Non-perturbative study of the Higgs-Yukawa model using lattice field theory

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Outline

- Motivation.
- Do Higgs and Yukawa live close to a critical point?
 --- ideas and strategy (non-chiral example).
- Preliminary results (chiral theory) from our ongoing study.
- Outlook.

Motivation



P. Gerhold and K. Jansen, 2011

* Constraints on the masses of extra-generation fermions from the 125 GeV scalar.

The I25 GeV scalar

• It may be a dilaton in a strongly-coupled theory:

Does it have to be walking technicolour?
 HY model exhibits quasi scale inavariance?

P.Q. Hung and C. Xiong, 2009

• It may be the Standard Model Higgs:

Evade the hierarchy problem w/o SUSY?

• Both require non-perturbative studies:

Second-order non-thermal phase transitions.

Hierarchy and triviality problem: perturbation theory (misleading)

- Scalar mass operator is of dimension 2 and is not protected by chiral symmetry.
- The one-loop beta-functions for the scalar and Yukawa coupling are positive.
- Perturbation theory over-simplifies the problem and may lead to misleading statements.

The scalar field theory as a spin model Scalar theory on the lattice (a=1),

$$S_{\varphi} = -\sum_{x,\mu} \varphi_x^{\alpha} \varphi_{x+\hat{\mu}}^{\alpha} + \sum_x \left[\frac{1}{2} (2d+m_0^2) \varphi_x^{\alpha} \varphi_x^{\alpha} + \frac{1}{4} \lambda_0 (\varphi_x^{\alpha} \varphi_x^{\alpha})^2 \right]$$

• Perform the change of variables,

$$\Phi^{\alpha} = \sqrt{2\kappa}\phi^{\alpha}, \quad \lambda_0 = \frac{\hat{\lambda}}{\kappa^2}, \quad \bar{m}_0^2 = \frac{1 - 2\hat{\lambda} - 8\kappa}{\kappa}$$

Bulk phase structure of the resulting spin model,

$$Z_{\phi} = \int \prod_{x,\alpha} d\phi_x^{\alpha} \exp(-S_{\phi}) = \int \prod_{x,\alpha} d\mu(\phi_x^{\alpha}) \exp\left(2\kappa \sum_{x,\mu} \phi_x^{\alpha} \phi_{x+\hat{\mu}}^{\alpha}\right),$$

$$d\mu(\phi_x^{\alpha}) = d\phi_x^{\alpha} \exp\left[-\phi_x^{\alpha} \phi_x^{\alpha} - \hat{\lambda}(\phi_x^{\alpha} \phi_x^{\alpha} - 1)^2\right].$$

Fermions

- The overlap fermion (exact chiral symmetry).
- The lattice Yukawa operator takes the same form as its continuum counterpart.
- Extremely computationally demanding.

What is it like with the Yukawa coupling



* Question: Is the theory non-trivial in 4D?

At stronger bare Yukawa coupling



The bulk phase structure (3D)



Only the Gfp remains in 4D scalar sector.... *The hierarchy problem is a consequence of triviality in 4D

The 4D bulk phase structure



Evidence for a tri-critical point? If so, is the Yukawa coupling non-trivial there?

Our target

- Study the chiral theory.
- Investigate the phase structure in detail.
- Make contact with phenomenology.

Finite-size scaling (a'la M. Fisher)

- Renormalisation Group near fixed points.
- Central statement: "Universal" function

$$\frac{P_L(t)}{P_{\infty}(t)} = f\left(\frac{L}{\xi_{\infty}(t)}\right) \text{, with observable P.}$$

Magnetic susceptibility and Binder's cumulant:

$$\chi_m(t,L) \cdot L_s^{-\gamma/\nu} = g\left(\hat{t}L_s^{1/\nu}\right), \text{ with } \hat{t} = \left[T/\left(T_c^{(L=\infty)} - C \cdot L_s^{-b}\right) - 1\right]$$
$$Q_L = g_{Q_L}\left(tL^{1/\nu}\right)$$

γ and ν are the critical exponents.
 How different are they from the mean-field values?

4D scaling test, susceptibility



	$T_c^{(L=\infty)}$	u	γ	C	b	fit interval
$\kappa = 0.06$	18.119(67)	0.576(28)	1.038(30)	4.7(1.6)	1.95(18)	17.5, 20.0
$\kappa = 0.00$	16.676(15)	0.541(22)	0.996(15)	10(2)	2.42(10)	15.0, 19.0
O(4)	0.304268(27)	0.499(12)	1.086(19)	N/A	N/A	0.300, 0.308



Concluding remarks and outlook

- Evidence for novel FP in the HY model.
- Complication in 4d (work in progress)

Gaussian FP in the scalar sector.
 Does it remain in the HY model?
 Logarithmic corrections to FSS.

• Spectrum calculation (on-going work).