoid activation function

$$g(x) = \frac{1}{1 + e^{-\beta x}}$$

... just another set of basis functions!

eg, a 1-2-1 NN: $\xi_1^{(3)}(\xi_1^{(1)}) = \frac{1}{1 + \exp \left[\theta_1^{(3)} - \frac{\omega_{11}^{(2)}}{1 + e^{\theta_1^{(2)} - \xi_1^{(1)} \omega_{11}^{(1)}}} - \frac{\omega_{12}^{(2)}}{1 + e^{\theta_2^{(2)} - \xi_1^{(1)} \omega_{21}^{(1)}}} \right]}$

Thm: any function can be represented by a sufficiently big neural network

Basis set

- Each independent PDF at the initial scale $Q_0^2 = 2\text{GeV}^2$ is parameterized by an individual NN.
- Little constraint on strange \rightarrow Flavor Assumptions:
 - Symmetric strange sea $s(x) = \bar{s}(x)$
 - Strange sea proportional to non-strange sea $\bar{s}(x) = \frac{C}{2}(\bar{u}(x) + \bar{d}(x))$ (C = 0.5)
 - Intrinsic heavy quarks contributions neglected.
- Parametrization of **(4+1)** combinations of PDFs at $Q_0^2 = 2 \text{ GeV}^2$:

Singlet : $\Sigma(x)$	$\mapsto \text{NN}_\Sigma(x)$	2-5-3-1 37 pars
Gluon : $g(x)$	$\mapsto \text{NN}_g(x)$	2-5-3-1 37 pars
Total valence : $V(x) \equiv u_V(x) + d_V(x)$	$\mapsto \text{NN}_V(x)$	2-5-3-1 37 pars
Non-singlet triplet : $T_3(x)$	$\mapsto \text{NN}_{T_3}(x)$	2-5-3-1 37 pars
Sea asymmetry : $\Delta_S(x) \equiv \bar{d}(x) - \bar{u}(x)$	$\mapsto \text{NN}_\Delta(x)$	2-5-3-1 37 pars

185 parameters

Normalization and sum rules

$$\Sigma(x, Q_0^2) = (1-x)^{m_\Sigma} x^{-n_\Sigma} \text{NN}_\Sigma(x),$$

$$V(x, Q_0^2) = A_V (1-x)^{m_V} x^{-n_V} \text{NN}_V(x),$$

$$T_3(x, Q_0^2) = (1-x)^{m_{T_3}} x^{-n_{T_3}} \text{NN}_{T_3}(x),$$

$$\Delta_S(x, Q_0^2) = A_{\Delta_S} (1-x)^{m_{\Delta_S}} x^{-n_{\Delta_S}} \text{NN}_{\Delta_S}(x),$$

$$g(x, Q_0^2) = A_g (1-x)^{m_g} x^{-n_g} \text{NN}_g(x).$$

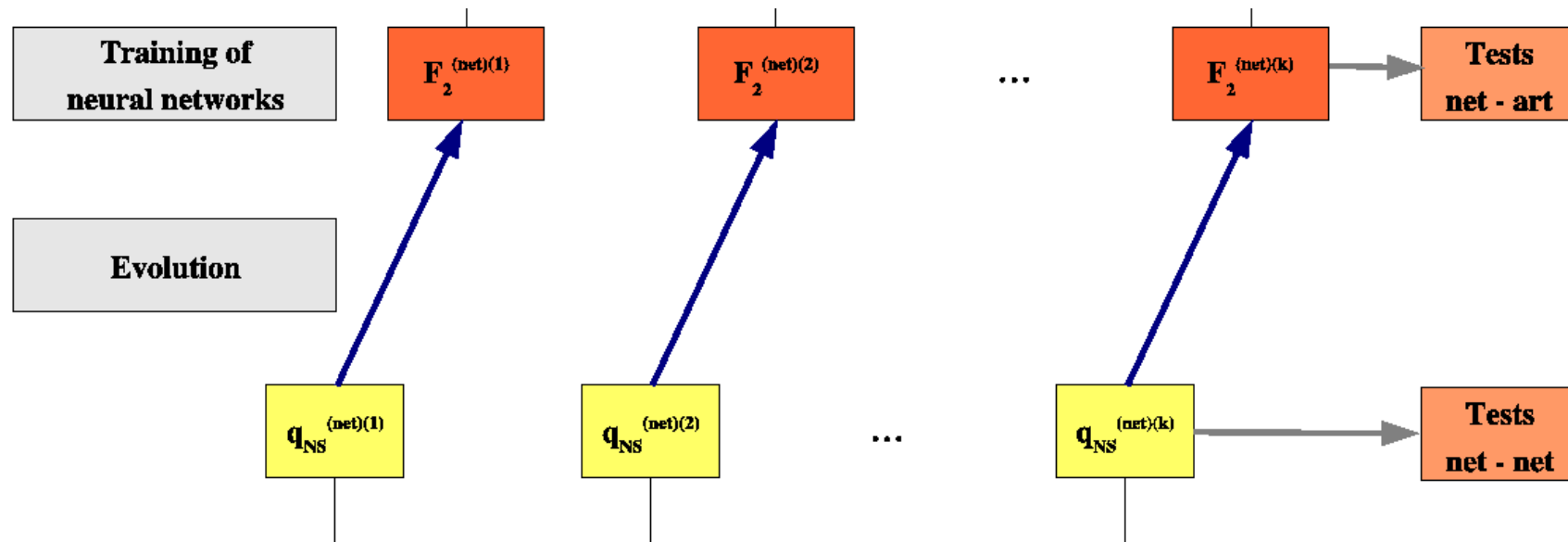
- **Polynomial Preprocessing** → Training Efficiency
- **Normalization** → Fixed by valence and momentum sum rules

$$\int_0^1 dx x (\Sigma(x) + g(x)) = 1$$

$$\int_0^1 dx (u(x) - \bar{u}(x)) = 2$$

$$\int_0^1 dx (d(x) - \bar{d}(x)) = 1.$$

Evolution



Observables are a convolution over x of PDFs and Coefficient Functions:

$$F_I(x, Q^2) = \sum_j C_{Ij}(x, \alpha_s) \otimes f_j(x, Q^2) = \sum_{j,k} C_{Ij}(x, \alpha_s) \otimes \Gamma_{jk}(x, \alpha_s, \alpha_s^0) \otimes f_k(x, Q_0^2)$$

Kernels for a physical observable

F_2 proton structure function

$$F_2^p = x \left\{ \frac{5}{18} C_{2,q}^s \otimes \Sigma + \frac{1}{6} C_{2,q} \otimes (T_3 + \frac{1}{3}(T_8 - T_{15})) + \frac{1}{5}(T_{24} - T_{35}) \right. \\ \left. + \langle e_q^2 \rangle C_{2,g} \otimes g \right\}$$

$$F_2^p = x \left\{ K_{F2,\Sigma} \otimes \Sigma_0 + K_{F2,g} \otimes g_0 + K_{F2,+} \otimes \left(T_{3,0} + \frac{1}{3}(T_{8,0} - T_{15,0}) \right) \right\}$$

