

# Implication of 126 GeV Higgs for Planck scale physics

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Higgs was discovered at  
 $M_H = 126 \text{ GeV}$

No evidence  
of “new” physics  
@ ATLAS, CMS & LHCb

**What is the implication of these two?**

together with some phenomena beyond SM  
( $\nu$  oscillation, Baryon asymmetry, Dark matter)

# 3 major hints towards the physics beyond SM

(1) Higgs mass at 126 GeV (Moriond)

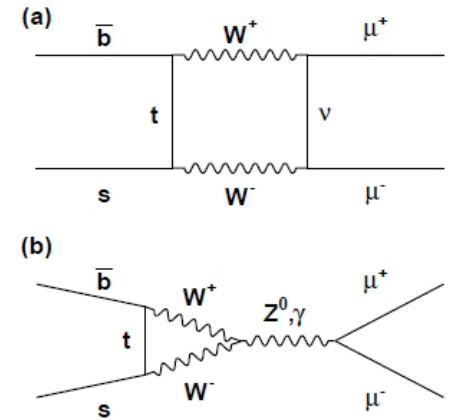
(2) No deviation from SM

Higgs decay

Flavor physics: LHCb, B-factory, MEG

$$\text{e.g. } \text{Br}(B_s \rightarrow \mu\mu) = 3.2^{+1.5}_{-1.2} \times 10^{-9}$$

which is consistent with SM  $(3.2 \pm 0.2) \times 10^{-9}$ .



(3) But we also know that **SM is not sufficient** to explain  
neutrino oscillation  
Baryon asymmetry  
dark matter

(also a big hint from **Cosmology**)

What can (or must) we do now  
to go beyond the SM?

Is there a possibility to connect  
with the string theory  
or Planck scale physics?

Most investigations of physics beyond the SM  
have been based on

“the central dogma” of particle physics  
GUT → hierarchy problem → TeV SUSY etc.

i.e. Unification below the Planck scale requires  
large symmetry enhancement at TeV scale.

It may be a good time to reconsider the basic  
strategy (central dogma) toward the physics BSM.

# Hierarchy problem

## Naturalness (Hierarchy problem)

$$\begin{aligned}\delta V(\phi) &= \frac{1}{2} \int \frac{d^4 k}{(2\pi)^4} \text{Str} \log(k^2 + M^2(\phi)) && \text{Quadratic divergence} \\ &= \frac{\Lambda^2}{32\pi^2} \text{STr} M^2(\phi) + \text{STr} \frac{M^4(\phi)}{64\pi^2} (\ln(M^2/\Lambda^2) - 1/2)\end{aligned}$$

$\text{STr} M^2(\phi) \neq 0$       Quadratic divergence in Higgs mass term

$\text{STr} M^2(\phi) = 0$       Cancellation of Quadratic divergence  
(**supersymmetry** etc.)

Bardeen (1995 @ Ontake summer school)

Standard model is **classically scale invariant** if Higgs mass term is absent.

$$T_{\mu}^{\mu} = 0$$

Quantum anomaly breaks the invariance (if not conformal )

$$T_{\mu}^{\mu} = \beta(\lambda_i) \mathcal{O}_i$$

But then, we cannot forbid the quadratically divergent mass term

$$T_{\mu}^{\mu} = \beta(\lambda_i) \mathcal{O}_i + \text{const. } \Lambda^2 \bar{h}h$$

**Bardeen argued that it should be**

$$T_{\mu}^{\mu} = \beta(\lambda_i) \mathcal{O}_i + \delta m^2 \bar{h}h$$

$$\delta m^2 = \text{const.} \times m^2 \neq \text{const.} \times \Lambda^2$$

# Hierarchy problem

## Is quadratic divergence the issue of hierarchy problem?

NO

Bardeen(1995)

H Aoki, SI PRD(2012)

There are 3 different types of divergences

1. Quadratic divergences  $\Lambda^2$
2. Logarithmic divergences  $m^2 \log \Lambda$
3. Logarithmic but looks like quadratic  $M^2 \log \Lambda$



(1) Quadratic divergence in mass can be simply subtracted, so it gives a **boundary condition** at UV cut off  $\Lambda$ .

→ If massless at  $\Lambda$ , it continues to be so in the IR theory.

(2) Logarithmic divergence gives a multiplicative renormalization.  
**No Higgs mass term is generated** if it is absent at UV scale.

$$\frac{dm^2}{dt} = \frac{m^2}{16\pi^2} \left( 12\lambda + 6Y_t^2 - \frac{9}{2}g^2 - \frac{3}{2}g_1^2 \right) + \boxed{\frac{M^2}{8\pi^2} \lambda_{mix}}$$

(3) If SM is coupled with a massive particle with mass  $M$ ,  
**logarithmic divergences give a correction to  $m$**  as

$$\delta m^2 = \frac{\lambda_{mix} M^2}{16\pi^2} \log(\Lambda^2/m^2)$$

In order to solve the “hierarchy problem” without a special cancellation like supersymmetry, we need to control

(a) “quadratic divergence”  $\rightarrow$  correct **boundary condition** at Planck

The most natural b.c. is **NO MASS TERMS at Planck**  
( = classical conformal invariance)

(b) “large logarithmic divergence” by mixing with a large mass  $M$   
**No intermediate scales** between EW (or TeV) and Planck

**“Classical conformal theory with no intermediate scale”**  
can be an alternative solution to hierarchy problem.

Bardeen (95)

Shaposhnikov (07)

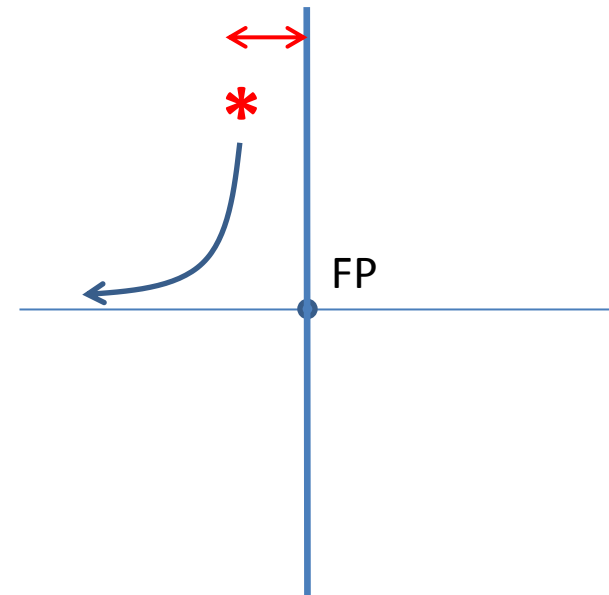
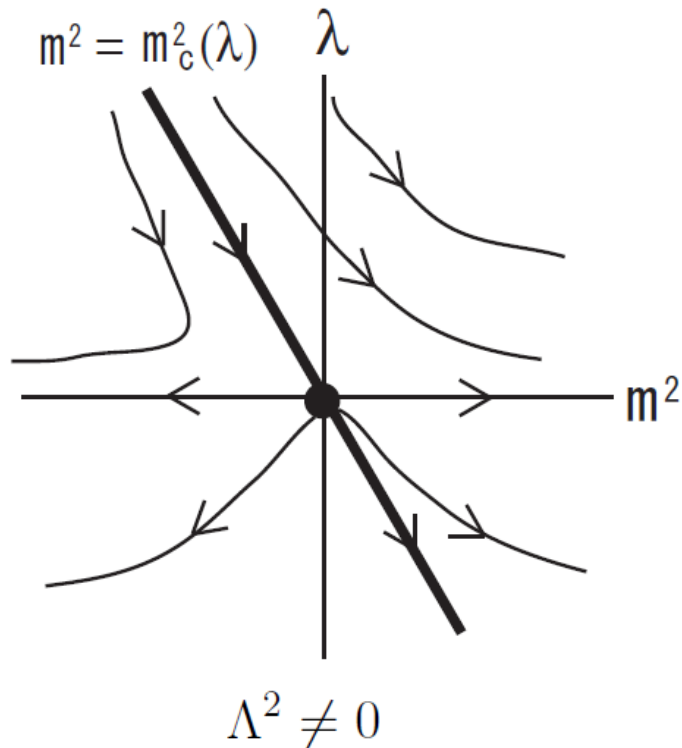
Meissner Nicolai (07)

SI, Okada, Orikasa (09)<sup>10</sup>

# Hierarchy problem in Wilsonian RG

H Aoki, SI PRD(2012)]

