Particle Creation in Charged Black Holes

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Outline

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- Particle Creation in Charged Black Holes
 - Boundary Conditions
 - Near Extremal Reissner-Nordström Black Holes

Conclusion

Overview: Spontaneous Pair Production

- Quantum Vacuum Fluctuation: virtual particles
 - Heisenberg's uncertainty principle: $\Delta E \Delta t \geq \hbar/2$
 - creation of particle-antiparticle pairs (virtual particles)
- Spontaneous Pair Production: from virtual to real
 - Schwinger mechanism: electric field

Schwinger, 1951

Hawking radiation: causal horizon (tunneling)

Parikh, Wilczek, [hep-th/9907001]





- Black Hole Thermodynamics
 - Hawking temperature: surface gravity κ

$$T = \frac{\hbar c^3}{kG} \frac{\kappa}{2\pi}$$

Beckenstein-Hawking entropy: area of horizon A

$$S = \frac{kc^3}{\hbar G} \frac{A}{4}$$

Holographic Principle

t' Hooft, 1993; Susskind, 1994

Gravity in Bulk *D* dimensions



Field Theory on Boundary D-1 dimensions

- Preliminary hint: Symmetry
 - Gravity: anti de Sitter (AdS) space appears in near horizon geometry of extremal black holes
 - Field Theory: Conformal Field Theory (CFT)

Symmetry group: SO(D-1,2)for both AdS_D and CFT_{D-1}



AdS/CFT Correspondence: (string theory)

Maldacena, 1997



$$N = 4$$
 Super Yang-Mills
on Boundary of AdS₅

 Kerr/CFT Correspondence (extremal) Guica, Hartman, Song, Strominger, arXiv:0809.4266 [hep-th] Compere, arXiv:1203.3561 [hep-th] (review)
 Near horizon geometry: warped AdS₃ SL(2, R) × U(1)
 (CFT) Central charge: c_L = c_R = 12J/ħ
 (CFT) Temperature: T_L = 1/2π; T_R = 0



Reissner-Nordström (RN)/CFT Correspondence (extremal) Hartman, Murata, Nishioka, Strominger: arXiv:0811.4393 [hep-th] Garousi, Ghodsi: arXiv:0902.4387 [hep-th] CMC, Huang, Zou: arXiv:1001.2833 [hep-th] CMC, Sun: arXiv:1106.4407 [hep-th] (summary)

• Near horizon geometry: $AdS_2 \times S^2$

■ Warped AdS₃/CFT₂ description:

The U(1) bundle of warped AdS₃ was recovered from the gauge field potential by uplifting the RN black hole into 5D.

- The period of extra coordinate is $y \sim y + 2\pi \ell$
- Gravitational constant: $G_5 = 2\pi \ell G_4$
- (CFT) Central charge: $c_L = c_R = 6Q^3/\ell$
- (CFT) Temperature: $T_L = \ell/2\pi Q; \quad T_R = 0$

- Kerr-Newman (KN)/CFTs Correspondence (extremal) CMC, Huang, Sun, Wu, Zou: arXiv:arXiv:1006.4097 [hep-th] CMC, Sun: arXiv:1201.4040 [hep-th] (summary)
- There are two different individual 2D CFTs holographically dual to the KN black hole.



microscopic hair theorem: For each "hair" of black hole, except the mass, there is an associated CFT discsrption.

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Overview: Motivation

Purpose: particle creation in charged black holes

- technical simplicity: constant electric field (exactly solvable)
- holographic description: anti de Sitter
- Reissner-Nordström (RN) Black Holes: near extremal

near horizon region: where the production occurs

 $\mathrm{A}dS_2\times S^2+\mathrm{constant}$ electric field

extremal RN

vanishing HAWKING temperature: stable (thermal)
 non-vanishing electric field: unstable (quantum)

CMC, S.P. Kim, Y.-J. Lin, J.-R. Sun, M.-F. Wu, arXiv:1202.3224 [hep-th]

Near Horizon Geometry of RN Back Holes

 Near horizon geometry of near extremal Reissner-Nordström Black Holes:



Near horizon geometry

- Horizon radius: $r_{\pm} = M \pm \sqrt{M^2 - Q^2}$
- Near extremal: $Q \rightarrow M \Rightarrow r_{-} \rightarrow r_{+}$

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Particle Creation: Boundary Conditions

- Pair Production: probe massive charged scalar
 - The ratios of fluxes exhibit particle (scalar) creation with suitable boundary conditions.

