

Higgs, Supersymmetry, and CP Violation

JAE SIK LEE

National Center for Theoretical Sciences, Hsinchu, TAIWAN





 \blacklozenge The $1^{\rm st}$ part: Introduction to

Higgs, Supersymmetry, CP Violation and the LHC

... Please don't be too picky

 \blacklozenge The 2nd part: My works on manifestations of

CP Violation through the Supersymmetric Higgs sector

at the LHC and other low-energy experiments

... Please don't feel too sleepy

♠ THE 1ST PART

THE 1^{st} PART

 \clubsuit Higgs (1/5)

• Why Particles Physics?

... to find answers to the Eternal Questions of

"What is the world made of?"

and

"What holds it together?"

The Particle Adventure

the fundamentals of matter and force
http://particleadventure.org/particleadventure/

\clubsuit Higgs (2/5)

• Answers to the Questions in 2011: Standard Model



They are all Massive except γ and g

The BIG question is :

How do they become massive?

\clubsuit Higgs (3/5)

• Spontaneous breaking of Electroweak symmetry Y. Nambu, PRL4(1960)380 Nobel Prize 2008 We are NOT living in "True Void" \Rightarrow The "Vacuum" is filled with Nonvanishing Field Strength of a scalar field called a Higgs boson : $\langle H \rangle_0 \neq 0$



Yukawa interactions $\Rightarrow m_l$ and m_q

Gauge interactions $\Rightarrow M_W$ and M_Z

* Something else may be needed for m_{ν}

\clubsuit Higgs (4/5)

- Status of the SM Higgs:
 - LEP bound: $M_H^{SM} \ge 114.4 \text{ GeV} (95 \% \text{ C.L.})$ ADLO, arXiv:hep-ex/0306033
 - Electroweak precision data: $M_H^{\rm SM} \lesssim 185$ GeV (95 % C.L.) direct-search limit included ACDDLOS, arXiv:0811.4682[hep-ex]



\clubsuit Higgs (4'/5)

- Status of the SM Higgs: ... continued
 - Tevatron exclusion: 158 GeV $\lesssim M_H^{\rm SM} \lesssim$ 175 GeV (95 % C.L.) CDF/D0, arXiv:1007.4587[hep-ex]



• *Higgs:* (4''/5)

• Status of the SM Higgs: *combining all...* Gfitter, arXiv:0811.0009



We anticipate the SM Higgs boson lighter than $\sim 200 \text{ GeV}$

Higgs = the God particle ? Leon Lederman, The God Particle: If the Universe Is the

Answer, What Is the Question?

- <u>elusive</u> ...
- controlling behind indirectly...



and ...

causing conceptual problems ...

 \clubsuit Supersymmetry (1/4)

• The Naturalness Problem

$$M_H^2 \sim m_0^2 - \frac{\Lambda^2}{16\pi^2}$$
 with $\Lambda \sim 10^{19} \text{ GeV}$ and $M_H \sim 10^2 \text{ GeV}$



 $\mathcal{P}\ll \mathsf{a} \ \mathsf{thunderbolt} \otimes \mathsf{a} \ \mathsf{lottery} \ \mathsf{ticket}$

Isn't it *natural* to assume there is *something* ?

\clubsuit Supersymmetry (2/4)

• Introducing Supersymmetry

 $\delta M_H^2 \sim -\chi^2$ (Particles) + χ^2 ("S"Particles) + $M_{\rm SUSY}^2 \log(\Lambda^2/M_{\rm SUSY}^2)$

 $M_{
m SUSY} \sim \mathcal{O}(1) \,\, {
m TeV}$

Great Discoveries at the LHC !



 \clubsuit Supersymmetry (3/4)

• Support for the existence of "S" Particles (I)



\blacklozenge Supersymmetry (4/4)

• Support for the existence of "S" Particles (II)



C. L. Bennett *et al.*, "First Year WMAP Observations: Preliminary Maps and Basic Results," Astrophys. J. Suppl. **148** (2003) 1 [arXiv:astro-ph/0302207]; M. Tegmark *et al.* [SDSS Collaboration], "Cosmological parameters from SDSS and WMAP," Phys. Rev. D **69** (2004) 103501 [arXiv:astro-ph/0310723]; P. Gondolo [arXiv:hep-ph/0501134] $\Omega_{\Lambda} = 0.72 \pm 0.08$ $\Omega_{m} = 0.27 \pm 0.016$

 $= \begin{cases} \Omega_b = 0.0448 \pm 0.0018 \\ \Omega_{\rm CDM} = 0.226 \pm 0.0016 \end{cases}$

Source: Robert Kindmer Source: NASA/WMAP Stience Teun

CDM = Lightest Supersymmetric Particle ?

 \blacklozenge CP violation (1/4)

• What is CP Violation?:



http://universe-review.ca

If Nature treated

Left-handed particle moving right and Right-handed antiparticle moving left differently, we are saying CP is violated

\blacklozenge CP violation (2/4)

Why CP Violation?: ⇒ There IS CP Violation! M. Kobayashi and T. Maskawa, Prog. Theor.
 Phys. 49 (1973) 652 Nobel Prize 2008



\blacklozenge CP violation (3/4)

• Then, who ordered "more" CP violation beyond the SM CKM phase? A. D. Sakharov, JETP Letters 5(1967)24



CP violation in the SM is too weak to explain the matter dominance of the Universe J. Cline, arXiv:hep-ph/0609145

The matter-dominated Universe did!

 \blacklozenge CP violation (4/4)

- Even the minimal SUSY model contains many possible sources of CP violation:
 - Gaugino mass terms: $3 \oplus 3 = 6$

$$-\mathcal{L}_{\text{soft}} \supset \frac{1}{2} (M_3 \,\widetilde{g} \widetilde{g} + M_2 \,\widetilde{W} \widetilde{W} + M_1 \,\widetilde{B} \widetilde{B} + \text{h.c.})$$

- Trilinear a terms $\mathbf{a}_{\mathbf{f}ij} \equiv \mathbf{h}_{\mathbf{f}ij} \cdot \mathbf{A}_{\mathbf{f}ij}$: $3 \times (3 \oplus 6 \oplus 9) = 54$

$$-\mathcal{L}_{\text{soft}} \supset (\widetilde{u}_R^* \, \mathbf{a_u} \, \widetilde{Q} H_2 - \widetilde{d}_R^* \, \mathbf{a_d} \, \widetilde{Q} H_1 - \widetilde{e}_R^* \, \mathbf{a_e} \, \widetilde{L} H_1 + \text{h.c.})$$

- Sfermion mass terms: $5 \times (3 \oplus 3 \oplus 3) = 45$

$$-\mathcal{L}_{\text{soft}} \supset \widetilde{Q}^{\dagger} \,\mathbf{M}_{\widetilde{\mathbf{Q}}}^{\mathbf{2}} \,\widetilde{Q} + \widetilde{L}^{\dagger} \,\mathbf{M}_{\widetilde{\mathbf{L}}}^{\mathbf{2}} \,\widetilde{L} + \widetilde{u}_{R}^{*} \,\mathbf{M}_{\widetilde{\mathbf{u}}}^{\mathbf{2}} \,\widetilde{u}_{R} + \widetilde{d}_{R}^{*} \,\mathbf{M}_{\widetilde{\mathbf{d}}}^{\mathbf{2}} \,\widetilde{d}_{R} + \widetilde{e}_{R}^{*} \,\mathbf{M}_{\widetilde{\mathbf{e}}}^{\mathbf{2}} \,\widetilde{e}_{R}$$

1 KM phase $(SM) \implies 1 + 45$ (Minimal SUSY Model)

♦ *LHC* (1/6)

• LHC = the Large Hadron Collider



♠ *LHC* (1'/6)

• LHC = the Large Hadron Collider



"Human beings can't see anything without wanting to destroy it, Lyra, *That's* original sin. And I'm going to destroy it."

... Lord Asriel

HIS DARK MATERIALS trilogy #1

♦ *LHC* (2/6)



- Two protons collide with a kinetic energy of 14 TeV $_1~\text{GeV}\simeq$ a proton mass
- then the two protons are *destroyed* to disappear and we will see something ...

♦ *LHC* (3/6)

- Past and present ...
 - 1989 : R&D starts
 - 1997 : Construction starts
 - 2003 : Underground installation starts
 - 2008 : Installation completed

11 Sep 2008 : First beam circulated 20 Sep 2008 : Magnet failures ...

23 Nov 2009 : First collisions at $\sqrt{s} = 900$ GeV 8 Dec 2009 : First collisions at $\sqrt{s} = 2.36$ TeV 30 Mar 2010 : First collisions at $\sqrt{s} = 7$ TeV



From Junichi Kanzaki's talk @ KEKPH2010

\clubsuit LHC (3'/6)



• ... and near/far future of the LHC

2010 ~ 2011 : 7 TeV (3.5+3.5 TeV) runs until 1 fb⁻¹ LHC 1/2 \Rightarrow For SUSY searches with 35 pb⁻¹, see next page

2012 ~ 2012 : Shutdown for repairing to achieve 7 (6.5) TeV/beam collisions 2013 ~ 2014 : Go to 7 (6.5) + 7 (6.5) TeV runs to get ~ 10 fb⁻¹ LHC

2015 \sim 202? : accumulate 3000 fb⁻¹ with 250 \sim 300 fb⁻¹ per year

again, from Junichi Kanzaki's talk @ KEKPH2010, or LHC Performance Workshop - Chamonix 2010 (Jan)

\blacklozenge Initial SUSY searches with 35 pb⁻¹ at the LHC

• CMS (jets + missing E_T ; left) and ATLAS (lepton + jets + missing E_T ; right) CMS, PLB698(2011)196,2011, arXiv:1101.1628 [hep-ex] and ATLAS, arXiv:1102.2357 [hep-ex]



ATLAS, more stringent ...