## Critical line

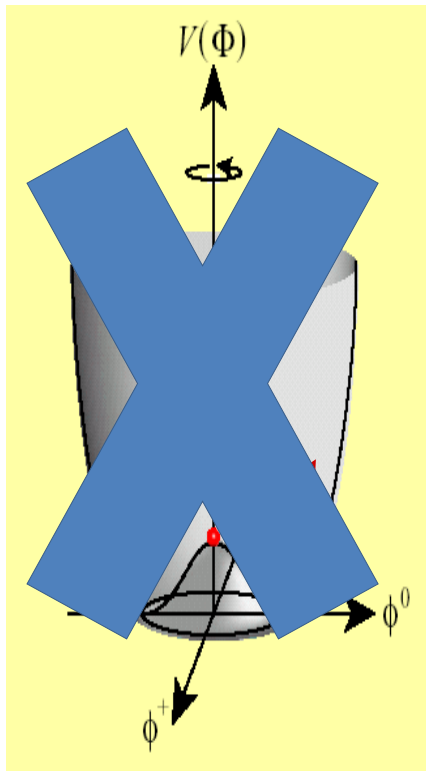


Fine - tuning of the **distance from the critical line** = **Low energy mass scale**

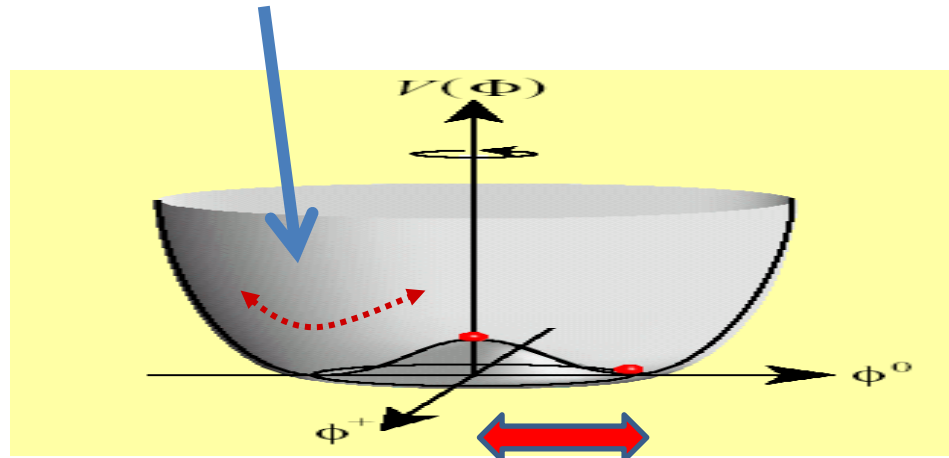
The difference is the **choice of the coordinates** of the parameter space.

# Stability of Vacuum

Another Hint of 126 GeV Higgs mass is  
**Stability bound** of the Higgs quartic coupling



$$m_H = 126 \text{ GeV}$$



$$v = 246 \text{ GeV}$$

RGE improved effective potential for large field ( $h \gg v$ )  $V_{\text{eff}}(h) = \frac{\lambda_{\text{eff}}(h)}{4} h^4$

$$\text{RGE @1-loop} \quad \frac{d\lambda_H}{dt} = \frac{1}{16\pi} \left( 24\lambda^2 - \underbrace{6Y_t^4 + \frac{9}{8}g^4 + \frac{3}{8}g_1^4 + \dots}_{\text{Already known}} \right)$$

It is related to Higgs mass as  $M_h^2 = 2\lambda v^2$

Higgs mass controls the behavior of Higgs potential at large values of  $h$ .

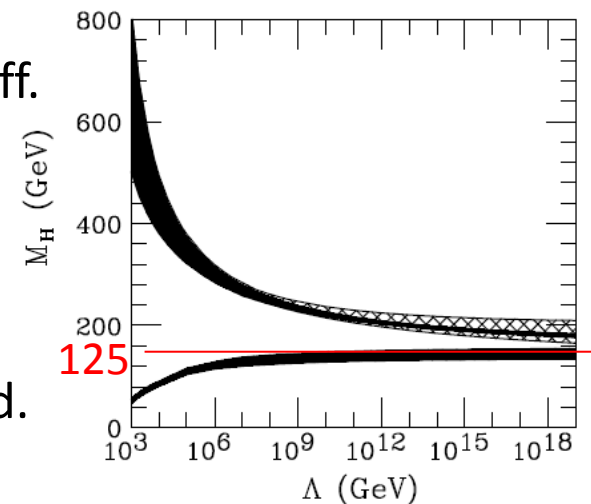
This gives two bounds for Higgs mass

(1) The quartic coupling does not blow up until UV cut-off.

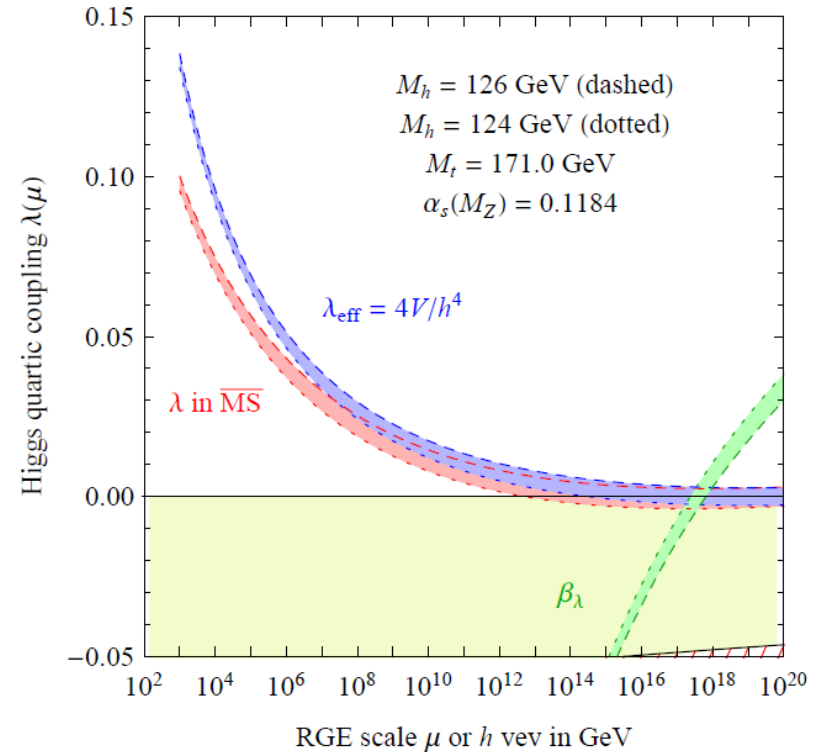
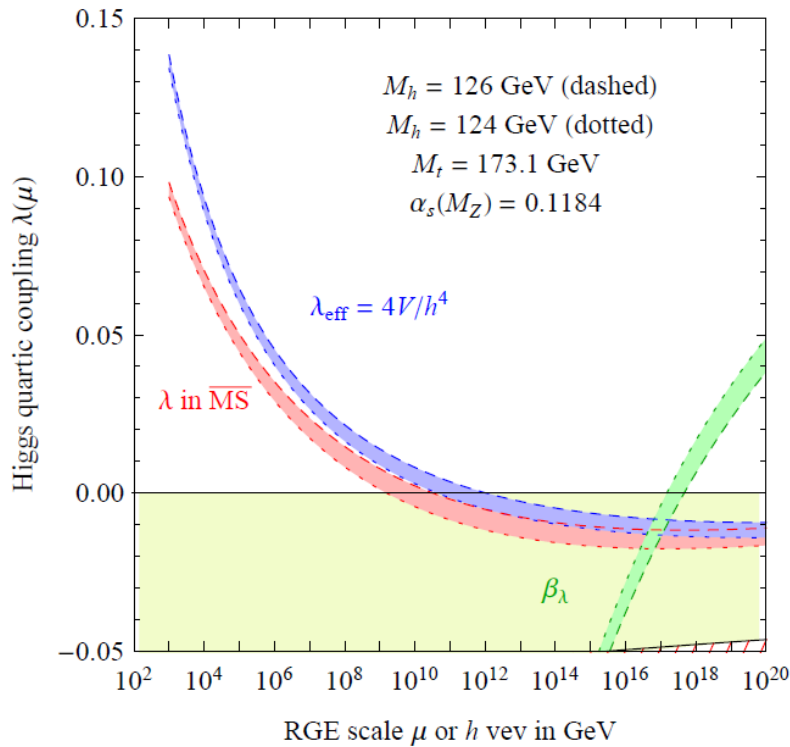
$M < 180 \text{ GeV}$  (triviality bound)

(2) The quartic coupling does not become negative until UV cut-off. (Stability bound)

$M = 125 \text{ GeV}$  Higgs is very close to the stability bound.



