World of ultracold atoms

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Early world \rightarrow Cosmology

Our ordinary world

Single (or few)

atoms:

Small world→ Particle phys.

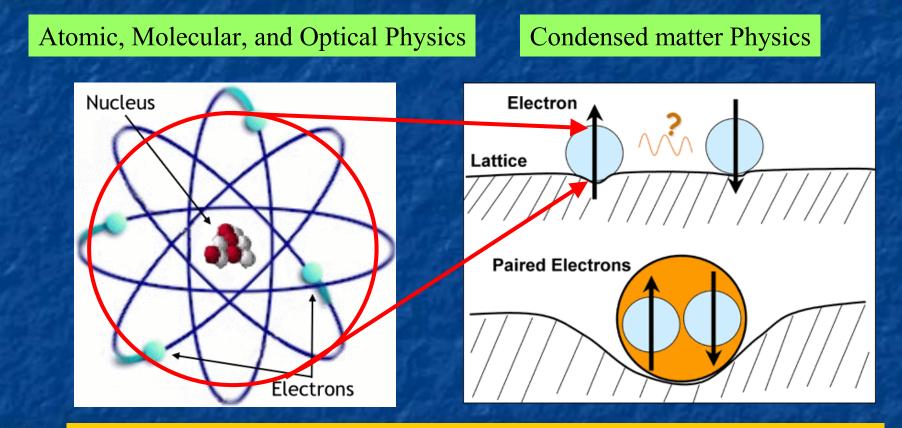
 $\begin{array}{l} \text{High } T \rightarrow \text{Plasma} \\ \text{AMO Phys.} \\ \text{Cold atoms} \end{array}$

Many-particle (or atoms):

Large world → Astrophysics

Future world → Applied Physics

Ultracold atoms as an emergent field....



Systems of ultracold atoms can be understood as a many-body system of atoms, which are strongly affected by the fruitful internal degrees of freedom of each single atom.

An Interdisciplinary field

Precise measurement

Traditional AMO

It Cosmology 8 High energ

Ultracold atoms Quantum Information

Nonlinear Physics

Soft-matter/ chemistry

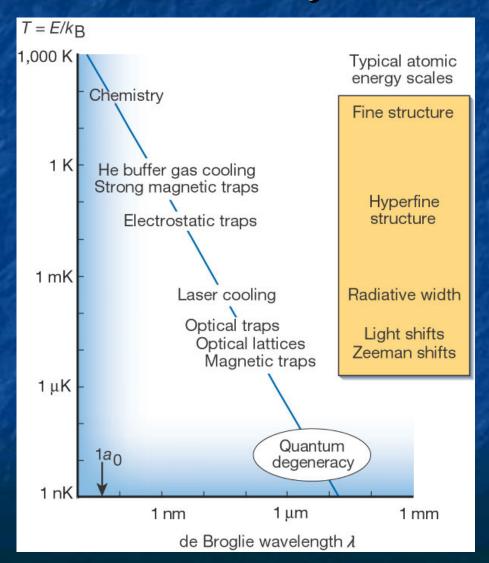
Condensed matter

How cold is an "ultracold system"?

Thermal wavelength

$$k_{B}T \sim \frac{\hbar^{2}}{2m} \left(\frac{2\pi}{\lambda}\right)^{2}$$

Quantum degeneracy: $\lambda \sim L \sim 1 - 100 \mu m$ $T < 1 \mu K$



Why low temperature ?

Ans: To see the quantum effects !

Uncertainty principle:

$$\Delta x \Delta p \ge \hbar$$

$$\frac{\Delta p^2}{2m} \sim k_B T \rightarrow \Delta x \sim \frac{\hbar}{\Delta p} \sim \frac{\hbar}{\sqrt{2mk_B T}} \equiv \lambda_T, \text{ Thermal wavelength}$$

Therefore, if
$$T \downarrow \Rightarrow \lambda_T \uparrow$$

Quantum regime when $\lambda_T \geq d \sim n^{-1/3}$

$$\lambda_T$$

Why strong interaction ?

