

FIMP Miracle of Sterile Neutrino by Scale Invariance



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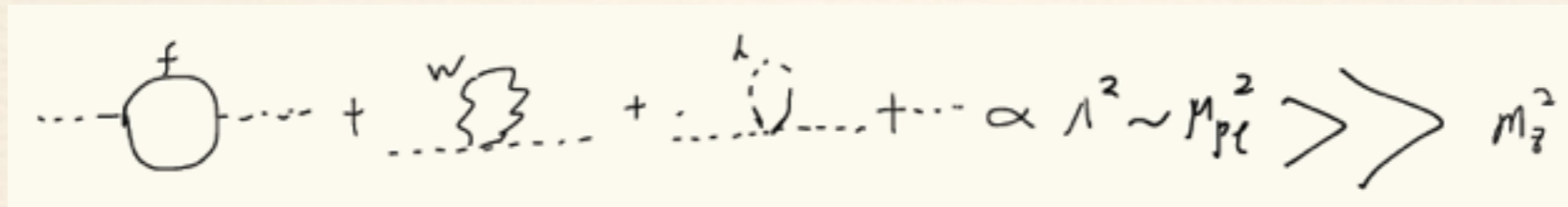
Outline

- ❖ Three leading new physics
- ❖ Scale invariant ν MSM
- ❖ Conclusions

Three leading new physics

- ❖ Theoretically: gauge hierarchy problem

Why does electro-weak (EW) scale exist facing M_{pl} ?



The image shows a series of Feynman diagrams representing fermion self-energy corrections. From left to right: a fermion line with a self-energy loop (circle with 'f' inside); a fermion line with a gluon loop (wavy line with 'g' inside); a fermion line with a photon loop (dashed line with 'h' inside); and a fermion line with a Higgs loop (dashed line with 'h' inside). To the right of these diagrams is the equation $\alpha \Lambda^2 \sim M_{\text{pl}}^2 \gg m_f^2$.

Solutions: supersymmetry, composite, ..., **scale invariance?**

- ❖ Phenomenologically: nonzero neutrino masses & the existence of dark matter (DM)

Adding right-handed neutrinos (RHNs) & neutral stable particle

Three leading new physics

❖ Address them in an unified framework?

SUSY: solving gauge hierarchy problem & offering LSP as DM candidate.
However, it is not a theory of neutrino

ν MSM: DM is the keV scale RHN, **unifying RHN and DM!!!** But it
can not address the hierarchy problem

❖ Their dark matter is beautiful

Dark matter is **predicted** instead of added!

Stability of DM is not required but **accident!**

Correct relic density is via a **miracle**, for LSP being WMIP...

Scale invariant ν MSM (ν SISM)

❖ ν MSM version 1.0

T. Asaka, S. Blanchet and M. Shaposhnikov, Phys. Lett. B 631 (2005) 151.

ν MSM=SM+RHNs=the canonical seesaw with very low seesaw scale

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + \bar{N}_I i \partial_\mu \gamma^\mu N_I - F_{\alpha I} \bar{L}_\alpha N_I \Phi - \frac{M_I}{2} \bar{N}_I^c N_I + \text{h.c.},$$

Accidental stability of the lightest RHN, N_1 :

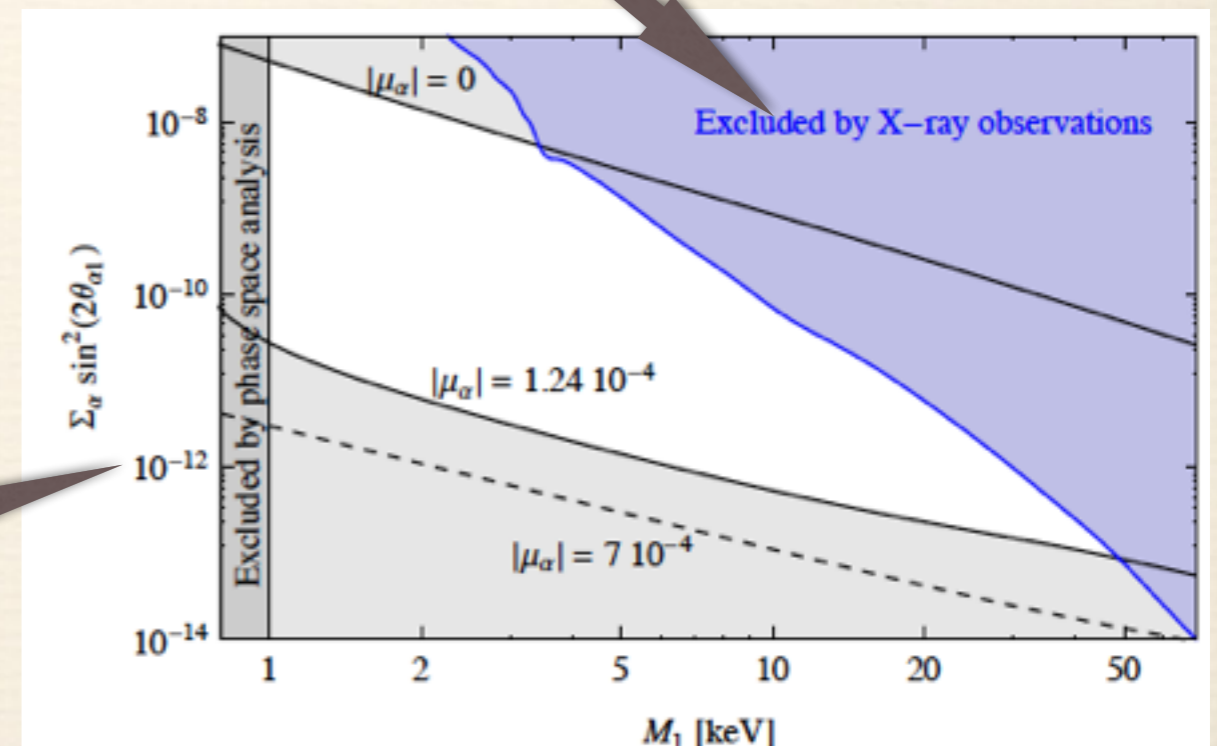
$$\Gamma_{N_1 \rightarrow \nu \gamma} \simeq \frac{9G_F^2 \alpha M_1^5}{256\pi^4} \times \sin^2 \theta_1,$$

Prediction: X-ray line——3.5keV line?

Relic density: sterile-active neutrino oscillation with(out) resonant effect

Lyman-alpha bound

production is problematic



Scale invariant ν MSM (ν SISM)

❖ ν MSM version 2.0

ν MSM 2.0 = ν MSM 1.0 with classical SI (no massive couplings)

$$\mathcal{L} = \frac{\lambda_1}{2}|H|^4 + \frac{\lambda_2}{2}|S|^4 + \lambda_3|H|^2|S|^2 + \lambda_4|H|^2(S^2 + S^{*2}) + \lambda_5|S|^2(S^2 + S^{*2}) + \frac{\lambda_6}{2}(S^4 + S^{*4}) + \left(\frac{\lambda_{sn}}{2}SN^2 + y_N\bar{\ell}HN + c.c. \right).$$

SI demands singlets with VEVs to give Majorana masses of RHNs; One complex singlet is required to accommodate Higgs phenomenology

similar singlet was introduced before, but only here it is a result of symmetry, namely in the spirit of naturalness

No need of any global symmetry, but here imposing CP-invariance in the Higgs sector in order to reduce parameters

Scale invariant ν MSSM (ν SISM)

S. Coleman and E. Weinberg, Phys. Rev. D 7, 1888 (1973).

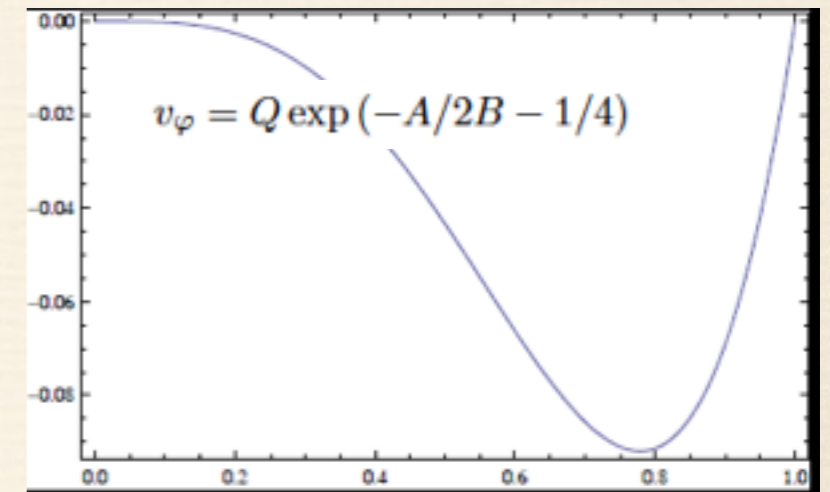
❖ SI spontaneously breaking: origin of VEVs

SI is anomaly, so a scaleless theory can break SI!

$$V_{\text{eff}} = A\phi_{\text{cl}}^4 + B\phi_{\text{cl}}^4 \ln \frac{\phi_{\text{cl}}^2}{Q^2},$$

$$A = \frac{\lambda}{8} + \frac{1}{64\pi^2} \sum_P n_P g_P^4 (-A_P + \ln g_P^2),$$

$$B = \frac{1}{64\pi^2} \sum_P n_P g_P^4,$$



The prediction of a light Higgs state $m_\phi^2 = 8B\langle\phi_{\text{cl}}\rangle^2$

Heavy top makes $B < 0$, so SM can not be SI. But extension by S works!