In Mellin space

$$K_{F2,\Sigma} = \frac{5}{18} C_{2,q}^s \Gamma_S^{qq} + \frac{1}{30} C_{2,q} (\Gamma_S^{24,q} - \Gamma_S^{35,q}) + \langle e_q^2 \rangle C_{2,g} \Gamma_S^{gq}$$

$$K_{F2,g} = \frac{5}{18} C_{2,q}^s \Gamma_S^{qg} + \frac{1}{30} C_{2,q} (\Gamma_S^{24,g} - \Gamma_S^{35,g}) + \langle e_q^2 \rangle C_{2,g} \Gamma_S^{gg}$$

$$K_{F2,+} = \frac{1}{6} C_{2,q} \Gamma_{NS}^+$$

Theoretical errors

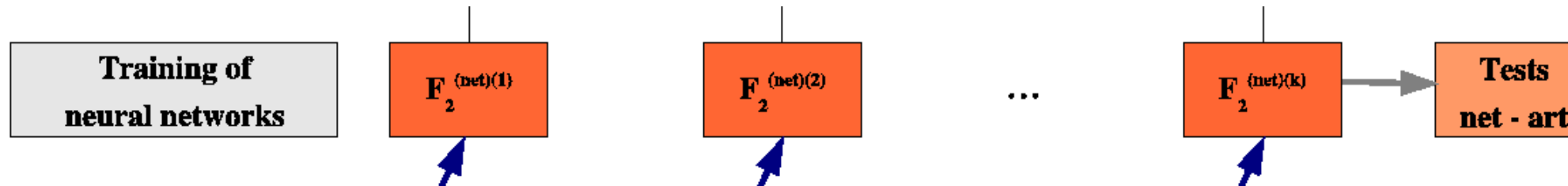
- Higher perturbative orders → **NLO** fit
- Heavy quark treatment → **Zero Mass Variable Flavor Number** scheme.
quarks are radiatively generated at thresholds. [Thorne,Tung, arXiv:0809.0714]
- **Target Mass Corrections** included and factorized into the hard kernels.

$$\tilde{F}_2(\xi, Q^2) = \frac{x^2}{\tau^{3/2}} \frac{F_2(\xi, Q^2)}{\xi^2} + 6 \frac{M_N^2}{Q^2} \frac{x^3}{\tau^2} I_2(\xi, Q^2)$$
$$\tau = 1 + \frac{4M_N^2 x^2}{Q^2}$$
$$\xi = \frac{2x}{1 + \sqrt{\tau}}$$
$$I_2(\xi, Q^2) = \int_{\xi}^1 \frac{dz}{z^2} F_2(z, Q^2).$$

Taking Mellin transforms with respect to ξ , defines a new target mass corrected coefficient function

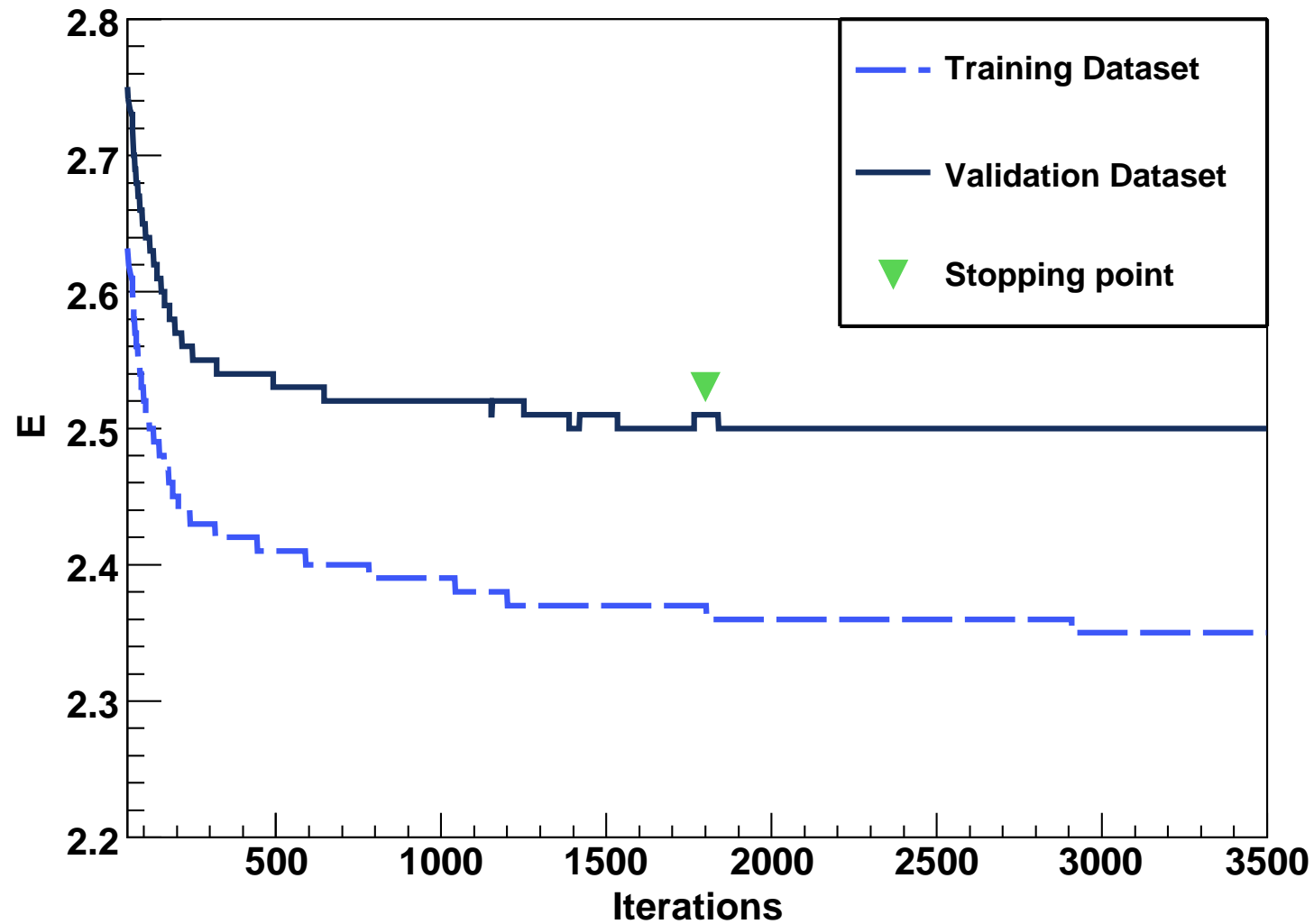
$$\tilde{C}_{2,j}(N, \alpha_s, \tau) = \frac{(1 + \tau^{1/2})^2}{4\tau^{3/2}} \left(1 + \frac{3(1 - \tau^{-1/2})}{N + 1} \right) C_{2,j}(N, \alpha_s)$$

Training strategy

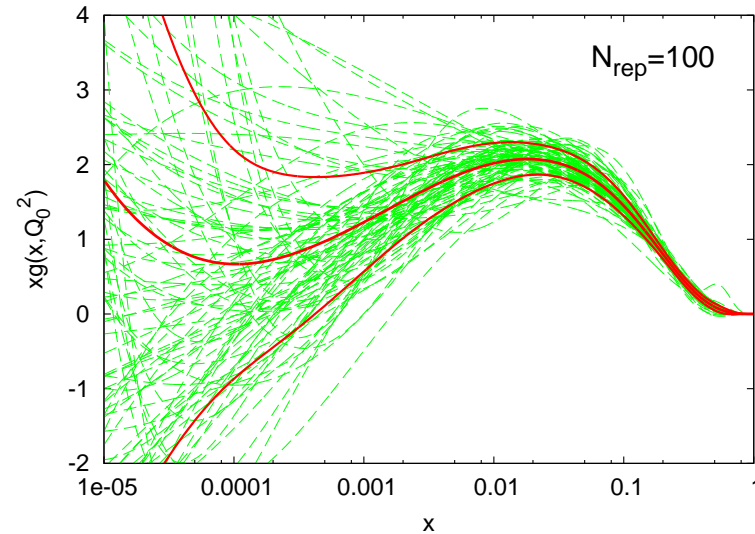
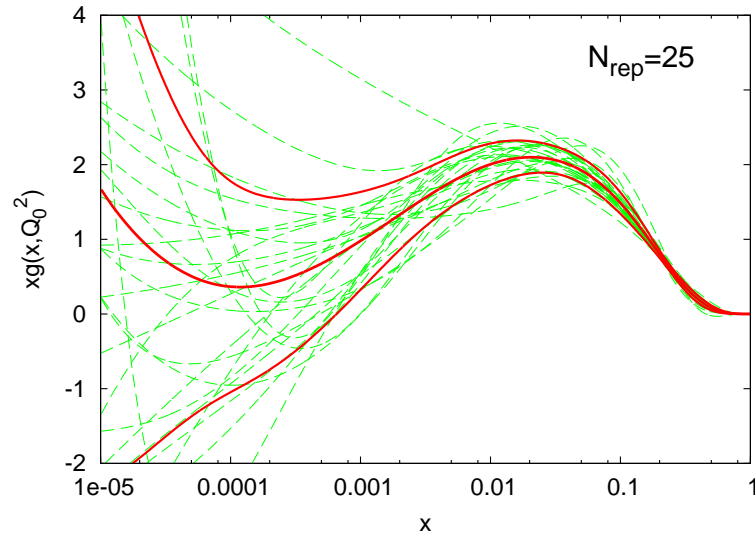


- unbiased basis of functions, parametrized by a large number of parameters
- genetic algorithms for minimization
- might accommodate statistical fluctuations of the data
- optimal training, beyond which the fit is just adjusting to statistical fluctuations
- dynamical stopping by cross validation
- for each replica divide the data randomly into **training** and **validation**
- minimization performed on the training set **only**
- when the training χ^2 still decreases while the validation χ^2 stops decreasing
→ **STOP**

Dynamical stopping



Ensemble of replicas



- individual replicas fluctuate significantly
- averages are smooth as the number of replicas is increased

PDFs uncertainties

- Monte Carlo prescription (**NNPDF**)

$$\sigma_{\mathcal{F}} = \left(\frac{N_{\text{set}}}{N_{\text{set}} - 1} (\langle \mathcal{F}[\{f\}]^2 \rangle - \langle \mathcal{F}[\{f\}] \rangle^2) \right)^{1/2}$$

- HEPDATA prescription (**CTEQ** and **MRST/MSTW**)

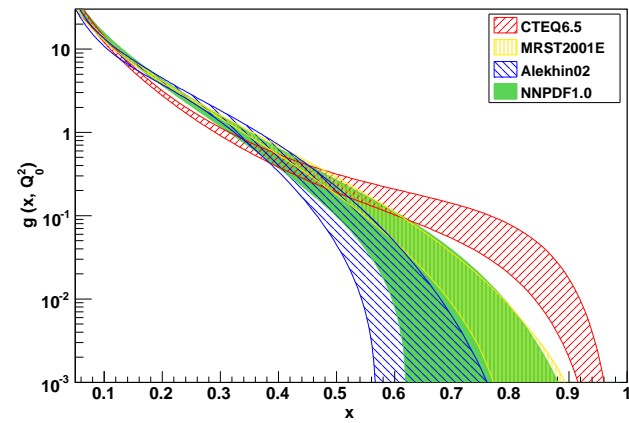
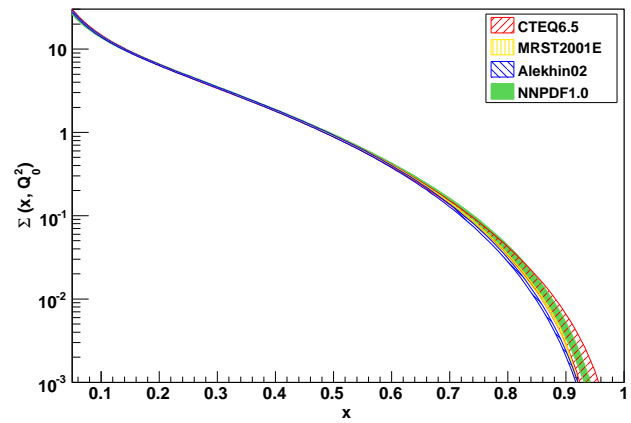
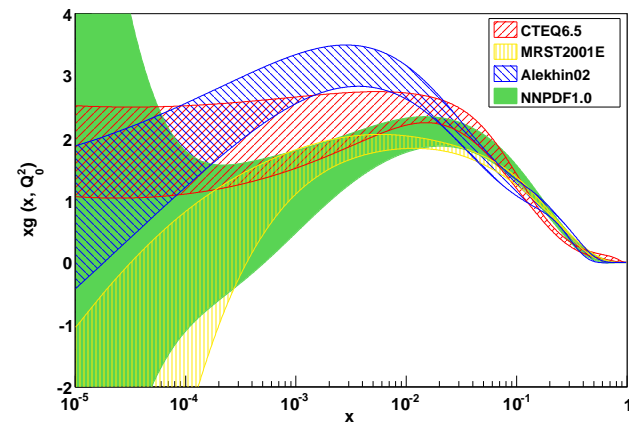
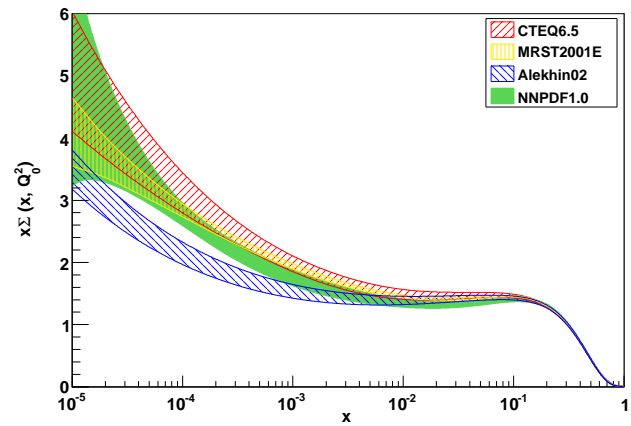
$$\sigma_{\mathcal{F}} = \frac{1}{2C_{90}} \left(\sum_{k=1}^{N_{\text{set}}/2} \left(\mathcal{F}[\{f^{(2k-1)}\}] - \mathcal{F}[\{f^{(2k)}\}] \right)^2 \right)^{1/2}, \quad C_{90} = 1.64485$$

C_{90} accounts for the fact that the upper and lower parton sets correspond to 90% confidence levels rather than to one- σ uncertainties.

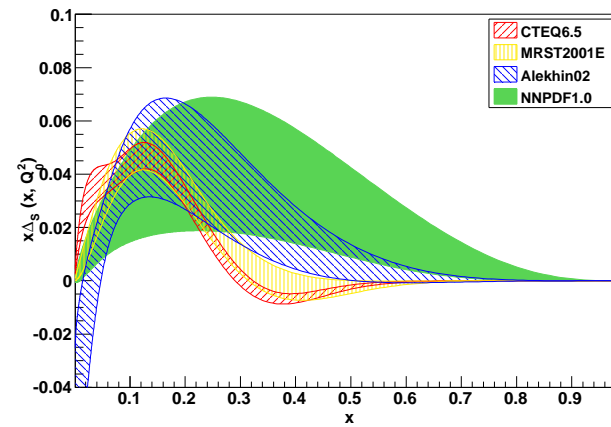
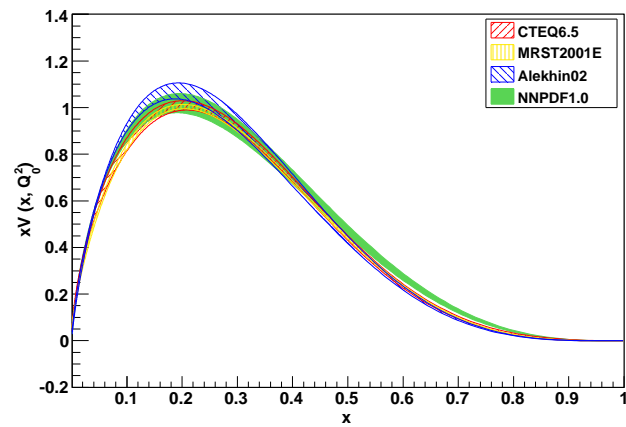
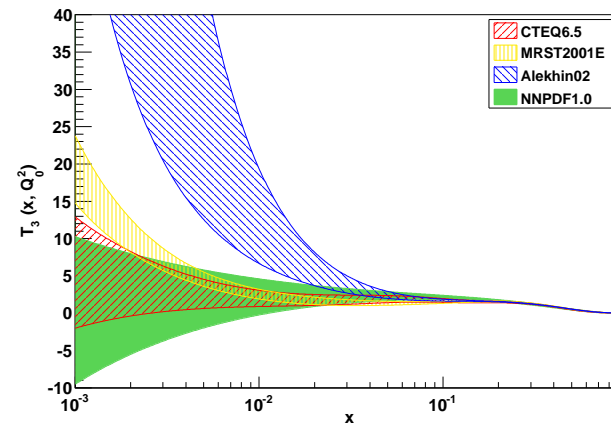
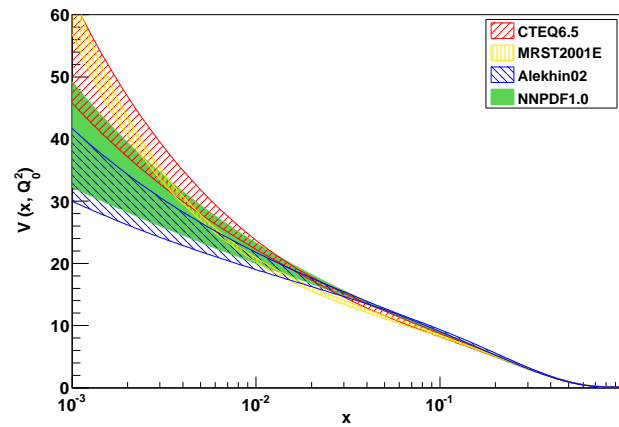
- HEPDATA* prescription (**Alekhin**)