P. Anderson: "Many is not more"

Because interaction can make "many" to be "different" !

Example: 1D interacting electrons → crystalization and no fermionic excitation

How to make interaction stronger?

$$H = \sum_{j=1}^{N} \left(\frac{\mathbf{p}_{j}^{2}}{2m} + V(x_{j}) \right) + \frac{1}{2} \sum_{i \neq j}^{N} U(x_{i} - x_{j})$$

1.
$$U(x)$$
 becomes stronger

2. $E_k \sim k_B T$ becomes smaller or *m* becomes smaller

3. V(x) changes to make lower dimension

4. *N* becomes larger (for short interaction); smaller for long range interaction

How to reach ultracold temperature ? See also: Prof. Yu's talk

1. Laser cooling ! (few $K \rightarrow mK$)

Use red detune laser + Doppler effect

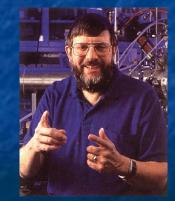
1997 Nobel Price



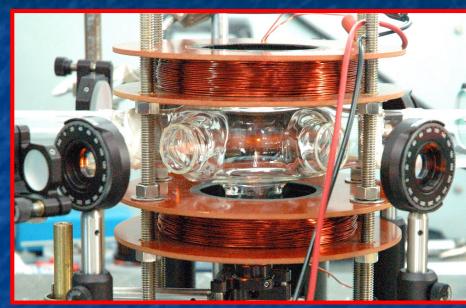
Steven Chu



Claude Cohen-Tannoudji



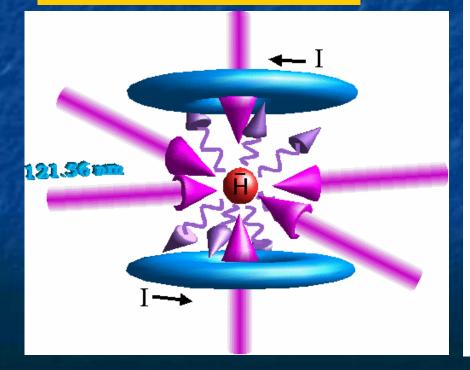
Williams D. Phillips



How to reach ultracold temperature ?

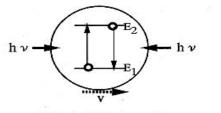
Laser cooling ! (1997 Nobel Price)

Use red detune laser + Doppler effect



II Principle of Laser Cooling & Trapping

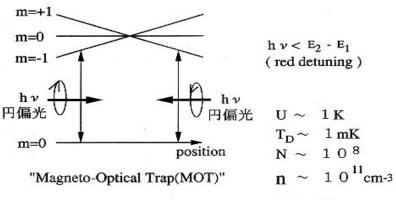
1) Doppler Cooling



 $h \nu < E_2 - E_1$ (red detuning)

 $\begin{array}{cc} T_{\rm D} \sim & 1 \text{ mK} \\ (v \sim 1 \text{ m/s}) \end{array}$

"Optical Molasses"

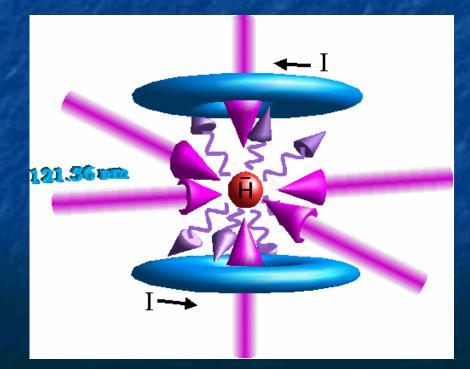


 $\tau \sim 1 \, \mathrm{hour}$

How to reach ultracold temperature ?

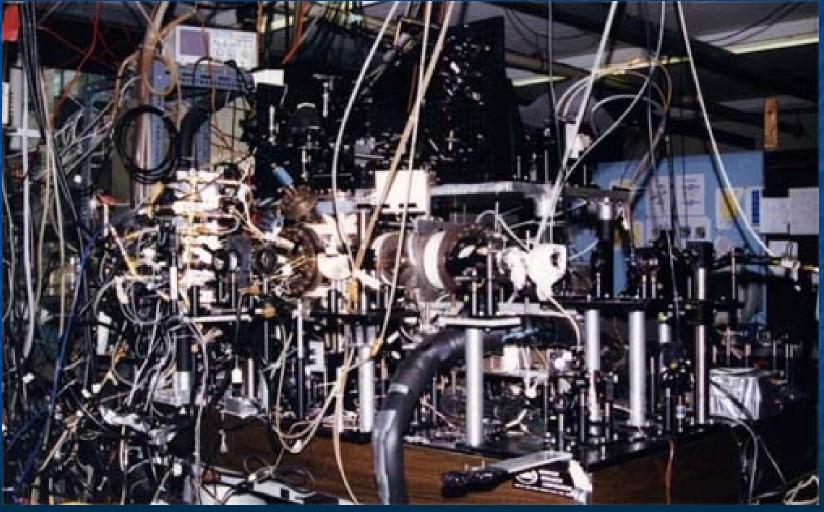
2. Evaporative cooling !

Reduce potential barrial +thermal equilibrium





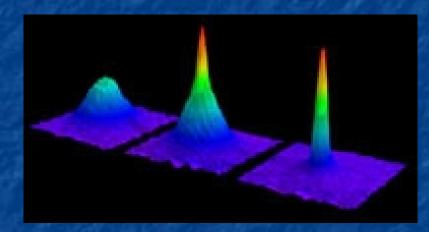
Typical experimental environment





How to do measurement ?

Trapping and cooling Perturbing Releasing and measuring







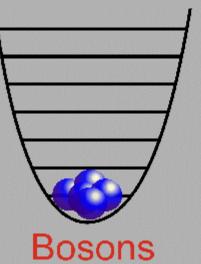
(2001 Nobel Price)

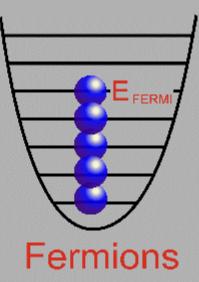
What is Bose-Einstein condensation ?

$$\Psi(x_1, x_2) = \pm \Psi(x_2, x_1)$$
, + for boson and - for fermion

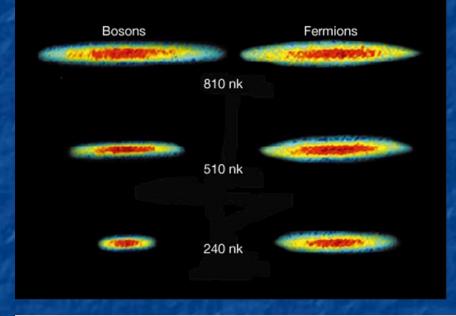
Therefore, for fermion we have $\Psi(x, x) = 0$, i.e. fermions like to be far away, but bosons do like to be close !

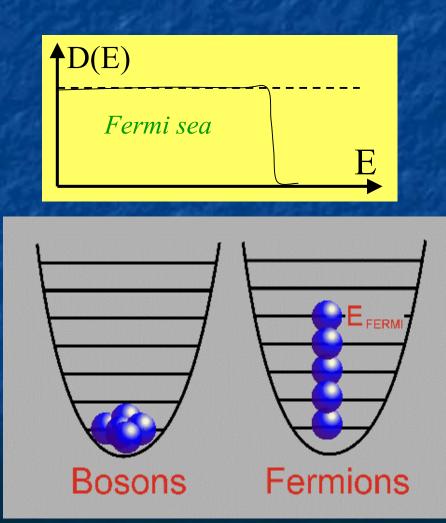
When *T* is small enough, noninteracting bosons like to stay in the lowest energy state, i.e. BEC





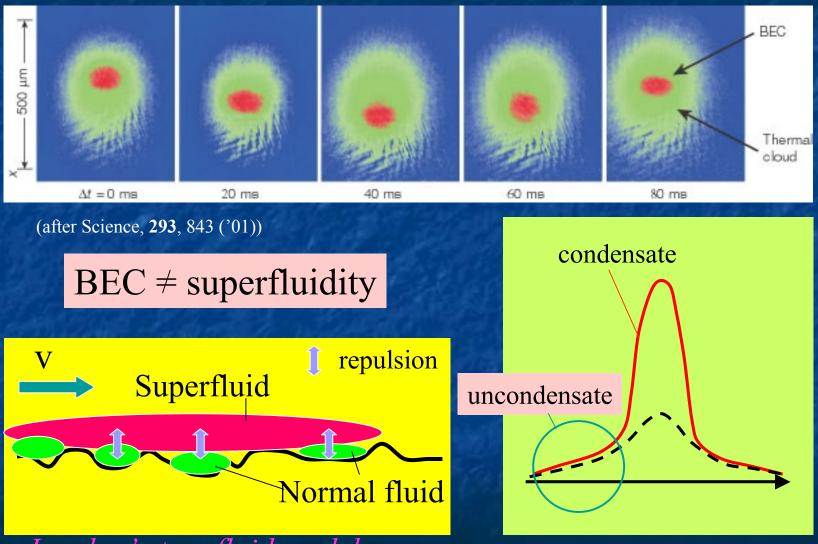
How about fermions in T=0 ?





When T-> 0, noninteracting fermions form a compact distribution in energy level.

BEC and Superfluidity of bosons



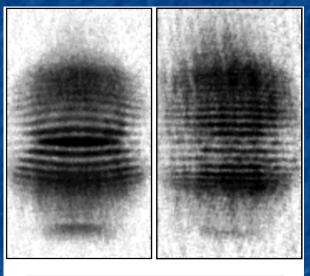
Landau's two-fluid model

Phonons and interference in BEC

Phonon=density fluctuation

0.5 mm $n_0 U$ V_{ph} m

Interference



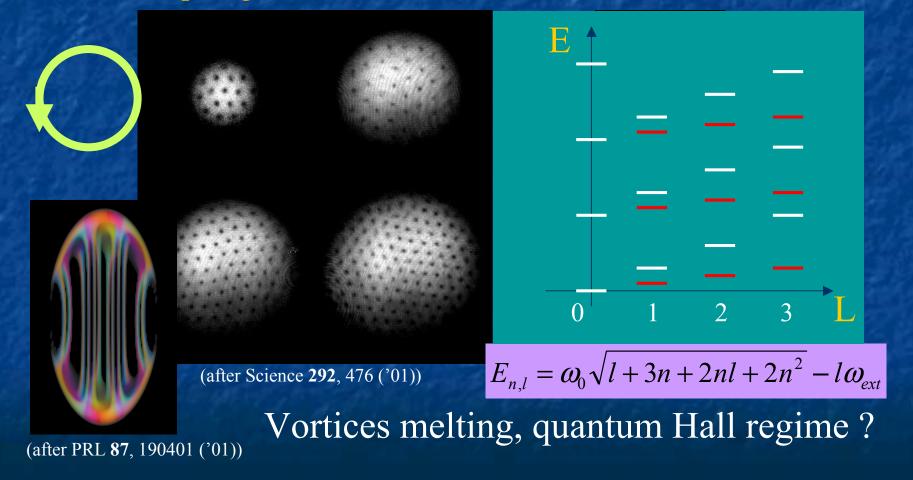


n

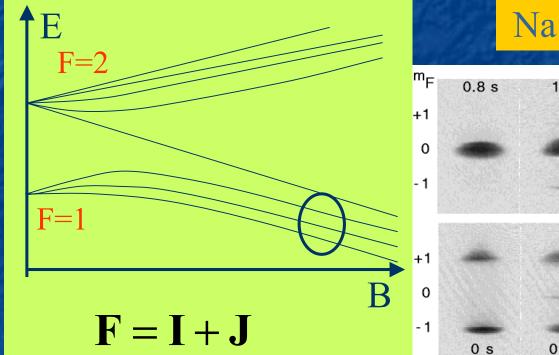
(after Science 275, 637 ('97))

Matter waves ?