• ... and near/far future of the LHC

2010 ~ 2011 : 7 TeV (3.5+3.5 TeV) runs until 1 fb⁻¹ $\underline{LHC 1/2}$ 2012 ~ 2012 : Shutdown for repairing to achieve 7 (6.5) TeV/beam collisions 2013 ~ 2014 : Go to 7 (6.5) + 7 (6.5) TeV runs to get ~ 10 fb⁻¹ \underline{LHC} 2015 ~ 202? : accumulate 3000 fb⁻¹ with 250 ~ 300 fb⁻¹ per year

> again, from Junichi Kanzaki's talk @ KEKPH2010, or LHC Performance Workshop - Chamonix 2010 (Jan)

• ... and near/far future of the LHC: ... continued

3000 fb⁻¹ \approx 2,400,000,000 $t\bar{t}$ pairs : Top factory

 \approx 100,000,000 130 GeV SM Higgs bosons

 \thickapprox 3,000,000 $\mathcal{O}(1)$ TeV gluinos + squarks

 \thickapprox 300,000 $\mathcal{O}(1)$ TeV charginos + neutralinos



♦ *LHC* (6/6)

- ... and near/far future of the L
 - 3000 fb⁻¹ \approx 2,400,000,000 $t\bar{t}$

 \approx 100,000,000 130

 \approx 3,000,000 $\mathcal{O}(1)$ 7

pprox 300,000 $\mathcal{O}(1)$ Te

♠ THE 2ND PART

THE 2^{nd} PART

\blacklozenge Preliminary (1/4)

• The Nobel Prize in Physics 2008: "Passion for symmetry" (broken?)



"for the discovery of the mechanism of spontaneous broken symmetry in subatomic physics" "for the discovery of the origin of the broken symmetry which predicts the existence of at least three families of quarks in nature"



\blacklozenge Preliminary (2/4)

• My works aim at ...:



\blacklozenge Preliminary (3/4)

- Introducing SUSY \implies Introducing lots of additional CP phases ! :
 - 1 KM phase (SM) \implies 1 + 45 (Minimal SUSY Model)

My Question:

How does the SUSY CP violation manifest itself through the Higgs sector ?

more specifically,

Where to find it? and Why through Higgs sector?

\blacklozenge Preliminary (4/4)



http://www.er.doe.gov/hep/vision/index.shtml

\blacklozenge Preliminary (4'/4)

- Why through Higgs sector?
 - Experimental Side: LHC
 - * Higgs boson will be searched most intensively
 - * Higgs boson will be studied most extensively
 - Theoretical Side: Quantum sensitivity
 - * Large quantum correction
 - * Resonant enhancement

♠ CPsuperH

• Tool: CPsuperH : \triangle RG-improved effective potential approach



For *experimentalists* as well as phenomenologists who are exploring the MSSM Higgs sector both in *CP*-conserving and CP-violating cases V. M. Abazov *et al.* [D0 Collaboration], "Search for neutral supersymmetric Higgs bosons in multijet events at $s^{**}(1/2) = 1.96$ -TeV," Phys. Rev. Lett. **95** (2005) 151801 [arXiv:hep-ex/0504018];

A. Abulencia *et al.* [CDF Collaboration], "Search for charged Higgs bosons from top quark decays in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ -TeV," Phys. Rev. Lett. **96** (2006) 042003 [arXiv:hep-ex/0510065].

 Consistent computational tool for both Energy- and Intensity-Frontier experiments including several B observables and EDMs

♦ Where to find it ?



http://www.er.doe.gov/hep/vision/index.shtml
\blacklozenge SM Higgs (1/2)

• LHC Discovery Significance of the SM Higgs ATLAS TDR, arXiv:0901.0512[hep-ex]



- $H \rightarrow \tau \tau$
- $H \rightarrow \gamma \gamma$
- $\triangleright H \to WW^* \to l\nu l\nu \text{ (GF)}$
- $\triangleright H \rightarrow WW^* \rightarrow l\nu l\nu \text{ (VBF)}$
- $\Box \quad H \to ZZ^{(*)} \to 4 \, l$

The WW channel plays a dominant role for 130 GeV $\lesssim M_H \lesssim 190$ GeV ! \bigstar SM Higgs (2/2)

• The 2 missing neutrinos: Transverse mass vs. MAOS Higgs mass



K. Choi, S. Choi, JSL, C.B. Park, PRD**80** (2009) 073010; K. Choi, JSL, C.B. Park, arXiv:1008.2690 [hep-ph], to appear in PRD

\bigstar MSSM Higgs (1/2)

• LEP Limits on the MSSM Higgs: A. Djouadi, hep-ph/0503173 and references there in: P. Bechtle, CPNSH Report, CERN-2006-009, hep-ph/0608079; ADLO, hep-ex/0602042, Combined results with FeynHiggs



 \bigstar MSSM Higgs (2/2)

• LHC Discovery of the MSSM Higgs: ATLAS TDR(1999); M. Schumacher, CPNSH Report, CERN-2006-009, hep-ph/0608079



\blacklozenge b-quark Fusion (1/2)

• Strong dependence of Higgs-boson production via *b*-quark fusion on CP phases: CPX Scenario with $\Phi_3 = 0^\circ$ or $\Phi_3 = 180^\circ$ and $M_{H_1} = 115$ GeV. F. Borzumati and JSL, PLB595(2004)347 and PLB641(2006)486



* Solid: $\Phi_3 = 180^{\circ}$ and Dashed: $\Phi_3 = 0^{\circ}$

* Around $\Phi_{A\mu} = 100^{\circ}$, $M_{H_2} - M_{H_1} \simeq 3$ GeV \Longrightarrow Distinguishable ?

\blacklozenge b-quark Fusion (2/2)

• Invariant mass distribution when Higgs bosons decay into two γ 's and/or $\mu^+\mu^-$

At the LHC, the experimental resolution $\delta M_{\gamma\gamma} \sim 1$ GeV and $\delta M_{\mu\mu} \sim 3$ GeV ATLAS TDR, Vol 2, CERN-LHCC-99-15, ATLAS-TDR-15 (1999); CMS Physics TDR, Volume II, CERN-LHCC-2006-021, 25 June 2006



\blacklozenge CP Asymmetries (1/3)

• <u>CP asymmetry in Higgs decays in the tau leptons</u>: J. Ellis, JSL and A. Pilaftsis, D71(2005)075007



 σ_{RR} σ_{LL}

$$\mathcal{A}_{CP} = \frac{\sigma_{RR} - \sigma_{LL}}{\sigma_{RR} + \sigma_{LL}} \quad \text{with} \quad \begin{cases} \sigma_{RR} \equiv \sigma(pp \to H \to \tau_R^+ \tau_R^- X) \\ \sigma_{LL} \equiv \sigma(pp \to H \to \tau_L^+ \tau_L^- X) \end{cases}$$

\blacklozenge CP Asymmetries (2/3)

• Resonant transitions between Higgs bosons enhance the asymmetry!:



The neutral Higgs bosons should not be treated separately !

J. Ellis, JSL and A. Pilaftsis, D70(2004)075010

 \blacklozenge CP Asymmetries (3/3)

• Integrated CP asymmetry \mathcal{A}_{CP}^{WW} with $\sigma_{tot}^{WW} > 0.2$ pb:



CP asymmetry is large over the whole ranges of CP phases !

\blacklozenge Diffraction (1/3)

<u>Diffraction</u> is a process with a *t*-channel exchange leading to a Large Rapidity Gap
 M. Arneodo, M. Diehl, hep-ph/0511047



The differential cross section

 $\left. \frac{d\sigma \left(pp \to pp \right)}{d \, t} \right|_{\text{elastic}}$

has a resemblance to the diffraction pattern in optics:



\blacklozenge Diffraction (2/3)

• Exclusive double diffractive process



Outgoing protons remain intact & scatter through small angles with $p_{\perp} \ll 1 \text{ GeV}$

- J. Ellis, JSL and A. Pilaftsis, D71(2005)075007
- * Need TOTEM/ATLASFP detectors

† $M^2 = s - 2\sqrt{s}(E'_1 + E'_2) + (p'_1 + p'_2)^2$ ⇒ δM ~ 1 GeV !

Clean environment due to the large rapidity gap

A. De Roeck, *et al.*, EPJC**25** (2002) 391; M. Boonekamp, *et al.*, PLB**598** (2004) 243; K. A. Assamagan *et al.* hepph/0406152; B. E. Cox, hep-ph/0409144, hep-ph/0501064; <u>B.E. Cox*, J.R. Forshaw, JSL, J.W. Monk* and A. Pilaftsis,</u> PRD**68** (2003) 075004; V.A. Khoze, A.D. Martin and M.G. Ryskin, EPJC**34** (2004) 327

\blacklozenge Diffraction (3/3)

•
$$pp \to pH_ip; H_i \to \tau^+\tau^- (\Phi_A, \Phi_3) = (90^\circ, -90^\circ) (\Phi_A, \Phi_3) = (90^\circ, -10^\circ)$$



Invariant mass distribution can be studied !

\blacklozenge Beyond the Minimal SUSY Model (1/2)

• The μ problem J. E. Kim and H. P. Nilles, Phys. Lett. B 138, 150 (1984)



K. Cheung, T. J. Hou, JSL and E. Senaha, PRD82(2010)075007

\blacklozenge Beyond the Minimal SUSY Model (2/2)

- We are working on ...
 - Constraints from EDMs on the new CP violation
 - ElectroWeak Phase Transitions
 - New CP-phase driven baryogenesis
 - B- and K-meson observables
 - Collider signatures of the light Higgs bosons
 - Higgs Inflation(?)
 - etc

♠ Summary (LHC)

- Phases of the LHC :
 - I (2013 ~ 2014): LHC with $\mathcal{L} \sim 10 \ fb^{-1}$

 $\begin{array}{ll} & \mbox{SM Higgs Discovery,} \\ & \mbox{Non-SM-like Higgs decays and productions, Multiple Higgs peaks} \\ & \mbox{II} & (2015 \sim 2016): & LHC \ with \ \mathcal{L} \sim 300 \ fb^{-1} \\ & \mbox{I} \oplus \ \mbox{Strongly-coupled Higgs, CP asymmetries, Final-state spin-spin correlations} \end{array}$

III (2017 ~ 202?) : LHC with $\mathcal{L} \gg 300 \ fb^{-1}$

 $I \oplus II \oplus \mathsf{Exotic}$ Higgs productions

Conclusions and Future Prospects

- Spontaneous breaking of the electroweak symmetry *may* be responsible for the masses of SM particles
- Supersymmetry *may* stabilize the Higgs-boson masses against the Quantum corrections
- The additional CP phases in the models extended by Supersymmetry *may* explain the matter dominance of the Universe
- The Supersymmetric CP phases *may* manifest themselves through the Higgs sector in Energy-, Intensity-, and Cosmic-Frontier experiments
- The LHC is to give definite answers to these *may*'s

♦ Where to find it ?



http://www.er.doe.gov/hep/vision/index.shtml

$\blacklozenge B \ Observables \ (1/4)$

• Higgs-mediated B-meson mixing and rare decays: LHCb and SuperB J. Ellis, JSL and A. Pilaftsis, PRD76(2007)115011

$$\left(-\mathcal{L}_{\text{eff}}^{d} \right)^{H} = \frac{g}{2M_{W}} \overline{d} \, \widehat{M}_{d} \, \mathbf{g}_{H_{i}\bar{d}d}^{L} P_{L} \, d \, H_{i} + \text{h.c.}$$
$$\mathbf{g}_{H_{i}\bar{d}d}^{L} \propto V_{\text{CKM}}^{\dagger} \left[\mathbf{1} + \tan\beta \, \Delta_{d} \right]^{-1} V_{\text{CKM}}$$



Non-universal Quantum interactions inducing the mixing and rare decays !

$\blacklozenge B \ Observables (2/4)$

• $\Delta M_{B_d}^{\rm SUSY}$ and $\Delta M_{B_s}^{\rm SUSY}$ as functions of $\tan \beta(M_{\rm SUSY})$ for three values of $\Phi_M \equiv \Phi_1 = \Phi_2 = \Phi_3$ SPS1a-like scenario $\widetilde{M}_{L,E} = 200$ GeV and $\Phi_A^{\rm GUT} = 0^\circ$



$\blacklozenge B \ Observables (3/4)$

• $B(\overline{B}^0_s \to \mu^+ \mu^-)$ and the ratio $R_{B\tau\nu}$ as functions of $\tan \beta(M_{SUSY})$ for three values of $\Phi_M \equiv \Phi_1 = \Phi_2 = \Phi_3$ SPS1a-like scenario $\widetilde{M}_{L,E} = 200$ GeV and $\Phi_A^{GUT} = 0^\circ$



$\blacklozenge B \ Observables (4/4)$

• $B(B \to X_s \gamma)$ and $\mathcal{A}_{CP}^{dir}(B \to X_s \gamma)$ as functions of Φ_M for four values of $\tan \beta(M_{SUSY}) = 10, 20, 30, 40$ SPS1a-like scenario $\widetilde{M}_{L,E} = 200$ GeV and $\Phi_A^{GUT} = 0^\circ$



♦ Where to find it ?



$\bullet EDMs (1/4)$

• Electric Dipole Moments (EDMs): T violation ⇒ CP violation (under CPT)



- $|d_{\rm Tl}| < 9 \times 10^{-25} {\rm e\,cm}, |d_{\rm Hg}| < 3.1 \times 10^{-29} {\rm e\,cm}, |d_{\rm n}| < 3 \times 10^{-26} {\rm e\,cm}$ B. C. Regan, E. D. Commins, C. J. Schmidt and D. DeMille, PRL**88** (2002) 071805; W. C. Griffith, M. D. Swallows, T. H. Loftus, M. V. Romalis, B. R. Heckel and E. N. Fortson, PRL**102** (2009) 101601; C. A. Baker *et al.*, PRL**97** (2006) 131801
- No large CP phases? -No! But, one needs to introduce some hierarchy between first two and third generations and/or some moderate tuning among parameters

\clubsuit EDMs (2/4)

• EDM constraints and *cancellation*: CPX Scenario with $\Phi_{A_{t,b}} = \Phi_3 = 90^\circ$, $\Phi_{A_{u,d,c,s}}$ $= 0^{\circ}$ $\tan\beta$ = 5, $M_{H^{\pm}}$ = 300 GeV, $|M_2|$ = $2|M_1|$ = 100 GeV, and $M_{\rm SUSY}$ = 0.5 TeV The hierarchy factor ρ : $M_{\tilde{X}_{1,2}} = \rho M_{\tilde{X}_3}$ with X = Q = U = D = L = E J. Ellis, JSL and A. Pilaftsis, JHEP08(2008)049 10 10³ $[(\Phi_1, \Phi_2) = (0^{\circ}, 90^{\circ})]$ $[(\Phi_1, \Phi_2) = (0^{\circ}, 90^{\circ})]$ $(\Phi_1, \Phi_2) = (0^{\circ}, 90^{\circ})$ EXP $|\mathbf{d}_{\mathrm{TI}}/\mathbf{d}^{\mathrm{EXP}}|$ 10² 10^{2} 10² d_{Hg} 10 10 10 1 1 1 -1 10 -1 10 10 10 10 10 1 1 1 ρ ρ ρ -24 10³ 10 10 -23 $(\Phi_1, \Phi_2) = (\mathbf{0}^0, \mathbf{90}^0)$ d^E_e [e cm] $(\Phi_1, \Phi_2) = (0^{\circ}, 90^{\circ})$ $(\Phi_1, \Phi_2) = (0^{\circ}, 90^{\circ})$ cm EXP -24 -25 10² **_ 10** $\mathbf{d}^{\mathbf{G}}$ $\mathbf{H}^{\mathbf{0}}$ -26 $\mathbf{H}^{\mathbf{0}}$ 10 10 -26 $\mathbf{d}_{\mathbf{u},\mathbf{d}}^{\mathbf{C}}$ 10 -27 $\mathbf{d}_{\mathbf{u},\mathbf{d}}^{\mathbf{E}}$ χ^{0} 1 10 -27 10 -28 -28 -1 10 10 10 10 10 1 1 1 ρ ρ ρ



$\bullet EDMs (4/4)$

• Optimal direction : the higher-dimensional generalization with N CP phases and n EDM constraints

$$\Phi_{\alpha} = \epsilon_{\alpha\beta_{1}\cdots\beta_{n}\gamma_{1}\cdots\gamma_{N-n-1}} E_{\beta_{1}}^{(1)}\cdots E_{\beta_{n}}^{(n)} B_{\gamma_{1}\cdots\gamma_{N-n-1}}$$

where (N-n-1)-dimensional B form is

$$B_{\gamma_1\cdots\gamma_{N-n-1}} = \epsilon_{\gamma_1\cdots\gamma_{N-n-1}\sigma\beta_1\cdots\beta_n} O_{\sigma} E_{\beta_1}^{(1)}\cdots E_{\beta_n}^{(n)}$$

♦ Where to find it ?



http://www.er.doe.gov/hep/vision/index.shtml

\blacklozenge Dark Matter (1/1)

• LEP Holes \Rightarrow Very light 7.5 GeV Higgs signals through the inclusive process $\chi\chi \rightarrow H \rightarrow \gamma$, \bar{p} JSL and S. Scopel, Phys. Rev. D **75** (2007) 075001 [arXiv:hep-ph/0701221]



 \bigstar MAÔS momenta of the missing neutrinos (1/2)

• $MAOS = M_{T2}$ -Assisted On-shell Scheme:

$$H \to W(p + \mathbf{k}) W(q + \mathbf{l}) \to l(p) \nu(\mathbf{k}) l'(q) \nu'(\mathbf{l})$$

$$(M_H^2)^{\mathrm{MA}\hat{\mathrm{O}}\mathrm{S}} = (p + \mathbf{k}^{\mathrm{MA}\hat{\mathrm{O}}\mathrm{S}} + q + \mathbf{l}^{\mathrm{MA}\hat{\mathrm{O}}\mathrm{S}})^2$$

$$\mathbf{k}_T^{\text{MAÔS}} = -\mathbf{q}_T \text{ and } \mathbf{l}_T^{\text{MAÔS}} = -\mathbf{p}_T$$

$$\boldsymbol{k}_{L}^{\text{MAÔS}} = \frac{|\mathbf{k}_{T}^{\text{MAÔS}}|}{|\mathbf{p}_{T}|} p_{L}, \quad \boldsymbol{l}_{L}^{\text{MAÔS}} = \frac{|\mathbf{l}_{T}^{\text{MAÔS}}|}{|\mathbf{q}_{T}|} q_{L}$$

 \bigstar MAÔS momenta of the missing neutrinos (2/2)

 \bullet Is the modified $MA\hat{O}S$ scheme working well?:



h M_{T2} : 1/2

• Transverse Mass M_T :

$$V(p) \qquad p^{\mu} = (\sqrt{m_{V}^{2} + |\mathbf{p}_{T}|^{2} + p_{L}^{2}}, \mathbf{p}_{T}, p_{L})$$

$$k^{\mu} = (\sqrt{m_{\chi}^{2} + |\mathbf{k}_{T}|^{2} + k_{L}^{2}}, \mathbf{k}_{T}, k_{L})$$

$$\mathbf{k}_{T} = \mathbf{p}_{T} (fully \ determined)$$

$$M_{X}^{2} = m_{V}^{2} + m_{\chi}^{2} + 2 \left(\sqrt{m_{V}^{2} + |\mathbf{p}_{T}|^{2} + p_{L}^{2}} \sqrt{m_{\chi}^{2} + |\mathbf{k}_{T}|^{2} + k_{L}^{2}} - \mathbf{p}_{T} \cdot \mathbf{k}_{T} - p_{L}k_{L}\right)$$

$$M_{T}^{2} = m_{V}^{2} + m_{\chi}^{2} + 2 \left(\sqrt{m_{V}^{2} + |\mathbf{p}_{T}|^{2}} - \sqrt{m_{\chi}^{2} + |\mathbf{k}_{T}|^{2}} - \mathbf{p}_{T} \cdot \mathbf{k}_{T} - p_{L}k_{L}\right)$$

$$M_{T}^{2} = m_{V}^{2} + m_{\chi}^{2} + 2 \left(\sqrt{m_{V}^{2} + |\mathbf{p}_{T}|^{2}} - \sqrt{m_{\chi}^{2} + |\mathbf{k}_{T}|^{2}} - \mathbf{p}_{T} \cdot \mathbf{k}_{T} - p_{L}k_{L}\right)$$

$$M_{T}^{2} = m_{V}^{2} + m_{\chi}^{2} + 2 \left(\sqrt{m_{V}^{2} + |\mathbf{p}_{T}|^{2}} - \sqrt{m_{\chi}^{2} + |\mathbf{k}_{T}|^{2}} - \mathbf{p}_{T} \cdot \mathbf{k}_{T} - p_{L}k_{L}\right)$$

$\spadesuit M_{T2}: 2/2$

• ... then M_{T2} ?: Generalization of M_T when there are 2 missing particles



$$\mathbf{k}_T + \mathbf{l}_T = \mathbf{p}_T$$

Only determined are ... the 2 among the 4 unknown degrees of freedom !

... but we know:

$$\max\{M_T^{(1)}, M_T^{(2)}\}(\mathbf{k}_T^{\text{true}}, \mathbf{l}_T^{\text{true}}) \le M_X$$

$M_{T2}: 2'/2$

• ... then M_{T2} ?: Generalization of M_T when there are 2 missing particles



$$\mathbf{k}_T + \mathbf{l}_T = \mathbf{p}_T$$

Only determined are ... the 2 among the 4 unknown degrees of freedom !

... but we know:

$$\max\{M_T^{(1)}, M_T^{(2)}\}(\mathbf{k}_T^{\text{true}}, \mathbf{l}_T^{\text{true}}) \le M_X$$

$$\max\{M_T^{(1)}, M_T^{(2)}\}(\mathbf{k}_T^{\min}, \mathbf{l}_T^{\min}) \le \max\{M_T^{(1)}, M_T^{(2)}\}(\mathbf{k}_T^{\text{true}}, \mathbf{l}_T^{\text{true}}) \le M_X$$

$$M_{T2} \equiv \min_{\mathbf{k}_T + \mathbf{l}_T = \mathbf{p}_T} \left[\max \left\{ M_T^{(1)}, M_T^{(2)} \right\} \right] \leq M_X$$

\blacklozenge $A_{\rm FB}$ at the Tevatron: 1/2

• $q\bar{q} \rightarrow t\bar{t}$ and $A_{\rm FB}$ at the Tevatron Shabamina(Blois2010)



 A_{fb} (raw) = 0.073 ± 0.028

$$A_{\rm FB} \equiv \frac{N_t(\cos\theta \ge 0) - N_{\bar{t}}(\cos\theta \ge 0)}{N_t(\cos\theta \ge 0) + N_{\bar{t}}(\cos\theta \ge 0)} = (0.15 \pm 0.05 \pm 0.024)$$

Pertinent $\sim 2\sigma$ deviation !

 \blacklozenge $A_{\rm FB}$ at the Tevatron: 2/2

• Effective Lagrangian approach to explain the anomaly D.-W. Jung, P. Ko, JSL, S.-h. Nam, PLB691(2010)238; D.-W. Jung, P. Ko, JSL, arXiv:1012.0102 [hep-ph]

$$\mathcal{L}_6 = \frac{g_s^2}{\Lambda^2} C_{8q}^{AB} \left(\bar{q}_A T^a \gamma_\mu q_A \right) \left(\bar{t}_B T^a \gamma^\mu t_B \right)$$



 $C_1 \equiv C_{8q}^{RR} + C_{8q}^{LL}, \quad C_2 \equiv C_{8q}^{LR} + C_{8q}^{RL}, \quad D, D_{\text{FB}} \propto (C_{8q}^{RR} - C_{8q}^{LL}) \pm (C_{8q}^{LR} - C_{8q}^{RL})$

🔶 Backup

• <u>CPX Scenario</u>:

Recall CPV mixing $\propto \Im m(A_f \mu)/M_{\rm SUSY}^2$

$$\begin{split} M_{\tilde{Q}_3} &= M_{\tilde{U}_3} = M_{\tilde{D}_3} = M_{\tilde{L}_3} = M_{\tilde{E}_3} = M_{\text{SUSY}} \\ |\mu| &= 4 \, M_{\text{SUSY}} \,, \quad |A_{t,b,\tau}| = 2 \, M_{\text{SUSY}} \\ |M_3| &= 1 \quad \text{TeV} \end{split}$$

Varying parameters are :

$$\tan \beta, \quad M_{H^{\pm}}^{\text{pole}}, \quad M_{\text{SUSY}}$$
$$(\Phi_{\mu}), \quad \Phi_{A_t}, \quad \Phi_{A_b}, \quad \Phi_{A_{\tau}}, \quad \Phi_3$$

* M_1 and M_2

* For simplicity, common Φ_A
🄶 Backup

• Large $\tan \beta$ and $M_{H^{\pm}}^{\text{pole}} \sim 150 \text{ GeV} \Rightarrow \text{Trimixing scenario}$ J. Ellis, JSL and A. Pilaftsis, PRD70(2004)075010, PRD71(2005)075007, PRD72(2005)095006, and NPB718(2005)247

$$\begin{aligned} &\tan \beta = 50, \quad M_{H^{\pm}}^{\text{pole}} = 155 \quad \text{GeV}, \\ &M_{\tilde{Q}_3} = M_{\tilde{U}_3} = M_{\tilde{D}_3} = M_{\tilde{L}_3} = M_{\tilde{E}_3} = 0.5 \quad \text{TeV}, \\ &|\mu| = 0.5 \quad \text{TeV}, \quad |A_{t,b,\tau}| = 1 \quad \text{TeV}, \quad |M_{1,2}| = 0.3 \quad \text{TeV}, |M_3| = 1 \quad \text{TeV}, \\ &\Phi_{\mu} = 0^{\circ}, \quad \Phi_{1,2} = 0^{\circ}. \end{aligned}$$

For $\Phi_{A_t,A_b,A_\tau} = 90^\circ$ and $\Phi_3 = -90^\circ$, we have (in GeV)

$$M_{H_1} = 117.6, \quad M_{H_2} = 118.7, \quad M_{H_3} = 122.6,$$

 $\Gamma_{H_1} = 1.48, \quad \Gamma_{H_2} = 8.12, \quad \Gamma_{H_3} = 8.21.$

All Higgs bosons are nearly degenerate with $\Gamma_i \gtrsim \Delta M_H$!

🔶 Backup

• SPS1a-like scenario

- At M_Z : Three gauge couplings $\alpha_1(M_Z)$, $\alpha_2(M_Z)$, and $\alpha_3(M_Z)$
- At m_t^{pole} : Quark and Lepton masses $m_{q,l}(m_t^{\text{pole}})$ and $V_{\text{CKM}}(m_t^{\text{pole}})$
- At M_{SUSY} : $\tan \beta(M_{\text{SUSY}})$
- At $M_{\rm MFV} = M_{\rm GUT}$: 19 MCPMFV Parameters

$$|M_{1,2,3}| e^{i \Phi_{1,2,3}}, \quad |A_{u,d,e}| e^{i \Phi_{A_{u,d,e}}}, \quad \widetilde{M}^2_{Q,U,D,L,E}, \quad M^2_{H_{u,d}}$$

Specifically, we have taken the parameter set:

$$\begin{split} |M_{1,2,3}| &= 250 \text{ GeV} \\ M_{H_u}^2 &= M_{H_d}^2 = \widetilde{M}_Q^2 = \widetilde{M}_U^2 = \widetilde{M}_D^2 = \widetilde{M}_L^2 = \widetilde{M}_E^2 = (100 \text{ GeV})^2 \\ |A_u| &= |A_d| = |A_e| = 100 \text{ GeV} \end{split}$$

This parameter set is equivalent to SPS1a when $\Phi_{A_{u,d,e}} = 180^{\circ}$ and $\Phi_{1,2,3} = 0^{\circ}$ if $M_{SUSY} = m_t^{pole}$ and $\tan \beta = 10$

🄶 Backup

• The *B*-meson observables strike against CPX!: $B(B_s \to \mu^+\mu^-)$ (95 %), $B(B \to X_s\gamma)$ (2 σ), $R_{B\tau\nu}$ (1 σ), $\Upsilon(1S) \to H_1\gamma$ CPX scenario with $\Phi_A = \Phi_3 = 90^\circ$ and $M_{SUSY} = 0.5$ TeV JSL, M. Carena, J. Ellis, A. Pilaftsis, C.E.M. Wagner, arXiv:0712.2360; JSL and S. Scopel, PRD75(2007)075001