$$\sigma_{\mathcal{F}} = \left(\sum_{k=1}^{N_{\text{set}}} \left(\mathcal{F}[\{f^{(k)}\}] - \mathcal{F}[\{f^{(0)}\}] \right)^2 \right)^{1/2}.$$

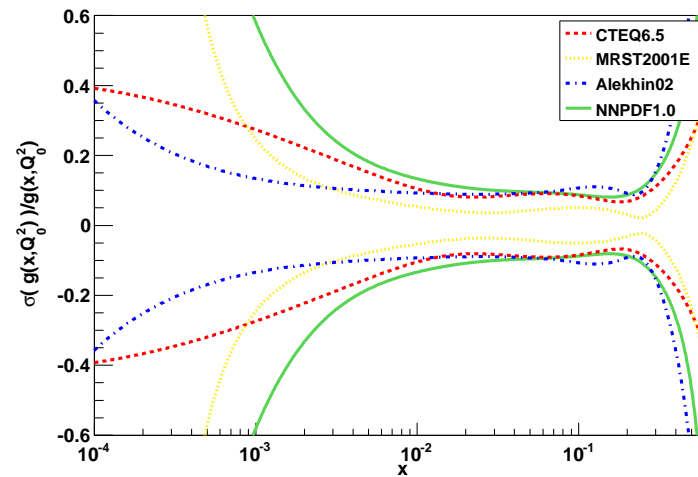
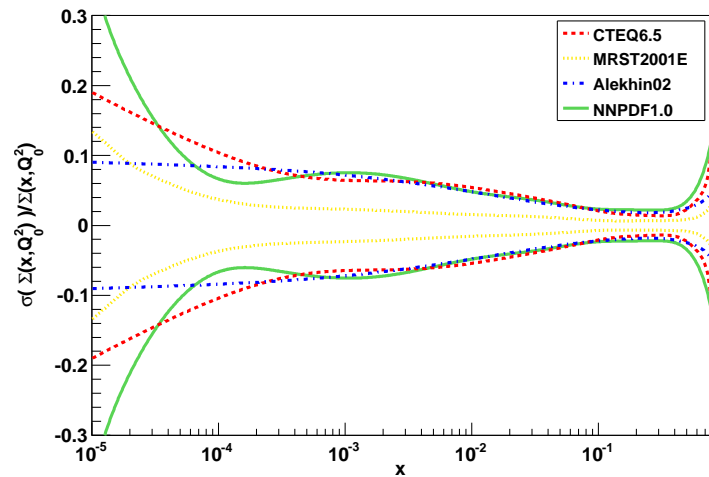
The NNPDF1.0 parton set



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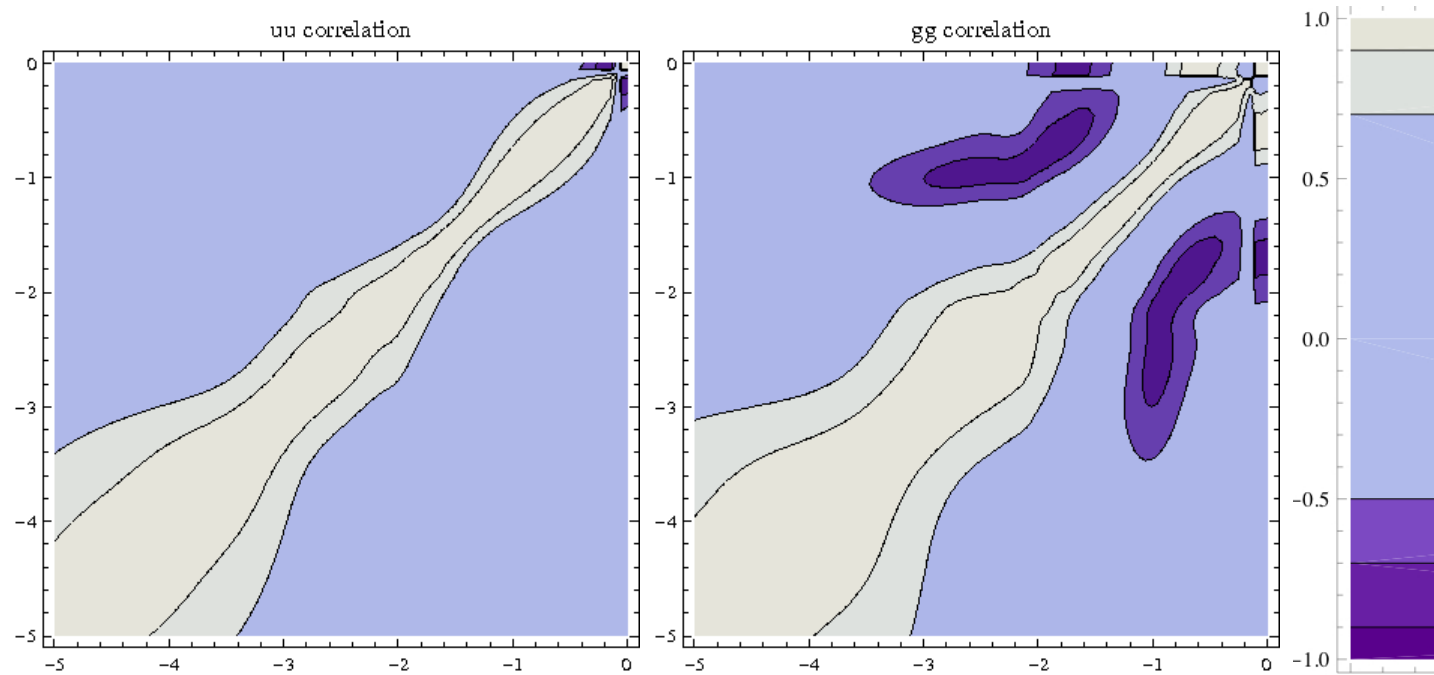
Relative uncertainties



- comparable uncertainty in the data region
- BUT larger uncertainties in the extrapolation

PDFs correlations

Correlations between $u - u$ and $g - g$ ($Q=85\text{GeV}$) [nadolsky 08]



$$\rho [f_a(x_1, Q_1^2) f_b(x_2, Q_2^2)] = \frac{\langle f_a(x_1, Q_1^2) f_b(x_2, Q_2^2) \rangle_{\text{rep}} - \langle f_a(x_1, Q_1^2) \rangle_{\text{rep}} \langle f_b(x_2, Q_2^2) \rangle_{\text{rep}}}{\sigma_a(x_1, Q_1^2) \sigma_b(x_2, Q_2^2)}$$

Distance between MC ensembles.

- Stability of the NNPDF parton set can be assessed by using standard statistical tools.
- Distances between two probability distributions: $\{f_{ik}^{(1)} = f_k^{(1)}(x_i, Q_0^2)\}$

$$\langle d[f] \rangle = \sqrt{\left\langle \frac{\left(\langle f_i \rangle_{(1)} - \langle f_i \rangle_{(2)}\right)^2}{\sigma^2[f_i^{(1)}] + \sigma^2[f_i^{(2)}]} \right\rangle_{\text{pts}}}$$

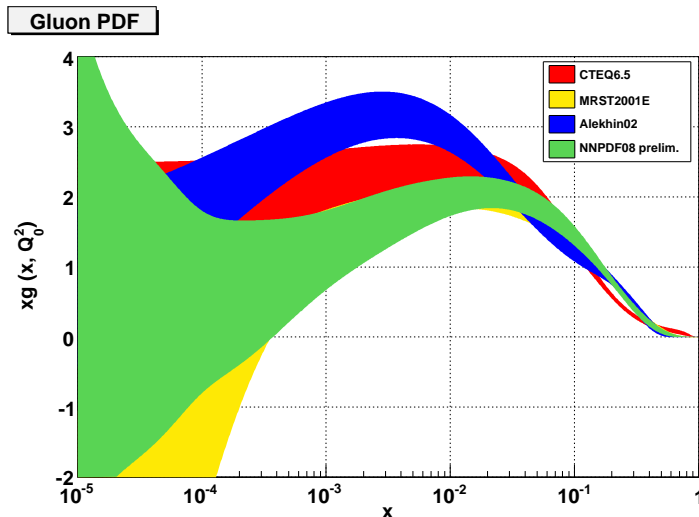
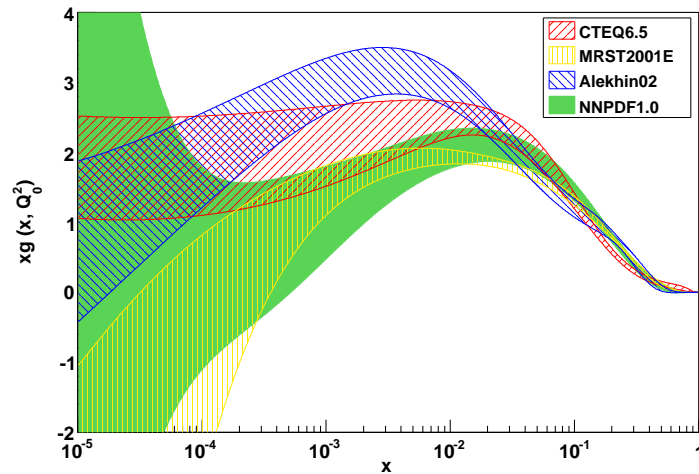
- where:

$$\langle f_i \rangle_{(1)} \equiv \frac{1}{N_{\text{rep}}^{(1)}} \sum_{k=1}^{N_{\text{rep}}^{(1)}} f_{ik}^{(1)},$$

$$\sigma^2[f_i^{(1)}] \equiv \frac{1}{N_{\text{rep}}^{(1)}(N_{\text{rep}}^{(1)} - 1)} \sum_{k=1}^{N_{\text{rep}}^{(1)}} \left(f_{ik}^{(1)} - \langle f_i \rangle_{(1)}\right)^2$$

- For statistically equivalent PDF sets: $\langle d[f] \rangle \sim \langle d[\sigma_f] \rangle \sim 1$

Stability under variation of the parametrization



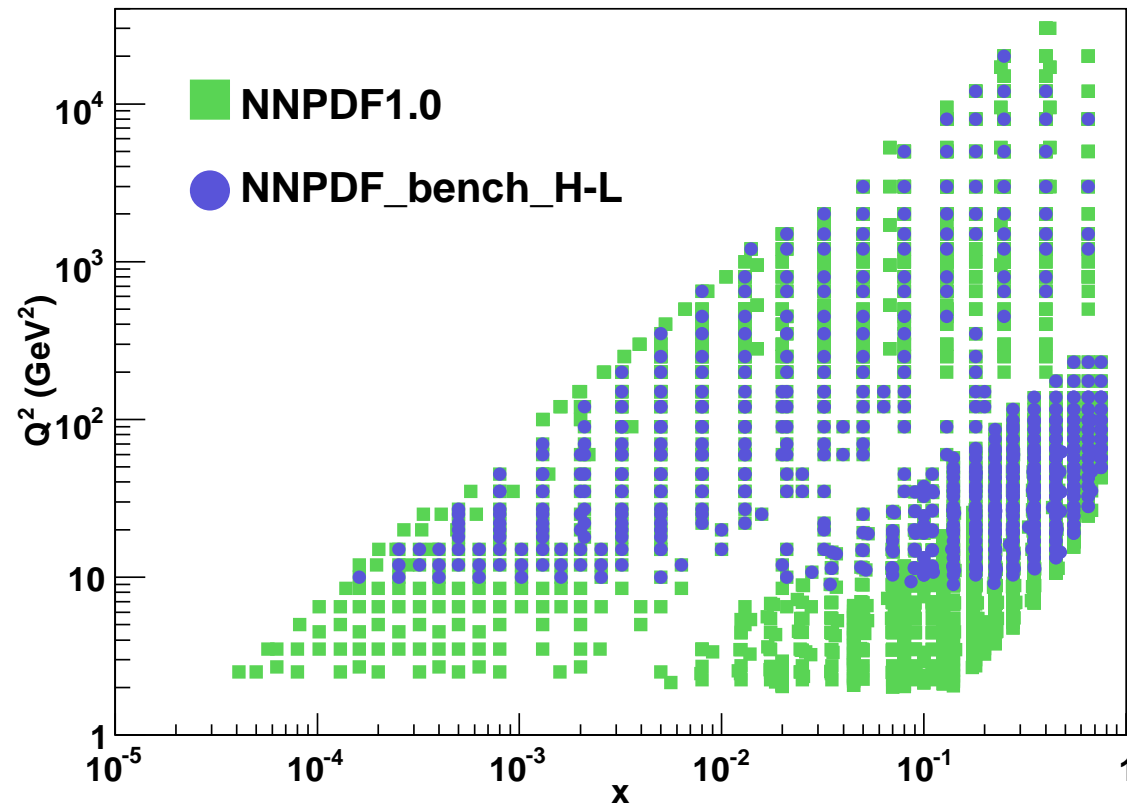
	Data	Extrapolation
$\Sigma(x, Q_0^2)$	$5 \cdot 10^{-4} \leq x \leq 0.1$	$10^{-5} \leq x \leq 10^{-4}$
$\langle d[f] \rangle$	0.98	1.25
$\langle d[\sigma] \rangle$	1.14	1.34
$g(x, Q_0^2)$	$5 \cdot 10^{-4} \leq x \leq 0.1$	$10^{-5} \leq x \leq 10^{-4}$
$\langle d[f] \rangle$	1.52	1.15
$\langle d[\sigma] \rangle$	1.16	1.07
$T_3(x, Q_0^2)$	$0.05 \leq x \leq 0.75$	$10^{-3} \leq x \leq 10^{-2}$
$\langle d[f] \rangle$	1.00	1.11
$\langle d[\sigma] \rangle$	1.76	2.27
$V(x, Q_0^2)$	$0.1 \leq x \leq 0.6$	$3 \cdot 10^{-3} \leq x \leq 3 \cdot 10^{-2}$
$\langle d[f] \rangle$	1.30	0.90
$\langle d[\sigma] \rangle$	1.10	0.98
$\Delta_S(x, Q_0^2)$	$0.1 \leq x \leq 0.6$	$3 \cdot 10^{-3} \leq x \leq 3 \cdot 10^{-2}$
$\langle d[f] \rangle$	1.04	1.91
$\langle d[\sigma] \rangle$	1.44	1.80

- Stability under change of architecture of the nets:
37 pars → **31 pars**
- Independence on the parametrization!

Dependence on data sets

HERA-LHC benchmark

Benchmark PDF fit to a reduced consistent set of DIS data [hep-ph/0511119]



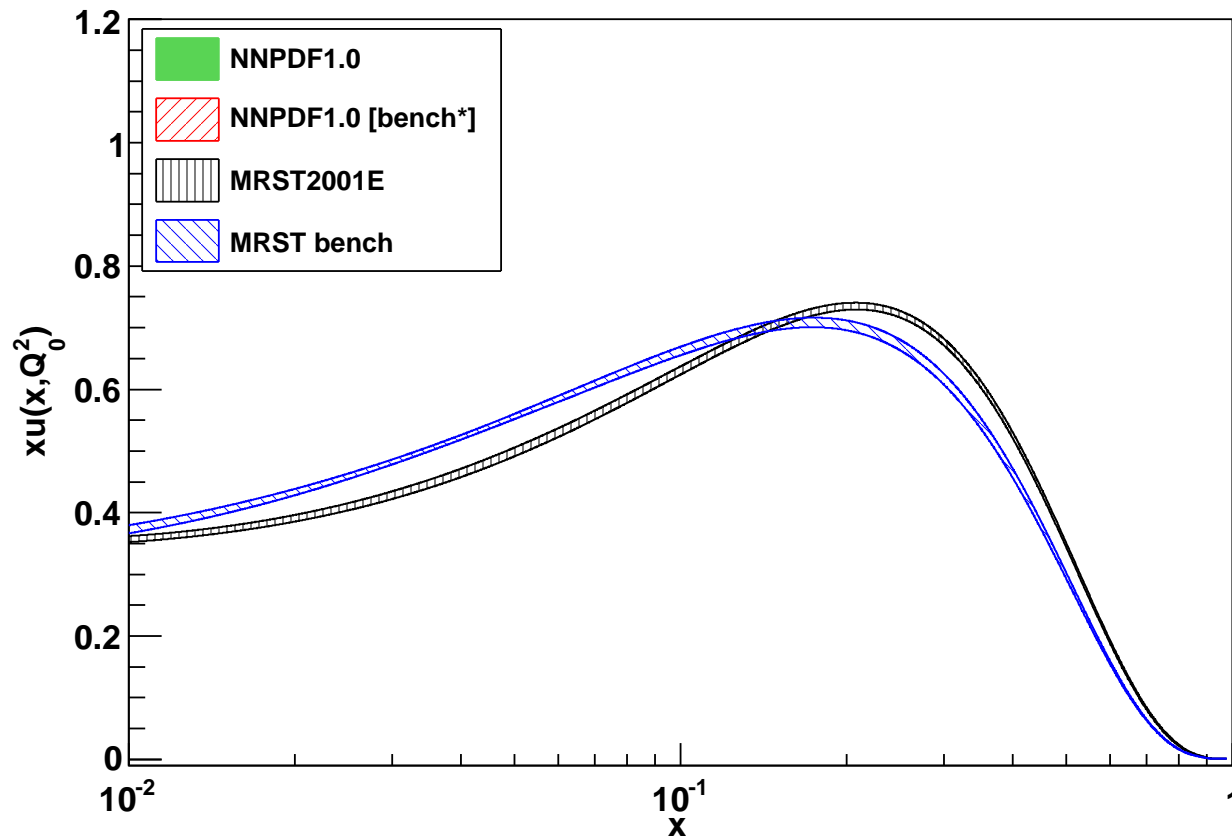
3163 data \longrightarrow 773 data

Dependence on data sets

HERA-LHC benchmark

Comparison between collaborations and between benchmark/global partons.

$u(x, Q^2 = 2\text{GeV}^2)$: MRST data region

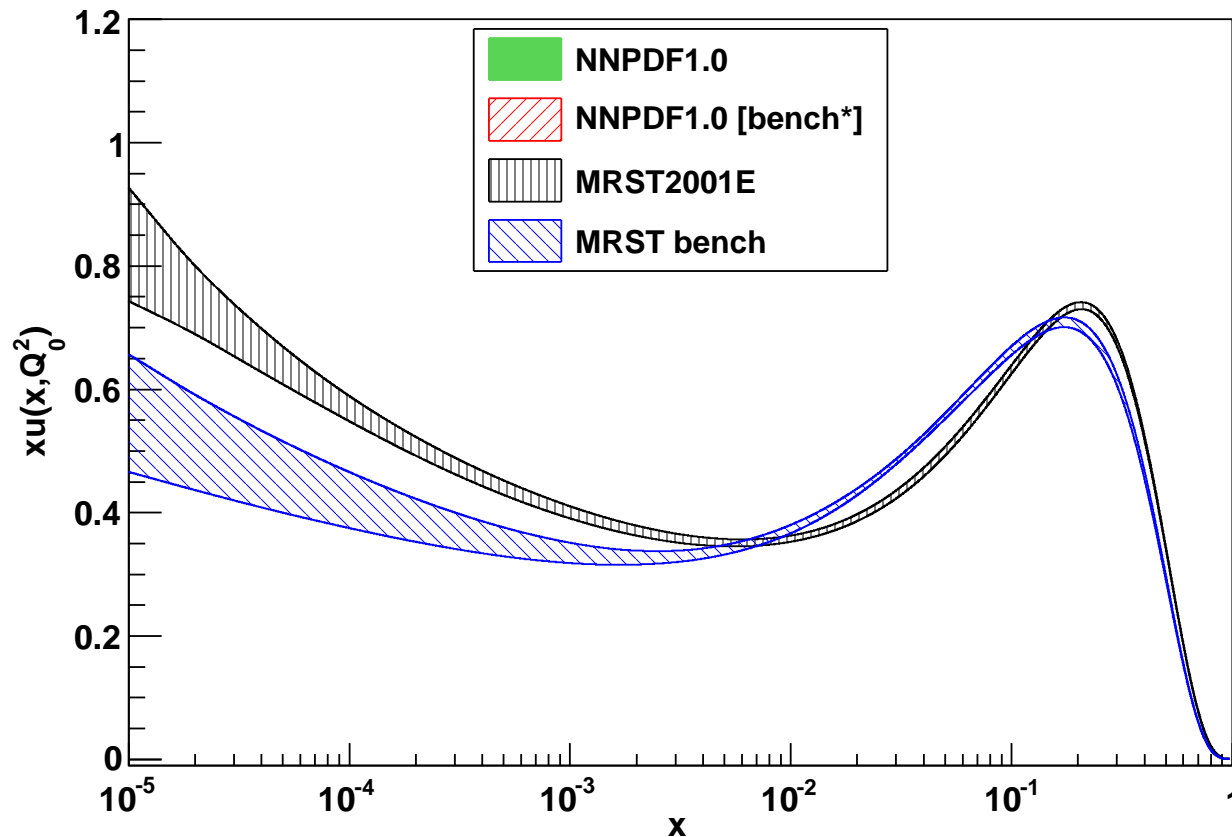


Dependence on data sets

HERA-LHC benchmark

Comparison between collaborations and between benchmark/global partons.

$u(x, Q^2 = 2\text{GeV}^2)$: MRST extrapolation region



Dependence on data sets

HERA-LHC benchmark

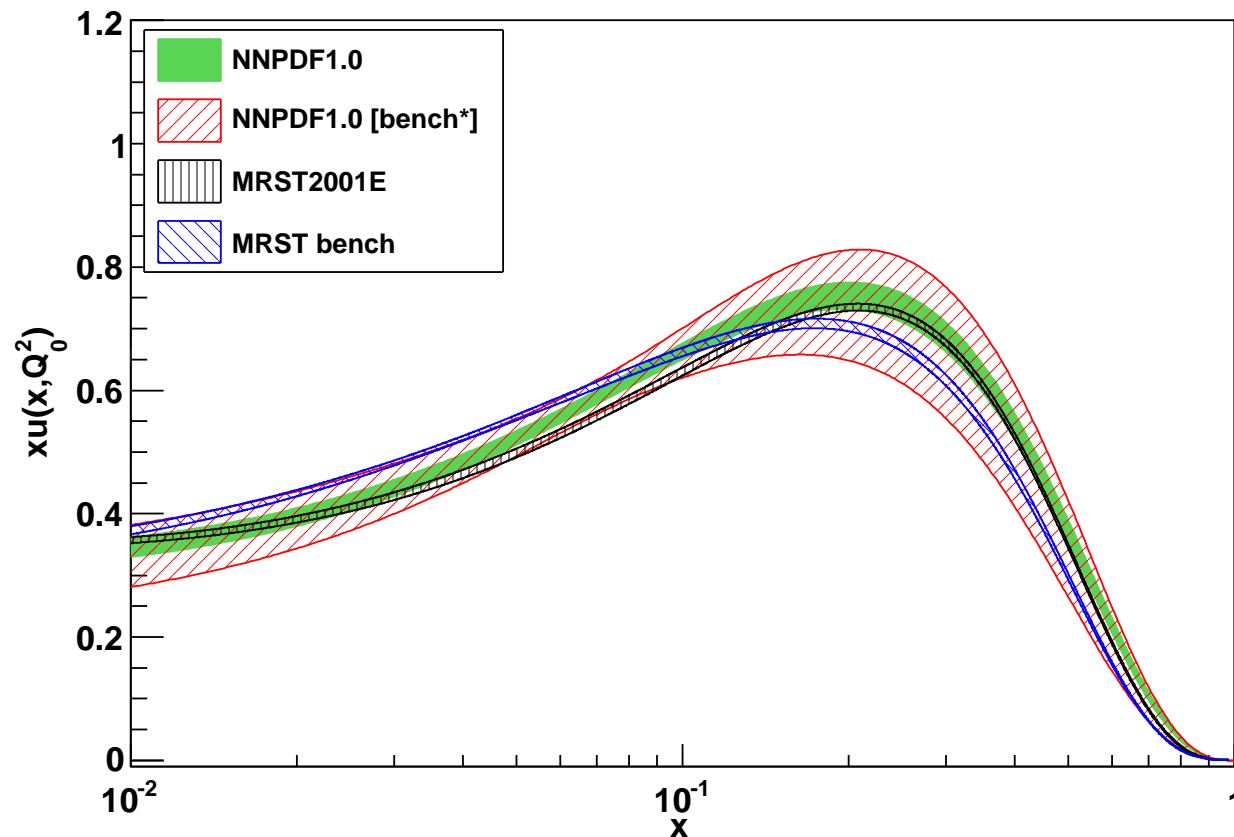
- benchmark partons and global partons do not agree within error!
- note that PDFs input parametrization, flavor assumptions and statistical treatment ($\Delta\chi^2_{\text{global}} = 50$, $\Delta\chi^2_{\text{bench}} = 1$) are tuned to data.
- not satisfactory especially to predict the behaviour of PDFs in the extrapolation region (LHC)

Dependence on data sets

HERA-LHC benchmark

Comparison between collaborations and between benchmark/global partons.

$u(x, Q^2 = 2\text{GeV}^2)$: data region

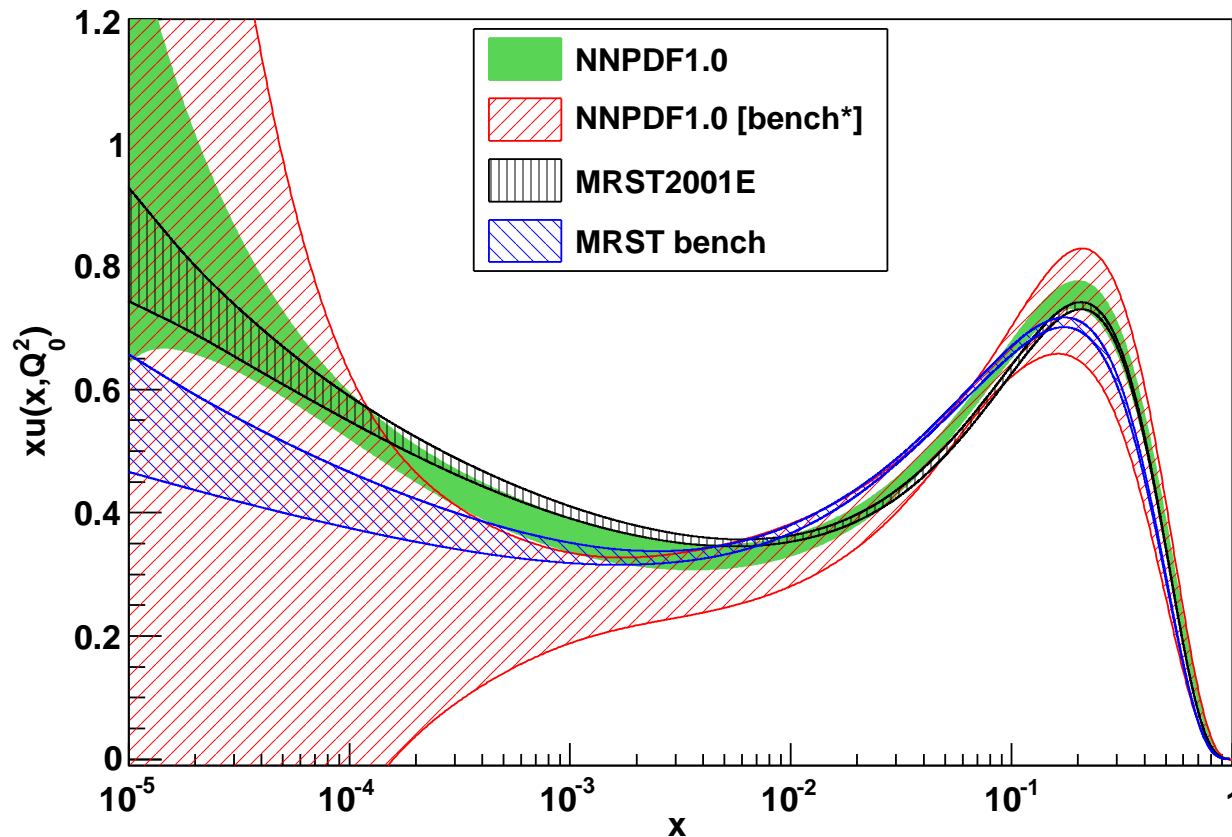


Dependence on data sets

HERA-LHC benchmark

Comparison between collaborations and between benchmark/global partons.

$u(x, Q^2 = 2\text{GeV}^2)$: extrapolation region

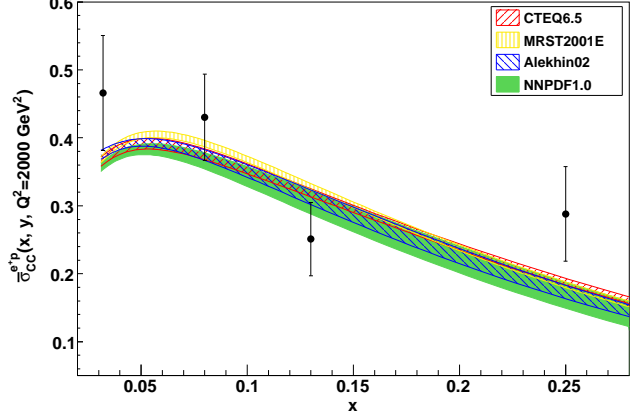
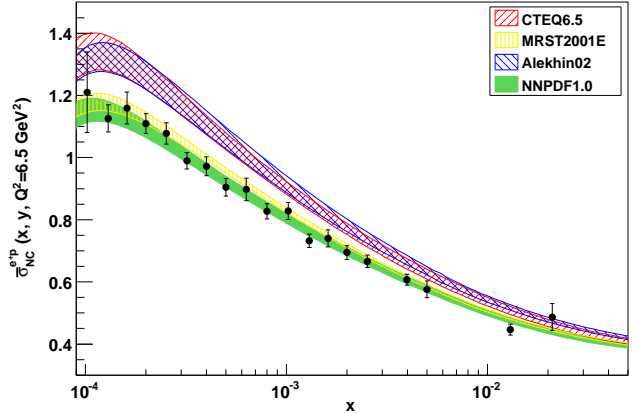
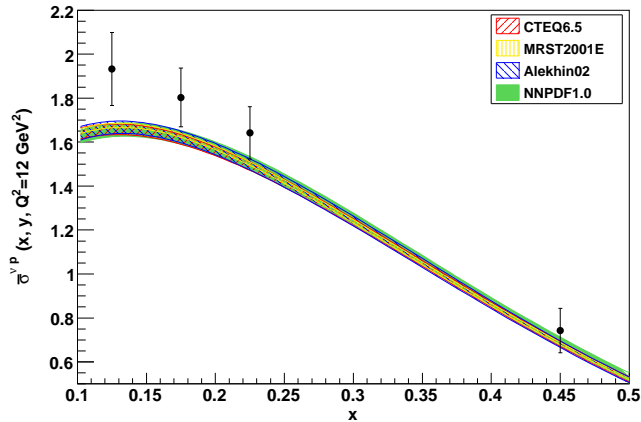
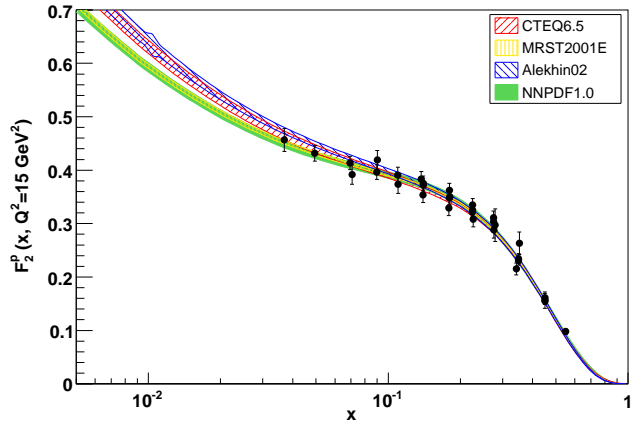


Dependence on data sets

HERA-LHC benchmark

- NNPDF1.0 is consistent with MRST global fit.
- NNPDFbench is consistent with NNPDF1.0 and MRST.
- Same parametrization and flavour assumption.
- Same statistical treatment.
- Underestimation of the error in the standard approach.

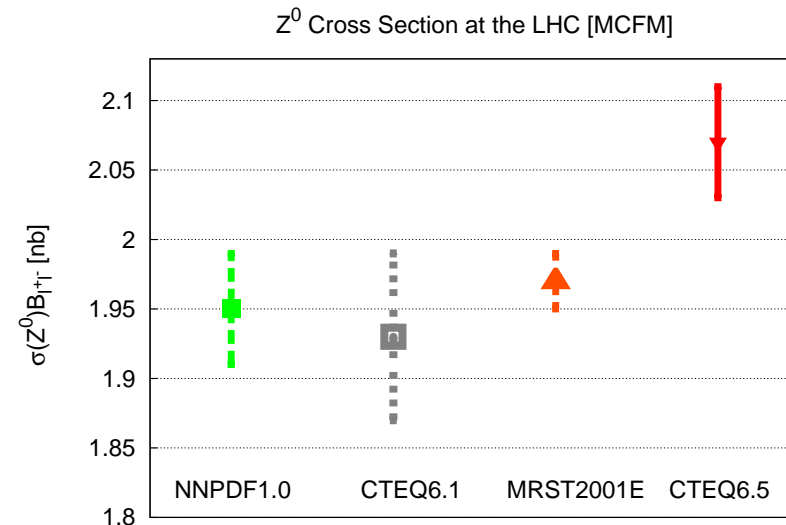
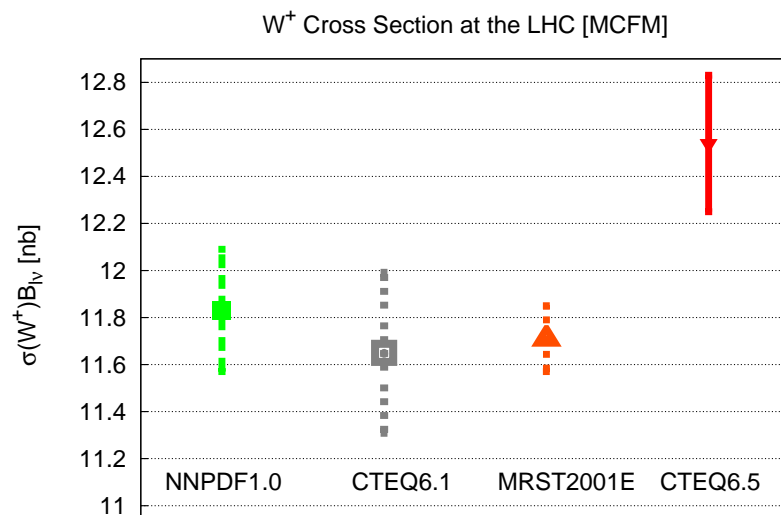
Comparison with present experimental data



LHC standard candle processes

- All quantities have been computed at NLO with MCFM [<http://mcfm.fnal.gov>]
- Quoted uncertainties are the 1σ bands due to the PDF uncertainty only.

	$\sigma_{W^+} \mathcal{B}_{l+\nu_l}$	$\Delta\sigma_{W^+} / \sigma_{W^+}$	$\sigma_Z \mathcal{B}_{l+l^-}$	$\Delta\sigma_Z / \sigma_Z$
NNPDF1.0	11.83 ± 0.26	2.2%	1.95 ± 0.04	2.1%
CTEQ6.1	11.65 ± 0.34	2.9%	1.93 ± 0.06	3.1%
MRST01	11.71 ± 0.14	1.2%	1.97 ± 0.02	1.0%
CTEQ6.5	12.54 ± 0.29	2.3%	2.07 ± 0.04	1.9%



Conclusions

- Standard approaches with fixed parametrization tend to underestimate uncertainties unless experimental errors are inflated by essentially arbitrary amount.
- **Monte Carlo** ensemble
 - Any statistical property of PDFs can be calculated using standard statistical methods.
 - No need of any tolerance criterion.
- The **Neural Network** parametrization
 - Small uncertainties come from an underlying physical law, not from parametrization bias.
 - Inconsistent data or underestimated uncertainties do not require a separate treatment and are automatically signalled by a larger value of the χ^2 .
- The first NNPDF parton set [arXiv:0808.1231] is available on the common LHAPDF interface [<http://projects.hepforge.org/lhapdf>].

Outlook

- Inclusion of **hadronic data** to
 - improve the accuracy of gluon at large x (jets)
 - determine the light antiquark sea asymmetry (Drell-Yan)
 - allow for a direct determination of the strange distribution (dimuon data)
- More accurate treatment of **Heavy Quark thresholds**.
- LO parton set in view of its use in Monte Carlo generators.
- More sophisticated theoretical treatment: NNLO parton distributions, large and small x resummation corrections should also be considered.
- Study of the impact of PDFs uncertainties on LHC phenomenology