# Why stability bound is important for Planck scale physics?



**New physics at  $10^{12}$  GeV**  
is necessary to stabilize the vacuum

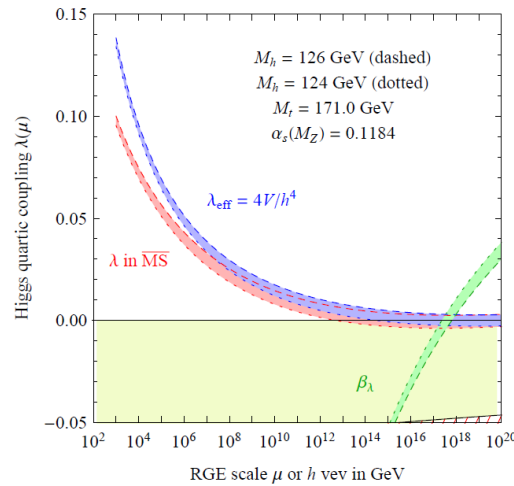
**Flat Higgs potential**  
at Planck scale

$$M_H \geq 129.2 + 1.8 \times \left( \frac{m_t^{\text{pole}} - 173.2 \text{ GeV}}{0.9 \text{ GeV}} \right) - 0.5 \times \left( \frac{\alpha_s(M_Z) - 0.1184}{0.0007} \right) \pm 1.0 \text{ GeV}.$$

**very sensitive to top quark mass**

Elias-Miro et.al.(12)  
Alkhin, Djouadi, Moch (12)

If this



is the case ?

$$\lambda(\Lambda_0) = \beta_\lambda(\Lambda_0) = 0$$

Direct window to Planck scale

M.Shaposhnikov (07)



# Emergence of Higgs potential

at the Planck scale

# Indication on the Higgs potential

$$V = -\mu^2 |H|^2 + \lambda (|H|^2)^2$$

Hierarchy  
(classical conformality)

Stability  
vanish at Planck

LHC data implies that

Higgs has a **flat potential**  $V(H)=0$  at Planck.

How can we achieve **EW symmetry breaking**  
from  $V(H)=0$  potential at Planck?

Everything should be  
radiatively generated.

# Two mechanisms of symmetry breaking

(1) SSB by negative mass term

$$V = \frac{\lambda}{4}h^4 + \frac{\mu^2}{2}h^2 \quad (\mu^2 < 0) \longrightarrow m_h^2 = 2|\mu^2| = 2\lambda\langle h \rangle^2$$

(2) Coleman-Weinberg mechanism (radiative breaking)

$$V_{eff} = \underbrace{\frac{\lambda h^4}{4}}_{\text{tree}} + \underbrace{Bh^4 \left( \ln \left( \frac{h^2}{\langle h \rangle^2} \right) - \frac{25}{6} \right)}_{\text{1-loop}} \quad B = \frac{3}{64\pi^2} \left( 3\lambda^2 + \frac{3g^4 + 2g^2g'^2 + g'^4}{16} - Y_t^4 \right)$$

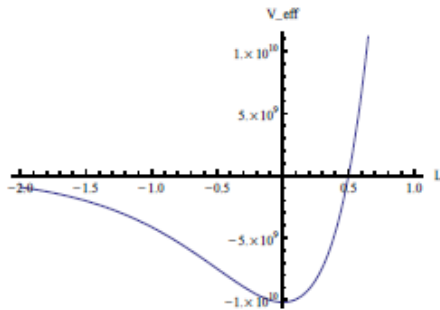


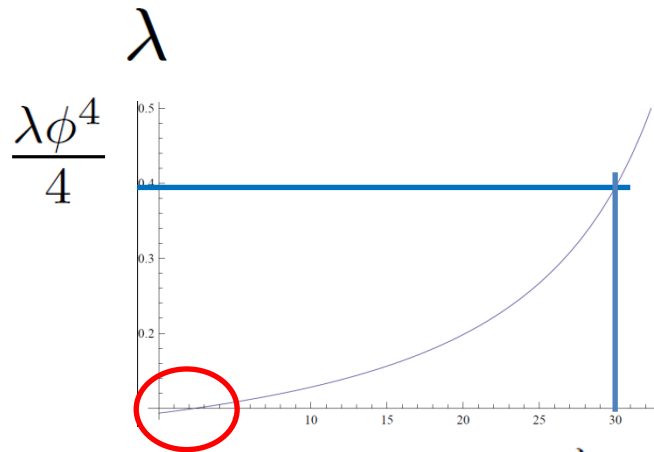
図 2.1 improve されていない素数がテンシヤル

Higgs mass is given by

$$V''|_{\langle h \rangle} = m_h^2 = 8B\langle h \rangle^2 = \frac{6}{11}\lambda\langle h \rangle^2$$

## (2') RG improved CW mechanisms

### Coleman-Weinberg radiative breaking



Symmetry is broken near the scale where the running coupling crosses zero.

$$M_{CW} = M_{UV} \exp\left(-\frac{\lambda_{UV}}{b} - \frac{1}{4}\right) \quad \text{Positive beta function}$$

Dimensional transmutation

cf. Dimensional transmutation in QCD

$$\Lambda_{QCD} = M_{UV} \exp\left(-\frac{2\pi}{b_0\alpha_s(M_{UV})}\right)$$

But CW does not work in SM.

the large top Yukawa coupling invalidates the CW mechanism



Extension of SM is necessary !

Meissner Nicolai (07)

(B-L) extension of SM with flat Higgs potential at Planck

SM



B-L sector

- U(1)B-L gauge
- SM singlet scalar  $\phi$
- Right-handed  $\nu$

N Okada, Y Orikasa,  
& SI

0902.4050 (PLB)

0909.0128 (PRD)

1011.4769 (PRD)

1210.2848(PTEP)

“Occam’s razor” scenario

that can explain

- 126 GeV Higgs
- hierarchy problem
- $\nu$  oscillation, baryon asymmetry

# Model: (B-L) extension of SM with Right Handed Neutrinos

	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)_{B-L}$
$q_L^i$	<b>3</b>	<b>2</b>	$+1/6$	$+1/3$
$u_R^i$	<b>3</b>	<b>1</b>	$+2/3$	$+1/3$
$d_R^i$	<b>3</b>	<b>1</b>	$-1/3$	$+1/3$
$\ell_L^i$	<b>1</b>	<b>2</b>	$-1/2$	$-1$
$\nu_R^i$	<b>1</b>	<b>1</b>	$0$	$-1$
$e_R^i$	<b>1</b>	<b>1</b>	$-1$	$-1$
$H$	<b>1</b>	<b>2</b>	$-1/2$	$0$
$\Phi$	<b>1</b>	<b>1</b>	$0$	$+2$

N Okada, Y Orikasa, SI  
PLB676(09)81,  
PRD80(09)115007  
PRD83(11)093011

$H$  Higgs doublet  
 $\Phi$  B-L sector scalar field

- **B-L** is the only anomaly free global symmetry in SM.
- $[U(1)_{B-L}]^3$  is anomaly free if we have **right handed fermion**.
- B-L gauge symmetry is broken by vev of an **additional scalar**.

See-saw mechanism

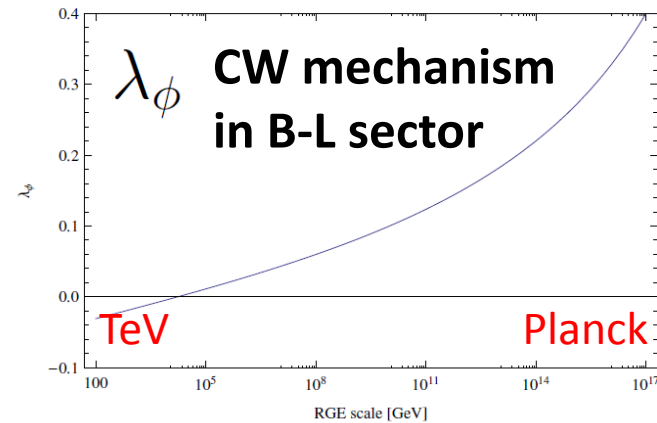
$$\mathcal{L} \supset -Y_D^{ij} \bar{\nu}_R^i H^\dagger \ell_L^j - \frac{1}{2} Y_N^i \Phi \bar{\nu}_R^i \nu_R^i + \text{h.c.}, \quad \begin{pmatrix} 0 & m \\ m & M_N \end{pmatrix} \longrightarrow m_\nu = \frac{m^2}{M_N}$$

$$m = Y_D \langle H \rangle \quad M_N = Y_N \langle \phi \rangle$$

B-L can be broken by CW mechanism at TeV.

$$V(\Phi, H)|_{UV} = \lambda_{\Phi}(\Phi^{\dagger}\Phi)^2$$

$$M_{B-L} \sim M_{Planck} \exp\left(-\frac{\lambda_{\phi}}{b}\right)$$





How about EWSB ?

$$\cancel{m_H^2 H^2} + \cancel{\lambda_H H^4} + \cancel{\lambda_{H\Phi} H^2 \Phi^2}$$

classically  
conformal

126 GeV



key to relate EW and TeV

Can the small scalar mixing be naturally generated?



YES

Radiatively generated scalar mixing  $\lambda_{\text{mix}}$  in  $V(H)$

$$V(H) = \lambda_H H^4 + \lambda_{\text{mix}} \Phi^2 H^2 \rightarrow \text{EWSB}$$

# Why is small negative $\lambda_{\text{mix}}$ generated?

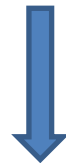
U(1) gauge mixing



$$\frac{dg_{\text{mix}}}{dt} \sim \frac{2}{3\pi^2} g_{B-L} g_Y^2 + g_{\text{mix}} \left( \text{*****} \right)$$



$$\frac{d\lambda_{\text{mix}}}{dt} = \frac{1}{16\pi^2} \left[ \lambda_{\text{mix}} \left( \text{*****} \right) + 12g_{\text{mix}}^2 g_{B-L}^2 \right]$$

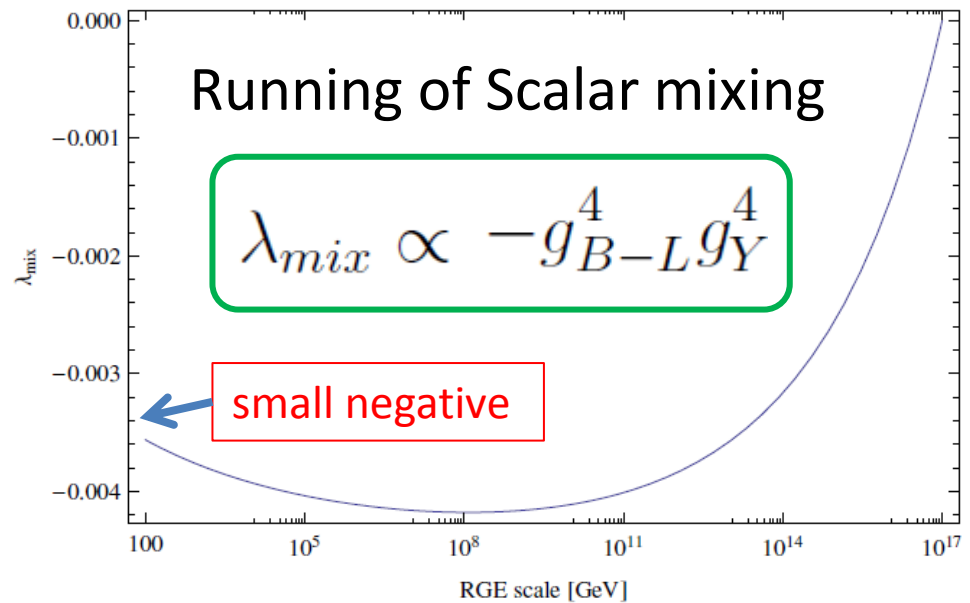


$$\lambda_{\text{mix}} \propto -g_{\text{mix}}^2 g_{B-L}^2 \propto -g_{B-L}^4 g_Y^4$$

# Radiative generation of scalar mixing

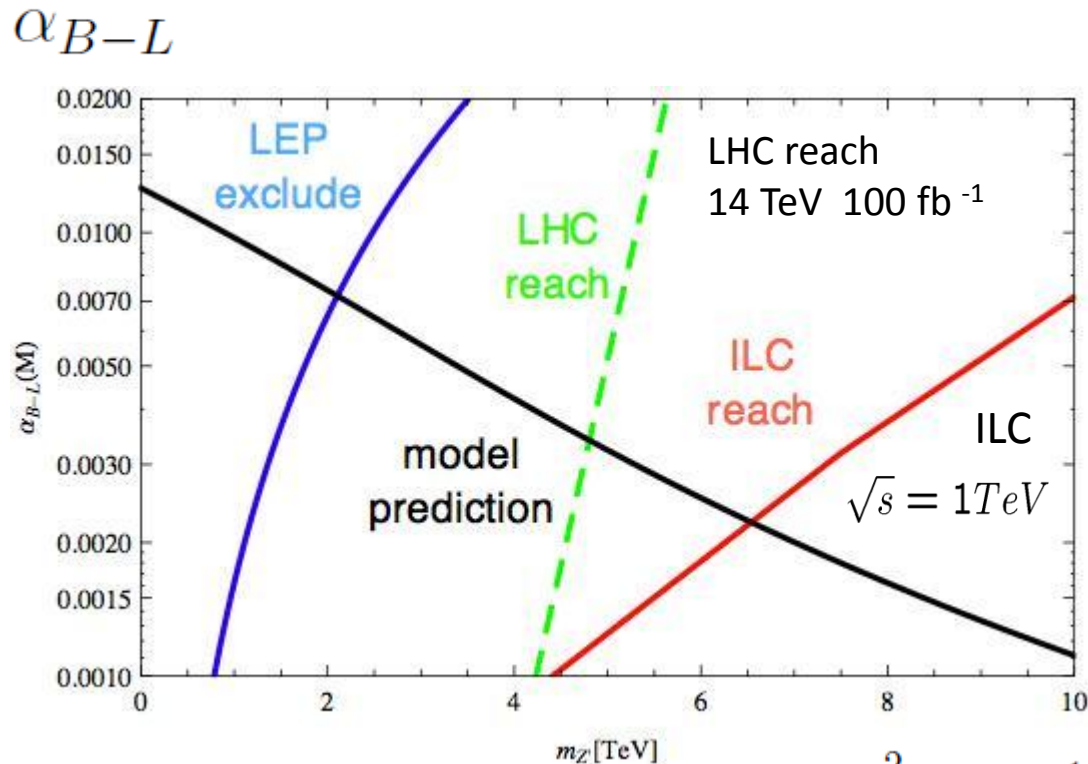
$$V(H) = \lambda_H H^4 + \lambda_{mix} \Phi^2 H^2$$

$$\begin{aligned}\langle H \rangle &= \sqrt{\frac{-\lambda_{mix}}{\lambda_H}} M_{B-L} \\ &\sim c \frac{\alpha_{B-L} \alpha_Y}{\sqrt{\lambda_H}} M_{B-L}\end{aligned}$$



# Prediction of the model

In order to realize **EWSB at 246 GeV**,  
B-L scale must be around TeV (for a typical value of  $\alpha_{B-L}$ ).



$$M_{B-L} \sim \frac{1}{\alpha_{B-L}} \times 35 \text{ GeV}$$

$$m_{Z'} \sim \frac{1}{\sqrt{\alpha_{B-L}}} \times 250 \text{ GeV}$$

$$m_\phi \sim 0.1 m_{Z'}$$

$$m_{Z'}^2 = 16\pi\alpha_{B-L}(0)M_{B-L}^2$$

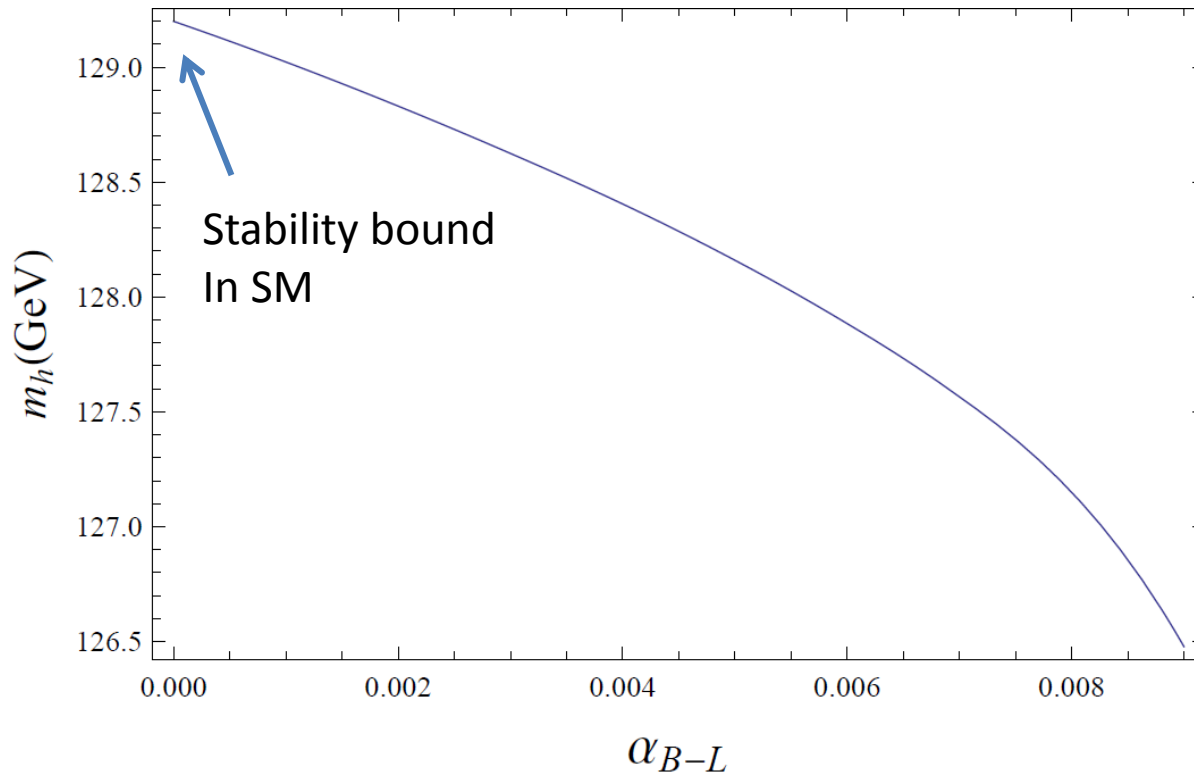
## Stability bound in TeV scale B-L model

$$\frac{d\lambda_H}{dt} = \frac{1}{16\pi} \left( 24\lambda^2 - 6Y_t^4 + \frac{9}{8}g^4 + \frac{3}{8}g_1^4 + \frac{3}{4}g^2g_1^2 + \boxed{\frac{3}{4}(g^2 + g_1^2)g_{mix}^2} + \dots \right)$$

An extra positive term is added



Lower the  
stability bound



# Summary

- 126 GeV Higgs = **border of the stability bound** of SM vacuum.
  - Direct window to Planck scale → **Flat Higgs potential @Planck**
  - Hint for the origin of Higgs in string theory
- **Occam's razor scenario beyond SM**
  - “Classically conformal B-L model” is proposed

(1) it can solve hierarchy problem

(2) it can explain why B-L breaking scale is around TeV.

(3) Stability bound can be lowered about 1 GeV

$$M_H \sim 128 \text{ GeV}$$

(4) phenomenologically viable

**Neutrino oscillation, resonant leptogenesis**

(5) Highly predictive (or excludable)

Prediction

$Z'$  around several TeV,  $M_\phi < M_{Z'}$ , Leptogenesis at TeV

## Ongoing and Future projects

- **Origin of flat Higgs potential at Planck**

Hierarchy problem &  $M_H = 126 \text{ GeV}$

→ PNGB ? Moduli ? Gauge/Higgs ? .....

Non-susy vacua of superstring with flat  $V(H)$

- **Resonant leptogenesis**

Garny, Kartavsev, Hohenegger (11)

Kadanoff-Baym equation (quantum Boltzman)

- **Non-susy GUT at Planck scale**

$SO(10)$  or  $E_6$  type or Pati-Salam

Gravity or string threshold correction to RGE

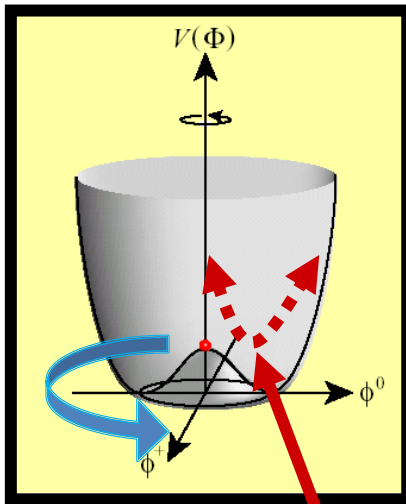




# Higgs potential

$$V = -\mu^2 |H|^2 + \lambda (|H|^2)^2 \longrightarrow |H| = \sqrt{\frac{\lambda}{2\mu^2}} \equiv \frac{v}{\sqrt{2}}$$