$\nearrow S = \frac{J+ib}{J^2}$

$$v(J, b, \lambda) = \frac{\lambda_J}{4!} J^4 + \frac{\lambda_b}{4!} b^4 + \frac{\lambda_\lambda}{4!} \lambda^4 + \frac{\lambda_{Jb}}{4} J^2 b^2 + \frac{\lambda_{\lambda J}}{4} \lambda^2 J^2 + \frac{\lambda_{\lambda b}}{4} \lambda^2 b^2$$

3-d classical field space, using the Gildener-Weinberg approach

PGSB of SI is dominated by singlet with largest VEV \sim TeV, with mass around 100 GeV

Scale invariant ν MSM (ν SISM)

❖ FImP (keV FIMP) miracle

EWSB favors singlets with $\text{VEVs} \sim \text{TeV}$, thus a keV RHN means

$$\frac{1}{2} \lambda S N^2 \xrightarrow{S = \frac{1}{\sqrt{2}}(v + s)} \left\{ \begin{array}{l} M_N = \lambda \langle s \rangle \Rightarrow \lambda = \frac{M_N}{\langle s \rangle} \sim 10^{-8} \\ \frac{1}{2\sqrt{2}} \lambda \cdot S N^2 \end{array} \right.$$

At the same time, RHN gains a feeble interaction, which is too weak to thermalize it. However, it is **just at the correct order to freeze-in RHN:**

$$\Omega_{\text{DM}} h^2 = 0.11 \times \sum_{H_a = \mathcal{P}, H_2} \left(\frac{m_{\text{DM}}}{10 \text{ keV}} \right)^3 \left(\frac{\text{TeV}}{v_J} \right)^2 \left(\frac{100 \text{ GeV}}{m_{H_a}} \right) \left(\frac{10^3}{g_*^S \sqrt{g_*^P}} \right),$$

By scale invariance, RHN mass and dynamics accounting for relic density share the same origin——FImP miracle

Scale invariant ν MSM (ν SISM)

❖ FIMP miracle shines in the X -ray line?

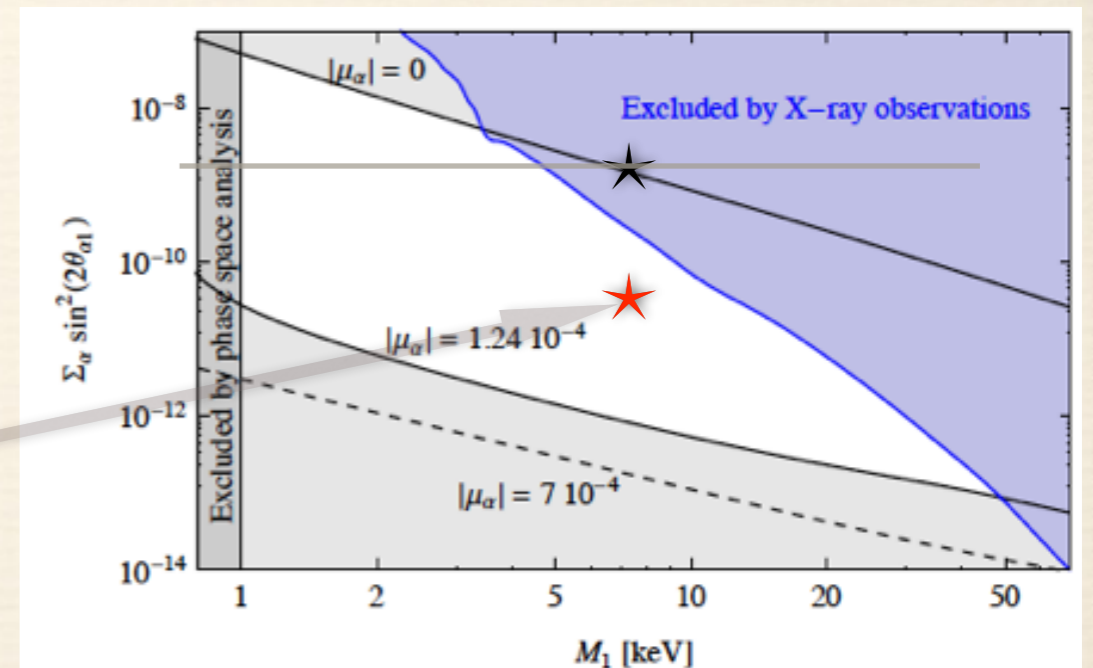
The X -ray line at 3.5 keV reported this year is explained by RHN

$$\Gamma_{\nu\gamma} \simeq 1.62 \times 10^{-28} s^{-1} \left(\frac{\sin^2 2\theta}{7 \times 10^{-11}} \right) \left(\frac{m_{\tilde{X}}}{7\text{keV}} \right)^5$$

But it has been excluded for RHN with conventional productions, even for resonant production which has been excluded by Lyman-alpha bound

A. Merle and A. Schneider, arXiv:1409.6311

The RHN from freeze-in with **colder spectrum** thus being favored!



Conclusions

- ❖ ν SISM is a good example to address three leading new physics in an unified framework
- ❖ The lightest RHN here is predicted to be an accidental DM, with FImP miracle

宇宙戰艦
台灣仁光號



Thanks...