Kim, Page, [arXiv:hep-th/0005078]

Kim, Lee, Yoon, [arXiv:0910.3363 [hep-th]]

Boundary condition I: particle view point



Particle Creation: Boundary Conditions

Boundary condition II: antiparticle view point



- incident: virtual particles
- reflected: re-annihilated
- transmitted: pair produced "antiparticles"

Equivalence:

particles and antiparticles should always appear in pairs

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Particle Creation: Physical Quantities

• vacuum persistence amplitude: (probability) $|\alpha|^2$ (* physically it should be $1/|\alpha|^2$ *)

$$|\alpha|^2 \equiv \frac{D_{\text{incident}}}{D_{\text{reflected}}}, \qquad \frac{1}{|\alpha|^2} \equiv \frac{D_{\text{reflected}}}{D_{\text{incident}}}$$

 $|\beta|^2$

mean number of produced pairs:

$$|\beta|^2 \equiv \frac{D_{\text{transmitted}}}{D_{\text{reflected}}}$$

absorption cross section: (probability) σ_{abs}

$$\sigma_{\rm abs} \equiv \frac{D_{\rm transmitted}}{D_{\rm incident}} = \frac{|\beta|^2}{|\alpha|^2}$$

flux conservation and Bogoliubov relation

 $|D_{\text{incident}}| = |D_{\text{reflected}}| + |D_{\text{transmitted}}| \quad \iff |\alpha|^2 - |\beta|^2 = 1$

Particle Creation: Physical Quantities

Particular values of physical quantities:

• Normalization: $D_{\text{incident}} = 1$

Production	$D_{\rm reflected}$	$D_{\mathrm{transmitted}}$	$ \alpha ^{-2}$	$ \beta ^2$	$\sigma_{ m abs}$
None	1	0	1	0	0
Half	1/2	1/2	1/2	1	1/2
Full	0	1	0	∞	1

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Particle Creation: RN Back Holes

Near horizon geometry of near-extremal RN

$$ds^{2} = -\frac{\rho^{2} - B^{2}}{Q^{2}}d\tau^{2} + \frac{Q^{2}}{\rho^{2} - B^{2}}d\rho^{2} + Q^{2}d\Omega_{2}^{2}$$
$$A_{[1]} = -\frac{\rho}{Q}d\tau; \qquad F_{[2]} = \frac{1}{Q}d\tau \wedge d\rho$$

"rescaled" deviation from extremality: B \propto M - Q
 geometric structure: AdS₂ × S² (radius Q for both)
 electric field: constant

probe massive charged scalar: Φ

$$S_{\Phi} = \int d^4 x \sqrt{-g} \left(-\frac{1}{2} D_{\alpha} \Phi^* D^{\alpha} \Phi - \frac{1}{2} m^2 \Phi^* \Phi \right)$$

• *m* and *q* are the mass and charge of Φ

$$D_{\alpha} \equiv \nabla_{\alpha} - iqA_{\alpha}$$

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Particle Creation: RN Back Holes

Field equation: Klein-Gordon (KG) equation

$$(
abla_{lpha} - iqA_{lpha})(
abla^{lpha} - iqA^{lpha})\Phi - m^{2}\Phi = 0$$

Flux:

$$D = i\sqrt{-g}g^{
ho
ho}(\Phi D_{
ho}\Phi^* - \Phi^*D_{
ho}\Phi)$$

Ansatz:

$$\Phi(\tau, \rho, \theta, \phi) = e^{-i\omega\tau + in\phi} R(\rho) S(\theta)$$

separated field equations (exactly solvable !!!)

$$\partial_{\rho} \left[(\rho^2 - B^2) \partial_{\rho} R \right] + \left[\frac{(q\rho - \omega Q)^2 Q^2}{\rho^2 - B^2} - m^2 Q^2 - \lambda_I \right] R = 0$$

$$\frac{1}{\sin \theta} \partial_{\theta} (\sin \theta \partial_{\theta} S) - \left(\frac{n^2}{\sin^2 \theta} - \lambda_I \right) S = 0$$

• $S(\theta)$ is spherical harmonics with the eigenvalue $\lambda_I = I(I+1)_{\text{respective}} \otimes \mathbb{Q}(\theta)$