Potential minimum



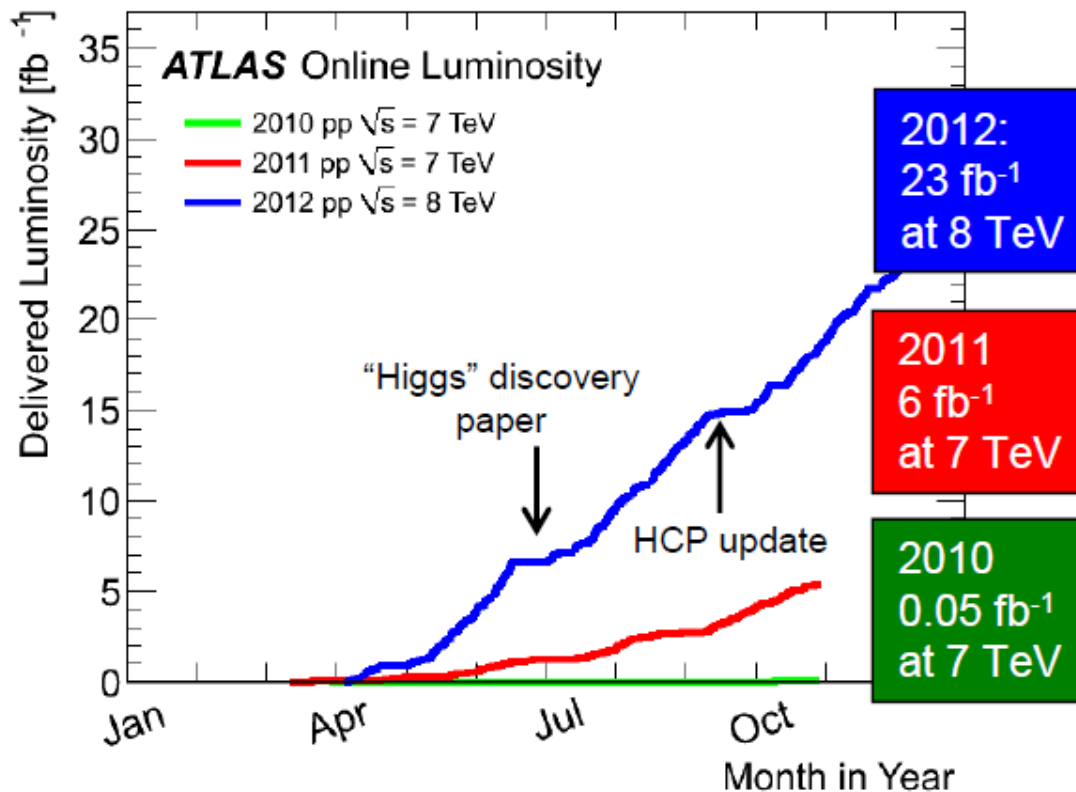
Spontaneous Symmetry  
Breaking

Higgs particle

# Brief summary of recent results reported at Moriond

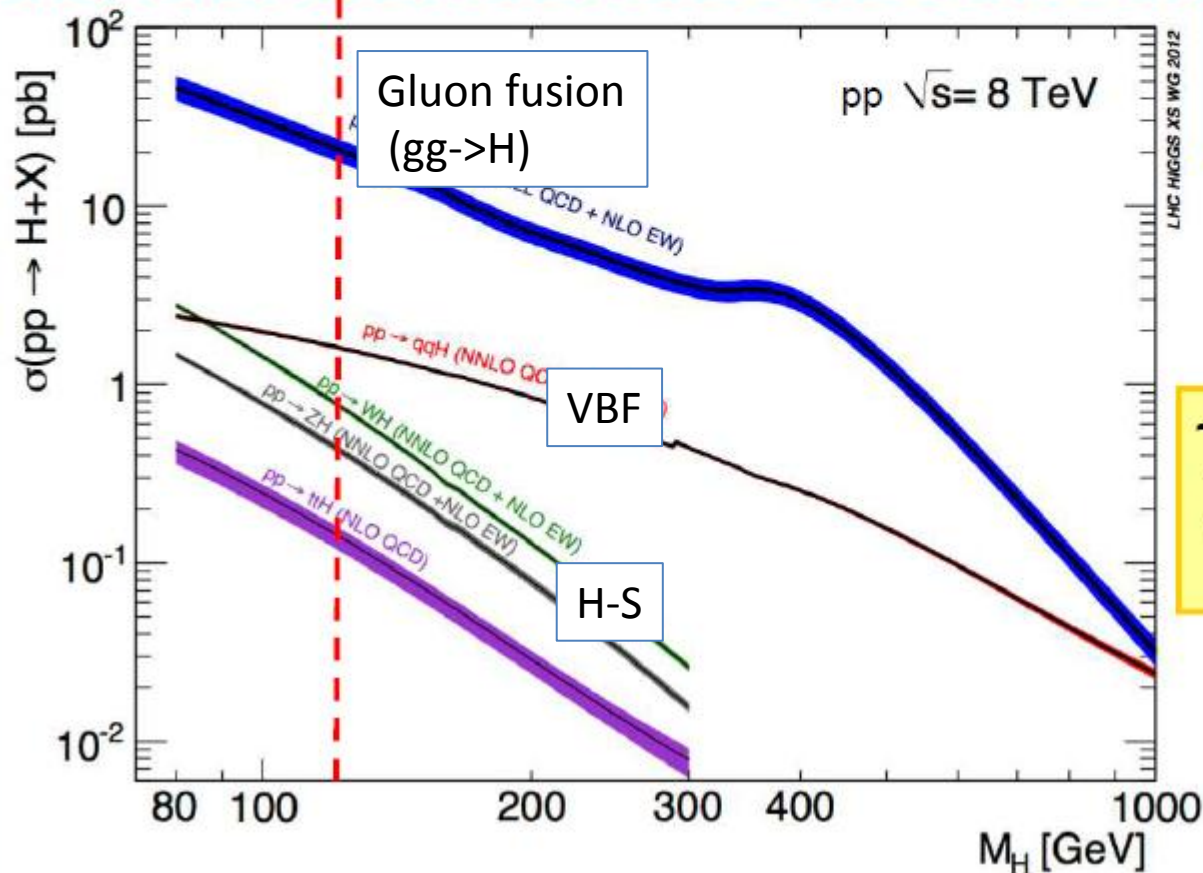
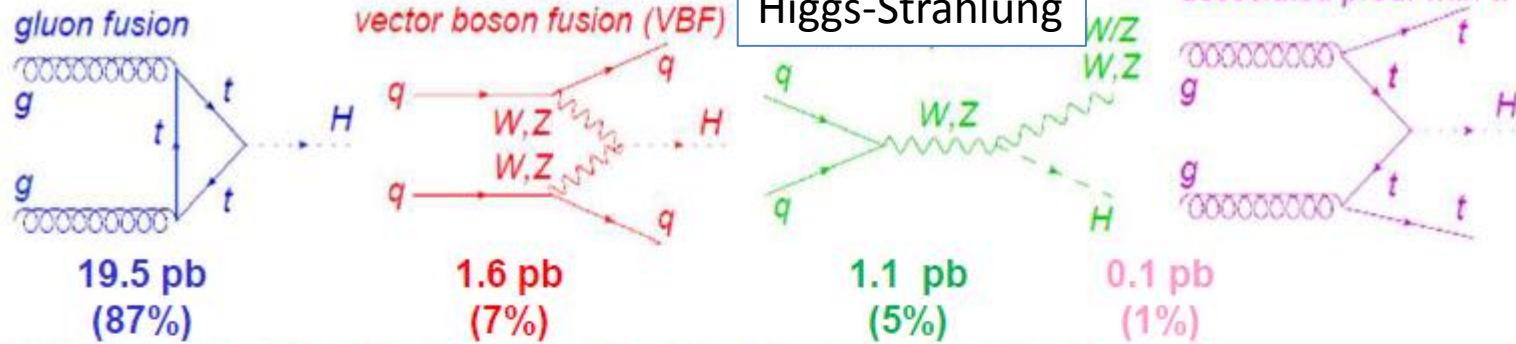
Discovery of Higgs(-like particle) was reported by CMS and ATLAS groups:

7 TeV  $L = 4.8 / \text{fb}$  (2011) + 8 TeV  $L = 20.7 / \text{fb}$



# Production of SM Higgs

## Higgs-Strahlung



**$\sigma(\text{pb})$  @125 GeV  
for  $pp \sqrt{s} = 8 \text{ TeV}$   
Total 22.3 pb**

**~2 Higgs bosons @  
125 GeV produced  
at LHC out of  $10^{10}$   
pp collisions**

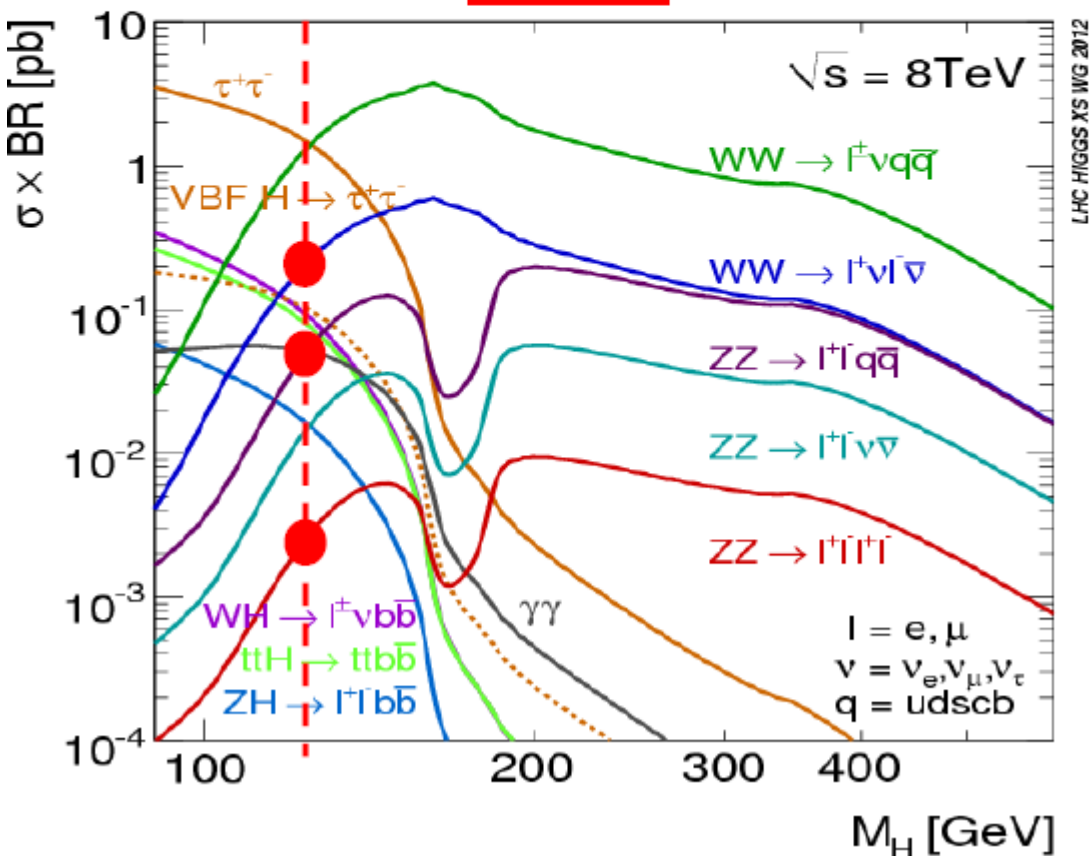
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# Decay modes of SM Higgs

## Branching ratios at 125 GeV:

$bb$ :	57.7%	$ZZ$ :	2.6%
$WW$ :	21.5%	$\gamma\gamma$ :	0.23%
$\tau\tau$ :	6.3%		

## $\sigma \times \text{BR}$ for « useable » final states



3 major channels for low mass Higgs

- $\gamma\gamma$  final state
  - $ZZ^* \rightarrow 4\ell$  final state
  - $WW^* \rightarrow 2\ell + 2\nu$  final state
- 1% mass resolution

$b\bar{b}, \tau\tau$

Huge backgrounds:

- high- $p_T$   $b$ -jet production:
  - $\rightarrow \sim 10^6$  larger than  $H \rightarrow b\bar{b}$  production
- Drell-Yan/ $Z \rightarrow \tau^+\tau^-$ :
  - $\rightarrow 10^5$  larger than  $H \rightarrow \tau^+\tau^-$  production

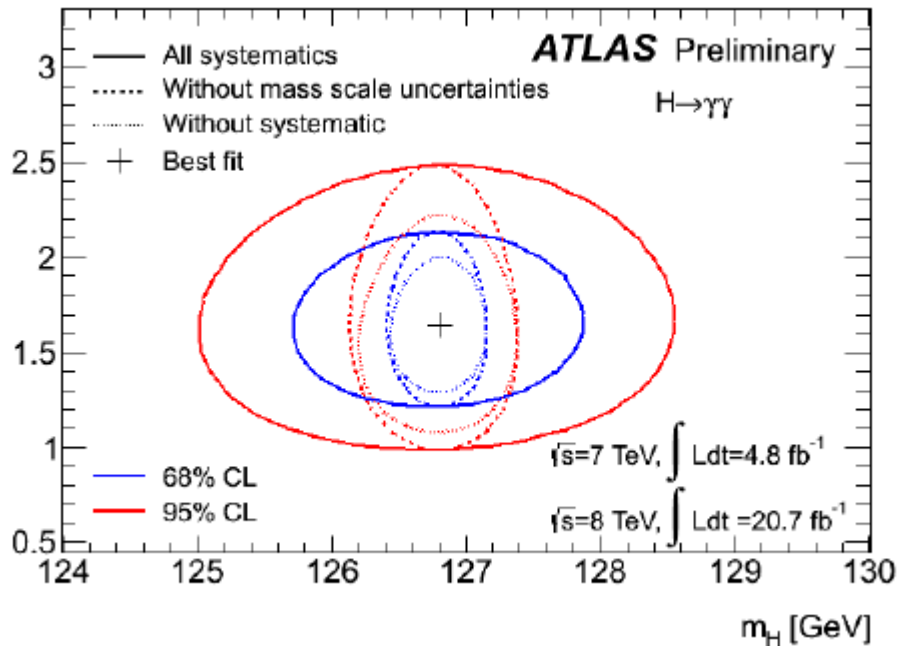
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$H \rightarrow \gamma\gamma$

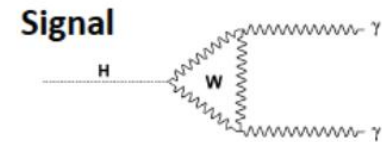
CS is low (50 fb), but relatively clean bkg

→ around 1250 events are expected @ 25 /fb  
identification efficiency 40 % (ATLAS)

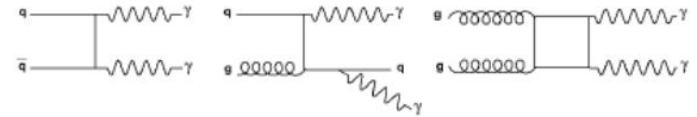
**$126.8 \pm 0.2$  (stat)  $\pm 0.7$  (syst) GeV**



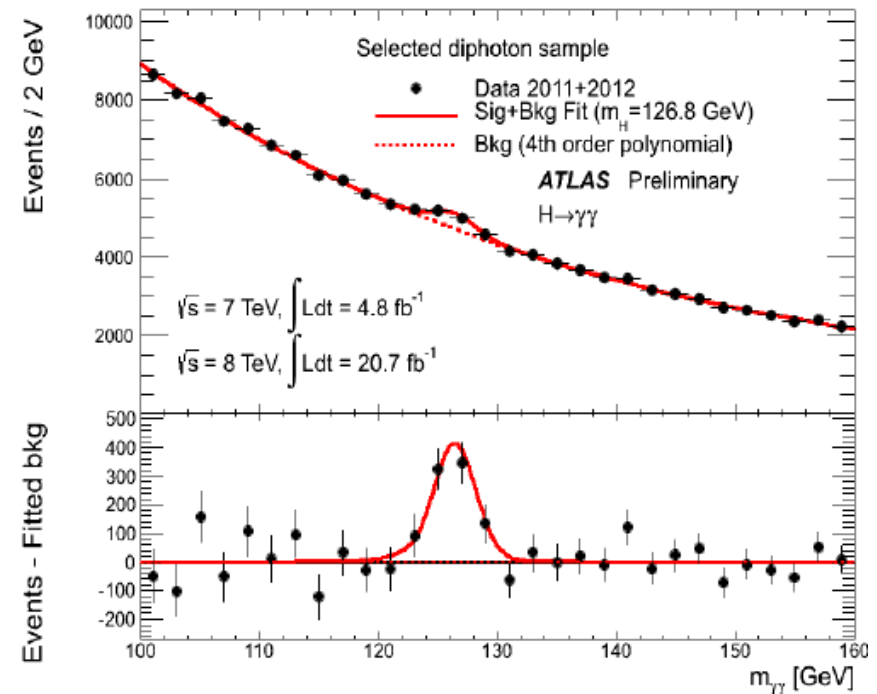
**$\mu = 1.65^{+0.34}_{-0.30} = 1.65 \pm 0.24$  (stat)  $^{+0.25}_{-0.18}$  (syst)**



**Irreducible background:  $pp \rightarrow \gamma\gamma + X$**

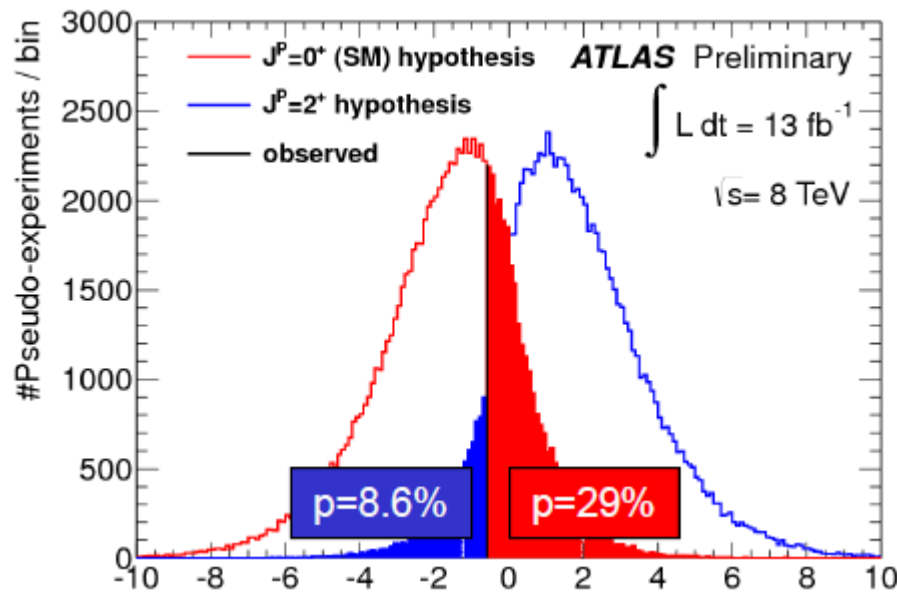


ATLAS



The data went closer to SM. Consistency with SM at  $2.3 \sigma$   
CMS did not report on  $H \rightarrow 2\gamma$

# Spin-Parity



From angular distribution of photons, we can identify whether it is spin 0 or 2

- Spin-2<sup>+</sup> hypothesis expected exclusion  $CL_s$  at 93% [for 100% gg spin-2 production]
- Observation compatible with spin-0<sup>+</sup>, slightly favored over spin-2<sup>+</sup> hypothesis

$$H \rightarrow ZZ^{(*)} \rightarrow 4l \quad (4e, 4\mu, 2e2\mu)$$

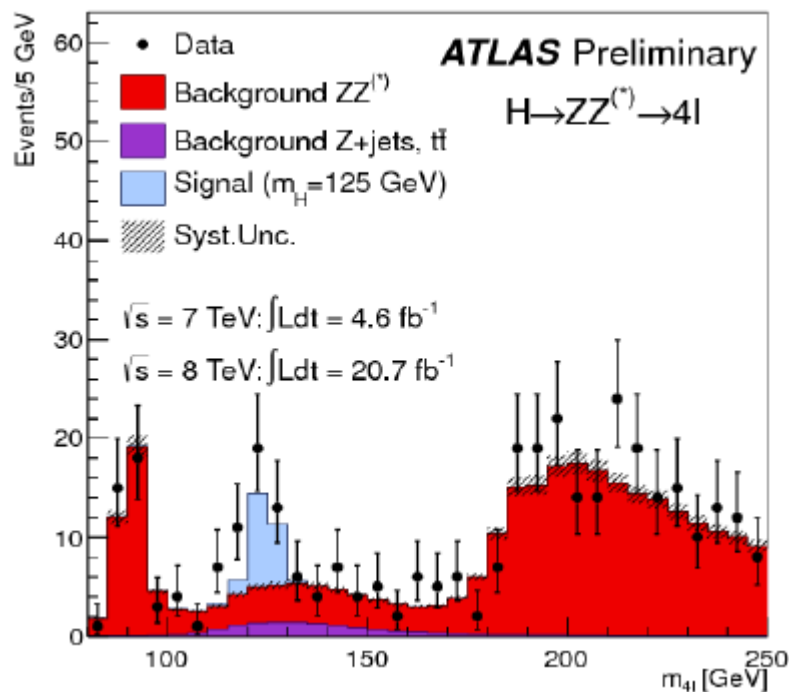
$$\sigma \times \text{BR} \sim 2.5 \text{ fb} \quad m_H \sim 126 \text{ GeV}$$

Gold plated mode

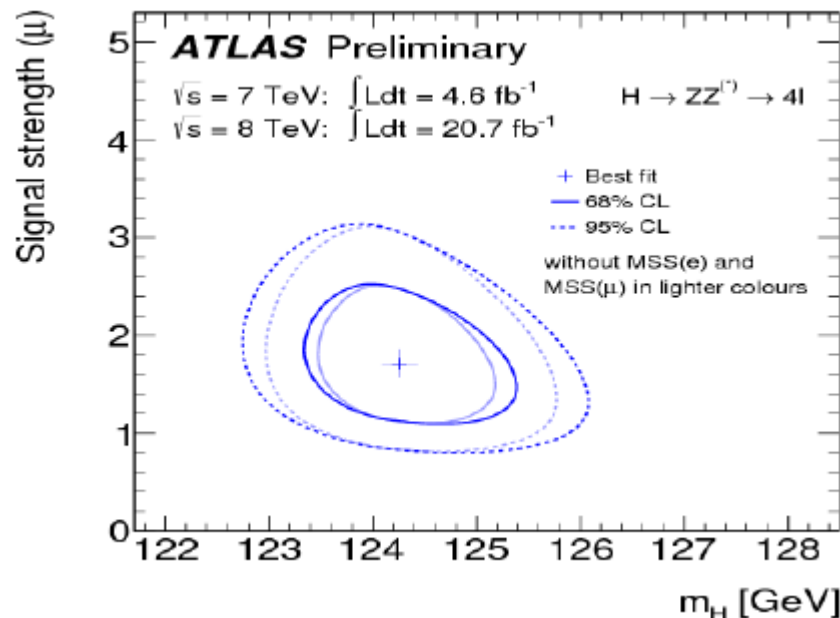
Clean signature: narrow peak, low background

□ Tiny rate, BUT:

- mass can be fully reconstructed
- pure: S/B  $\sim 1$



In region  $125 \pm 5 \text{ GeV}$ : 32 events  
observed  $[11.1 \pm 1.3]$  expected from  
bknd &  $15.9 \pm 2.1$  from SM Higgs]



Best mass fit  $124.3^{+0.6}_{-0.5} \text{ (stat)}^{+0.5}_{-0.3} \text{ (syst) GeV}$

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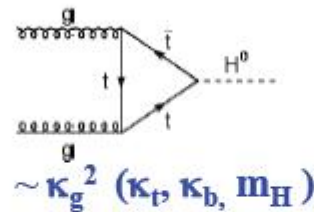


# Higgs coupling : information from production and decay rates

- For each observed final state, production and decay involve several couplings

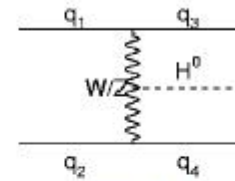
- Best example:  $\gamma\gamma$

– Production

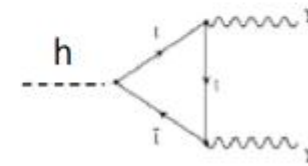
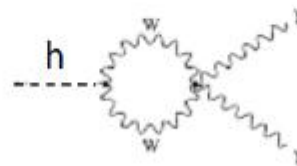
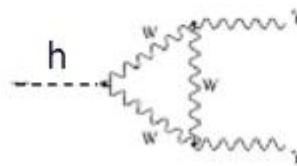


+

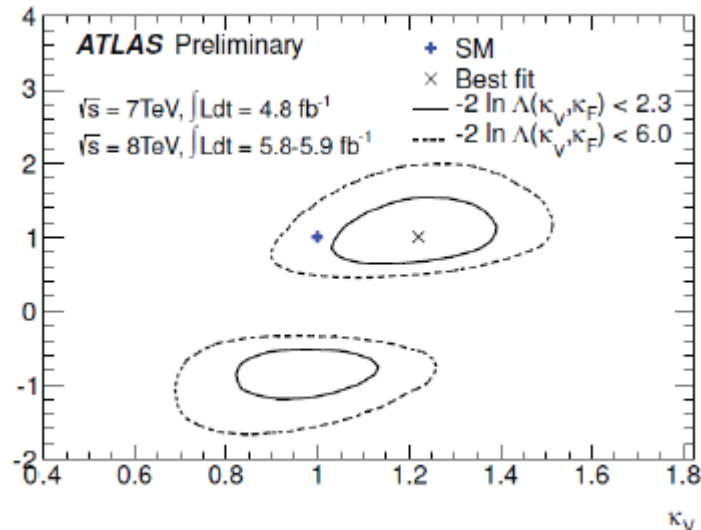
$\sim \kappa_W^2$



– Decay



Decay width :  $\sim (\kappa_W - 0.2 \kappa_t)^2$  [ note: interference ]



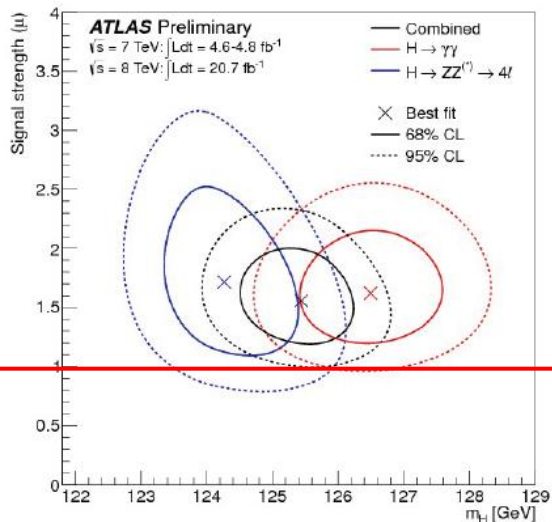
- Higgs couplings (partial dataset)  
Compatible with SM hypothesis.
- $\text{BR}(H \rightarrow \text{inv.}) < 65\% \text{ (95\% C.L.)}$

Consistent with SM  
No sign of BSM

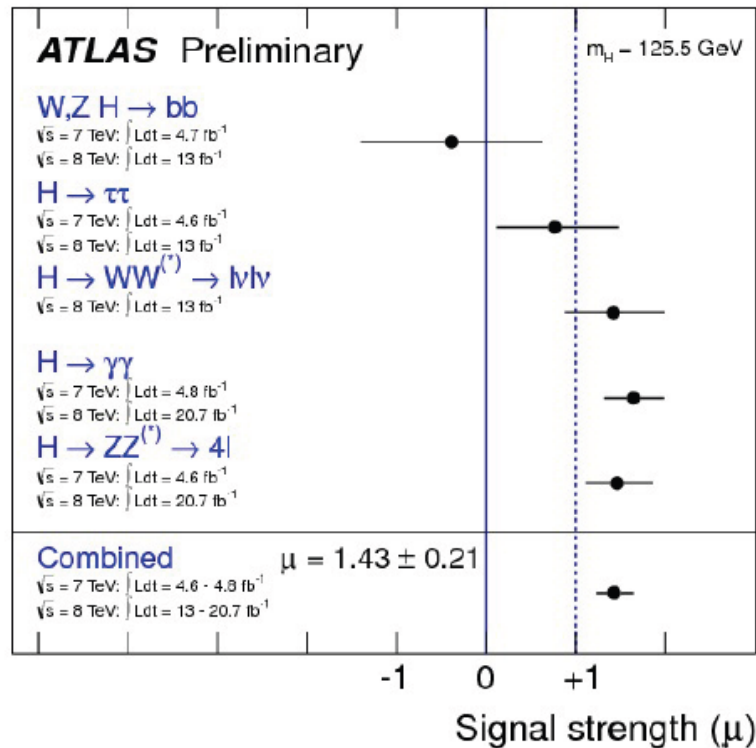
[back](#) to main



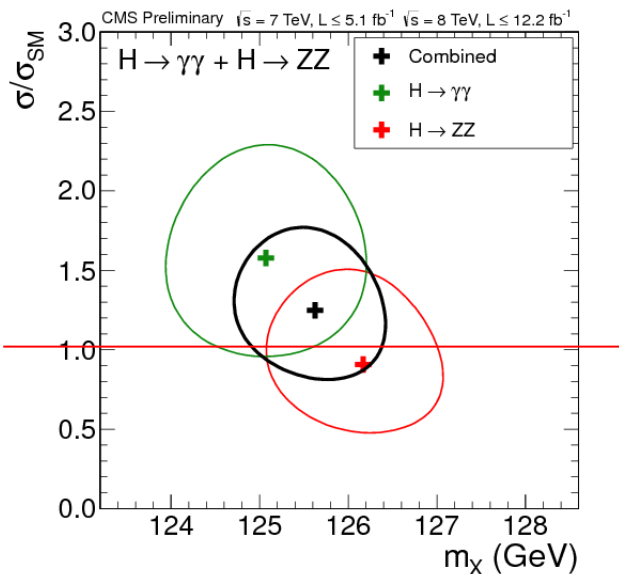
# ATLAS



SM



# CMS (Dec 12)



$$M_h = 125.8 \pm 0.4 \pm 0.4 \text{ GeV (CMS)}$$

$$125.5 \pm 0.2^{+0.5}_{-0.6} \text{ GeV (ATLAS)}$$

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