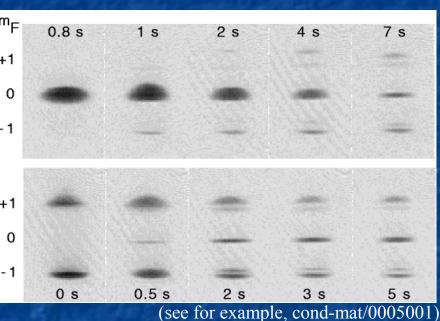
Vortices in condensate Vortex = topological disorder



Spinor condensation in optical trap

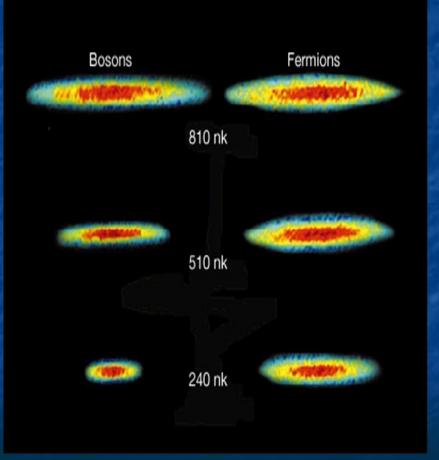


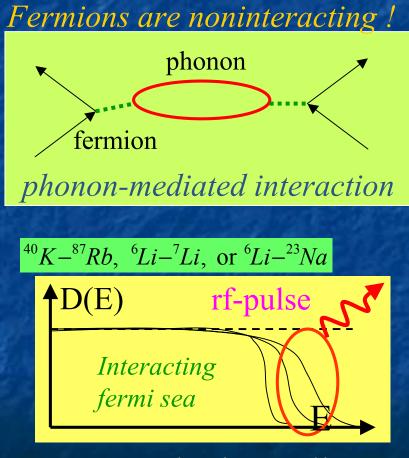
Na
$$|F=1, m_F=\pm 1, 0\rangle$$



$$H = \int d\mathbf{r} \left[-\hat{\psi}_i^+ \frac{\nabla^2}{2m} \hat{\psi}_i + \frac{g_0}{2} \hat{\psi}_i^+ \hat{\psi}_j^+ \hat{\psi}_j \hat{\psi}_i + \frac{g_2}{2} \left(\hat{\psi}_i^+ \mathbf{F}_{ij} \hat{\psi}_j \right) \cdot \left(\hat{\psi}_k^+ \mathbf{F}_{kl} \hat{\psi}_l \right) \right]$$

Boson-fermion mixtures

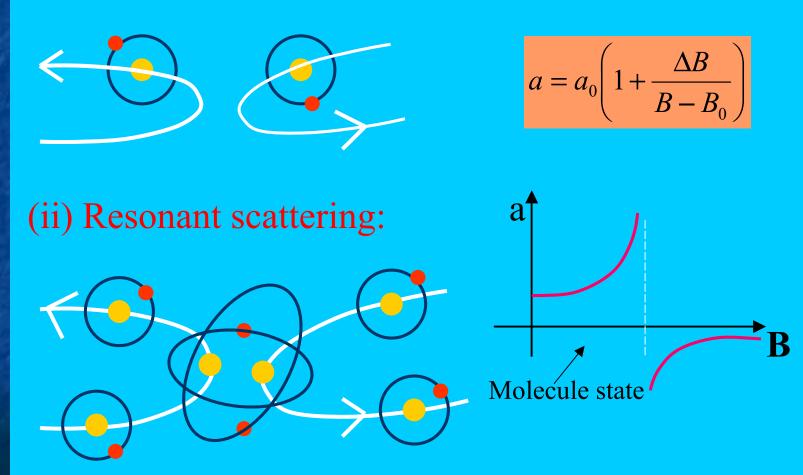




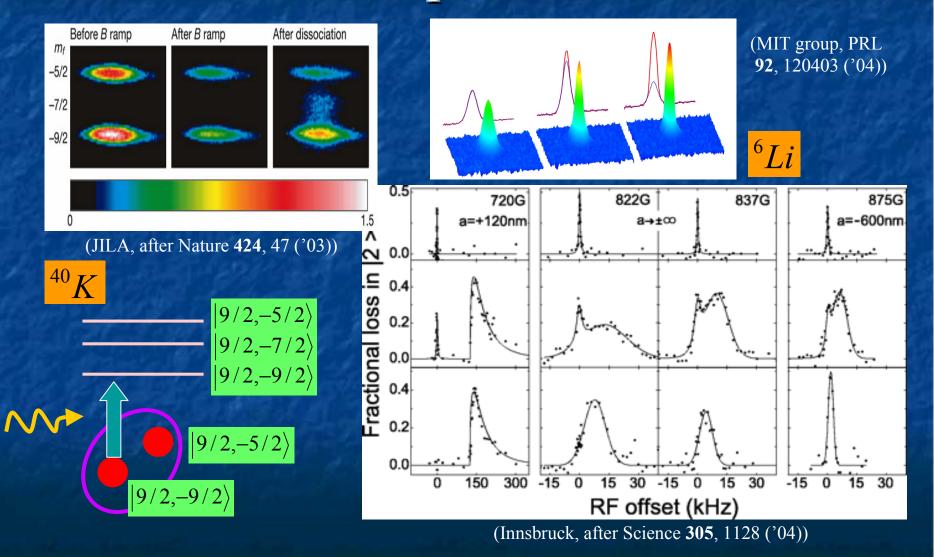
Sympathetic cooling

Feshbach Resonance

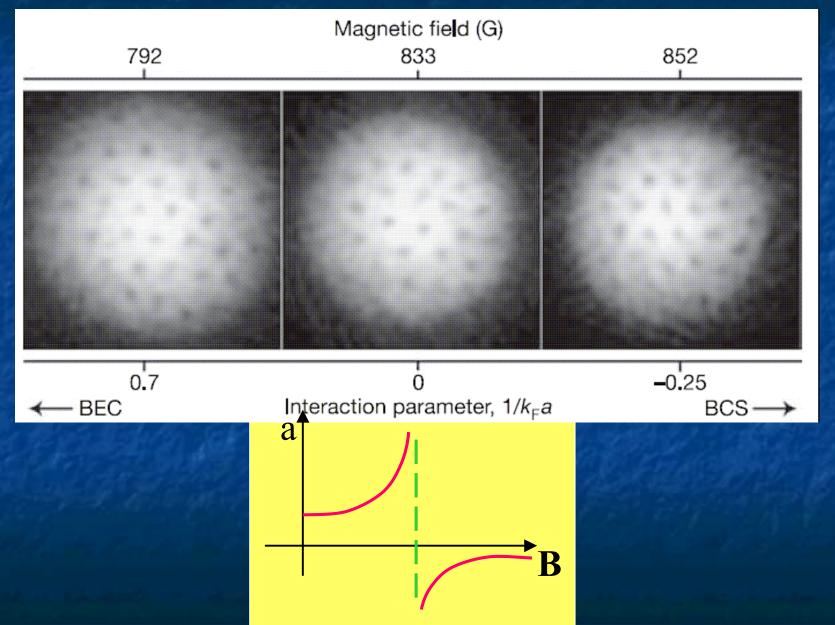
(i) Typical scattering:



Molecule and pair condensate

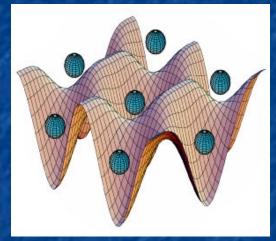


First evidence of superfluidity of fermion pairing

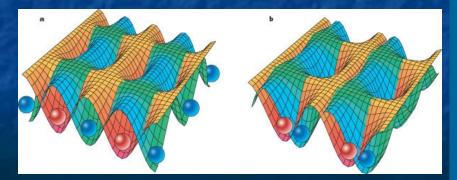


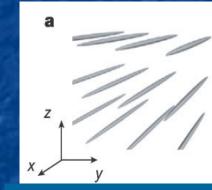
Optical lattice

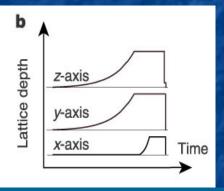
3D lattice

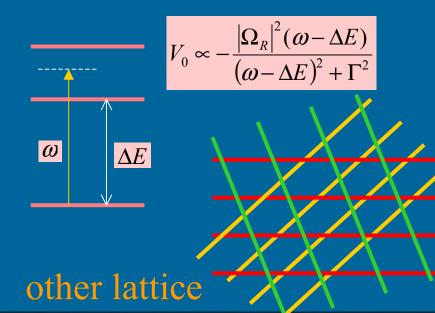


Entanglement control









Mott-Insulator transition

Bose-Hubbard model

$$H = -t \sum_{\langle i,j \rangle} a_i^+ a_j + U \sum_i a_i^+ a_i (a_i^+ a_i - 1) - \mu \sum_i a_i^+ a_i$$

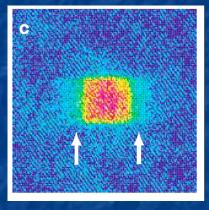
$$\begin{array}{c} \mu \\ n=3 \\ superfluid \\ n=2 \\ n=1 \\ t/U \end{array}$$

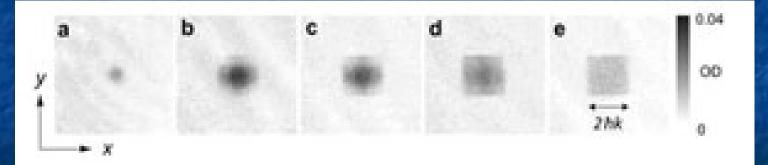
(after Nature **415**, 39 ('02))

Fermions in optical lattice

Fermi Hubbard model

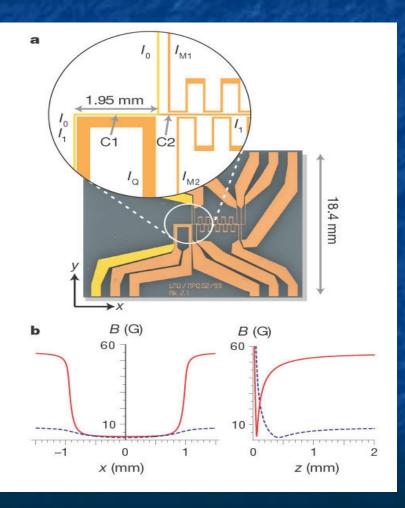
$$H = -t \sum_{\langle i,j \rangle} a_{i,s}^{\dagger} a_{j,s} + U \sum_{i} n_{i,\uparrow} n_{i,\downarrow}$$

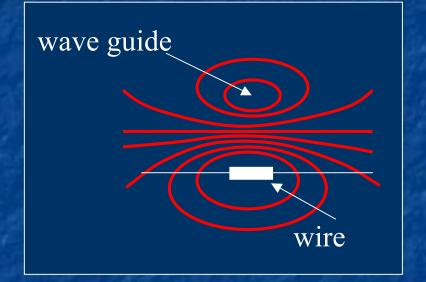


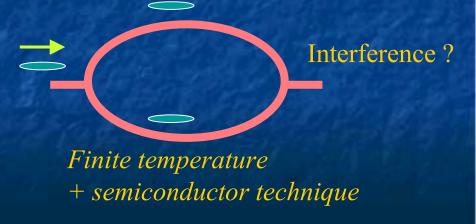


Superfluidity of fermion pairing in lattice is also realized.

Transport in 1D waveguide





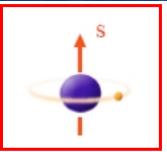


Cold dipolar atoms/molecules

(1) Heteronuclear molecules



(2) Atoms with large magnetic moment



(a) Direct molecules
p~ 1-5 D
(b) But difficult to be cooled

(Doyle, Meijer, DeMille etc.)

Small moment

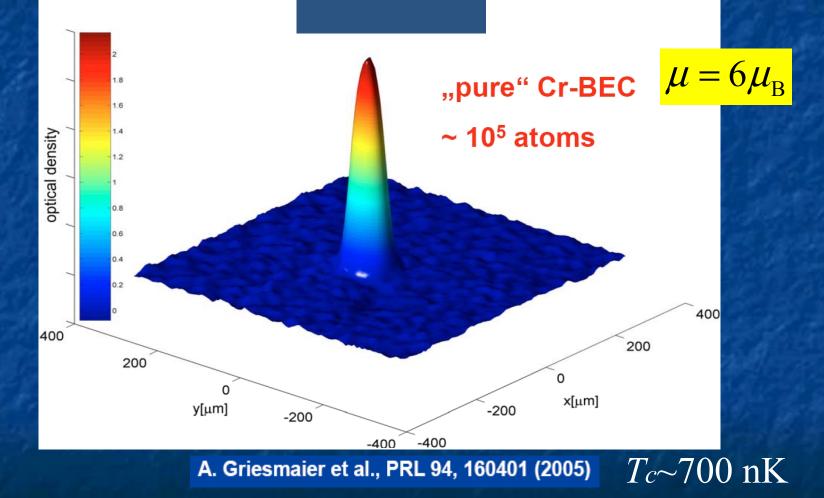
 $\mu \sim 6\mu_{R}$ (for Cr)

But it is now ready to go !

(Stuhler etc.)

 $p \sim 1D, U_{dd} \sim 10 \mu K, \ \mu = 1 \mu_B, U_{dd} \sim 1n K$

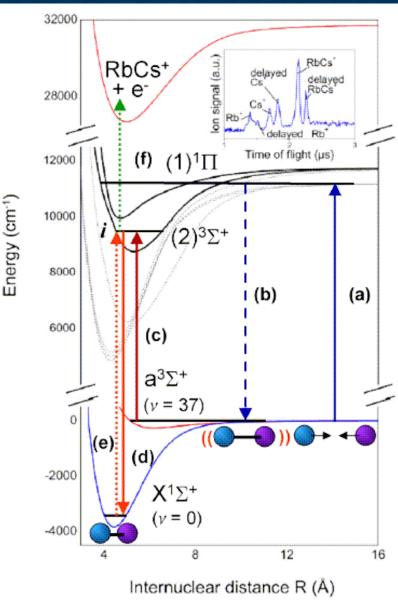
Condensate (superfluid)



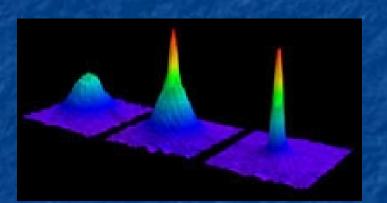
Artificial dipoles:

(KRb, JILA, ETH, etc.) But not in ground state weak dipole moment short life time (RbCs, Yale, etc.) Now in ground state But number of atoms are still small

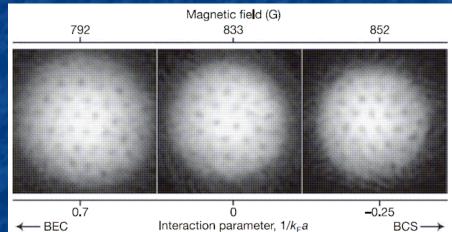
(J. Sage, et. al., PRL, 94, 2030



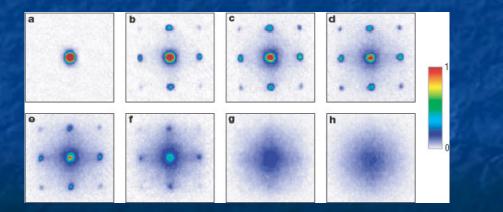
Gallery of pictures



Bose-Einstein Condensation



BEC-BCS crossover



Superfluid-Mott insulator transition

Fermi surface in optical lattice