Particle Creation: RN Back Holes

Condition for Schwinger mechanism and/or Hawking radiation

$$(m^2 - q^2)Q^2 + (l + 1/2)^2 < 0$$

violation of Breitenlohner-Freedman (BF) bound in AdS₂
 unstable mode

Cosmic censorship: necessary condition

$$q^2 > m^2$$

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avoiding naked singularity

Schwinger mechanism: extremal to non-extremal

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Particle Creation: Results

Bogoliubov coefficients

$$|\alpha|^{2} = \frac{\cosh(\pi a - \pi b)\cosh(\pi \tilde{a} + \pi b)}{\cosh(\pi a + \pi b)\cosh(\pi \tilde{a} - \pi b)}$$
$$|\beta|^{2} = \frac{\sinh(2\pi b)\sinh(\pi \tilde{a} - \pi a)}{\cosh(\pi a + \pi b)\cosh(\pi \tilde{a} - \pi b)}$$
$$\equiv qQ, \quad b \equiv \sqrt{(q^{2} - m^{2})Q^{2} - (l + 1/2)^{2}}, \quad \tilde{a} \equiv \frac{\omega Q^{2}}{B}$$

Absorption cross section:

$$\sigma_{
m abs} = rac{\sinh(2\pi b)\sinh(\pi ilde{a} - \pi a)}{\cosh(\pi a - \pi b)\cosh(\pi ilde{a} + \pi b)}$$

 \blacksquare Leading term of $|\beta|^2$ leads to the Schwinger formula

$$|\beta|^2 \approx e^{-\frac{\pi m^2 Q}{q}} \approx e^{-\frac{\pi m^2 r_H^2}{qQ}}$$

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Particle Creation: Results

An interesting observation:

$$\frac{|\beta(B=0)|^2}{|\beta(B\neq 0)|^2} = \frac{\cosh(\pi\tilde{a}-\pi b)}{\sinh(\pi\tilde{a}-\pi a)}e^{\pi b-\pi a} = \frac{1+e^{2\pi(b-\tilde{a})}}{1-e^{2\pi(a-\tilde{a})}} \ge 1$$

Production rate:

extremal (Schwinger) \geq near extremal (Schwinger+Hawking)

From the extremal to near extremal black holes

- Increasing attractive gravitational force reduces electromagnetic repulsive force for Schwinger mechanism.
- Schwinger mechanism is suppressed faster than the increasing part from the Hawking thermal radiation.
- Such kind of interaction generically prohibits to distinguish the Schwinger mechanism from the Hawking radiation.

Particle Creation: Holographic Description

CFT absorption cross section: standard formula

$$\sigma_{abs} \sim \frac{(2\pi T_L)^{2h_L-1}}{\Gamma(2h_L)} \frac{(2\pi T_R)^{2h_R-1}}{\Gamma(2h_R)} \sinh\left(\frac{\omega_L - q_L\Omega_L}{2T_L} + \frac{\omega_R - q_R\Omega_R}{2T_R}\right) \\ \times \left|\Gamma\left(h_L + i\frac{\omega_L - q_L\Omega_L}{2\pi T_L}\right)\right|^2 \left|\Gamma\left(h_R + i\frac{\omega_R - q_R\Omega_R}{2\pi T_R}\right)\right|^2$$

CFT dual to RN black holes

$$c_L = c_R = rac{6Q^3}{\ell}, \quad T_L = rac{\ell}{2\pi Q}, \quad T_R = rac{\ell B}{\pi Q^2}$$

 free parameter ℓ is related to measure of U(1) bundle. CMC, Y.-M. Huang, S.-J. Zou, arXiv:1001.2833 [hep-th]
 Conformal weights of dual operator: h_L = h_R = 1/2 ± ib complex conformal weight ⇒ unstable dual operator

Particle Creation: Holographic Description

Identification by thermodynamics:

$$\frac{\delta M}{T_H} - \frac{\Omega_H \delta Q}{T_H} = \frac{\omega_L - q_L \Omega_L}{T_L} + \frac{\omega_R - q_R \Omega_R}{T_R}$$

- Hawking temperature: $T_H = \frac{B}{2\pi Q^2}$
- chemical potential: $\Omega_H = A_\tau(B) = -B/Q$

•
$$\delta M = \omega$$
 and $\delta Q = -q$

$$\tilde{\omega}_L \equiv \omega_L - q_L \Omega_L = -q\ell \quad \text{and} \quad \tilde{\omega}_R \equiv \omega_R - q_R \Omega_R = 2\omega\ell$$

 The absorption cross section agrees with the CFT's result only up to some numerical factors.

Conclusion

- We discuss the spontaneous pair production of charged scalar field in RN black holes (both extremal and near extremal limits).
- This is one of the few examples that we exactly know the Bogoliubov coefficients.
- The pair production (unstable mode) is holographically dual to an operator with complex conformal weight.
- The production rate is suppressed when the black hole temperature is turned on. (It should be enhanced in dS space.)
- The effects of the Schwinger mechanism and the Hawking radiation generically cannot be distinguished by imposing different boundary conditions.

Conclusion

Phase Diagram:

