

## Quantum Mechanics (I): Homework 1

Due: October 5

### Ex.1

In early period of development of quantum mechanics, a few physicists (including Schrödinger himself) tried to develop so-called hidden-variable theories in which quantum fluctuations are attributed to some unknown variables in classical mechanics. In 1952, Bohm succeeded in identifying such kind of random forces in the Schrödinger equation: Suppose  $\psi$  satisfies the Schrödinger equation with the potential  $V(\mathbf{r})$ . Let  $\psi = |\psi| \exp(iS/\hbar)$ , show that

(a) **5**  $S$  is real

(b) **5** when  $\psi$  is a plane wave,  $\nabla S = \hbar \mathbf{k}$  is the momentum.

In general,  $\nabla S$  has the meaning of momentum. Let  $\mathbf{p} \equiv \nabla S$ . Here  $\mathbf{p}$  is a function of  $\mathbf{r}$  and is the momentum of the particle that arrives at  $\mathbf{r}$ . This is in the Euler description. In general, we want to follow the same particle in the so-called Lagrangian description. In this case, show that

(c) **10**  $\frac{d\vec{p}}{dt} = -\nabla(V + V_Q)$ , where  $V_Q = -\frac{\hbar^2}{2m|\psi|} \nabla^2 |\psi|$ .

$-\nabla V_Q$  is the so-called quantum mechanical force that was interpreted as the fluctuations due to the unknown hidden-variable. It is claimed that this theory can also explain the two-slit expt. (see Philippidis et al. Nuovo Cimento, **71B**, 75-87, 1982)

**Ex.2 10** Page 139, ex 4.2.2

**Ex.3 10** Page 139, ex 4.2.3

**Ex.4 10** Show that  $[p, x^n] = -ni\hbar x^{n-1}$ . In general, if  $f(x)$  is a differentiable function of  $x$ , this implies that  $[p, f(x)] = -i\hbar \frac{df(x)}{dx}$ .

**Ex.5** The state space of a certain physical system is three-dimensional. Let  $\{|u_1\rangle, |u_2\rangle, |u_3\rangle\}$  be an orthonormal basis of this space. The kets  $|\psi_1\rangle$  and  $|\psi_2\rangle$  are defined by :

$$|\psi_1\rangle = \frac{1}{\sqrt{2}}|u_1\rangle + \frac{i}{2}|u_2\rangle + \frac{1}{2}|u_3\rangle$$
$$|\psi_2\rangle = \frac{1}{\sqrt{3}}|u_1\rangle + \frac{1}{\sqrt{3}}|u_2\rangle$$

(a) **5** Are these kets normalized?

(b) **5** What is  $\langle\psi_1|\psi_2\rangle$ ?

**Ex. 6** Consider the probability amplitude

$$\Psi(x, t) = Ae^{-\lambda|x|}e^{i\omega t}$$

where  $A$ ,  $\lambda$ , and  $\omega$  are positive real constants. (a) **5** Normalize  $\Psi$ . (b) **5** Find  $\langle x \rangle$  and  $\langle x^2 \rangle$ .

**Ex. 7 10** Show that in the Schrödinger equation, the operator  $H = p^2/2m + V(x)$  is hermitian.

**Ex. 8 10** Consider two *independent* and positive classical random numbers:  $x_1$  and  $x_2$ , of which the probability densities are

$$p(x_i) = \lambda e^{-\lambda x_i}, \quad i = 1, 2,$$

where  $\lambda$  is positive and real. Let us form another random number  $x$  by defining  $x = (x_1 + x_2)/2$ . Find the probability density for  $x$ ,  $P(x)$ .

**Ex. 9 10** Suppose that at  $t = 0$ , a particle is described by the wavefunction

$$\Psi(x, 0) = \frac{1}{\sqrt{2L}} \quad |x| < L$$
$$= 0 \quad \text{otherwise.}$$

If, at the same instant, we measure the momentum of the particle. What are the possible values we will get and what are the corresponding probabilities?

**Ex. 10 10** Suppose that we do a measurement of the observable  $\hat{O}$  on some particle and get the value  $\alpha$ . Using the concept of "collapse of state", argue that after the measurement, the state of the particle has to be an eigenstate of  $\hat{O}$  with eigenvalue of  $\alpha$ .

**Ex. 11 10** Following the same argument used in the class, show that in the momentum space, the position operator  $\hat{x}$  is represented by  $i\hbar \frac{d}{dp}$ .

**Ex. 12 10** Page 46, ex.1.8.10.

**Ex.13 (a)5** Consider a particle governed by  $\hat{H} = \frac{-\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$ . Now, if the particle is further confined to be at  $x = na$  with  $n$  be integers and  $a$  being a positive constant, one needs to discretize the system: replacing any function  $f(x = na)$  by  $f_n$  and  $\left. \frac{d^2}{dx^2} \Psi(x) \right|_{x=na}$  by  $\frac{\Psi_{n+1} + \Psi_{n-1} - 2\Psi_n}{a^2}$ . Let us denote the state when the particle is localized at  $x = na$  by  $|n\rangle$ , and assume that  $\langle n|m\rangle = \delta_{nm}$ . Express  $\hat{H}$  in terms of the Dirac notations (ket and bra and etc.)  
**(b)5** In the above discrete system, the momentum operator should be changed accordingly. However, there are many possible ways to do: For example,  $\frac{\hbar}{i} \frac{d\Psi(x)}{dx} = \frac{\hbar}{i} \frac{\Psi_{n+1} - \Psi_n}{a}$  or  $\frac{\hbar}{i} \frac{d\Psi(x)}{dx} = \frac{\hbar}{i} \frac{\Psi_{n+1} - \Psi_{n-1}}{2a}$ , which one is correct? why? **(c)5** Suppose now that the particle is in the state  $|\phi\rangle \equiv |1\rangle + i|2\rangle + 2|3\rangle$ . Evaluate the averaged value of energy for  $|\phi\rangle$ .  
**(d)5** Now suppose that the particle is further confined to be at  $n = 1$  and  $n = 2$ , i.e.,  $\Psi_n = 0$  when  $n = 3, 4, 5, \dots$  and  $n = 0, -1, -2, \dots$ . The wavefunction  $|\phi(t)\rangle$  of the particle at time  $t$  can be thus written as  $|\phi(t)\rangle = \alpha(t)|1\rangle + \beta(t)|2\rangle$ , Find the differential equations that  $\alpha(t)$  and  $\beta(t)$  obey **(5)**. If at  $t = 0$ ,  $|\phi(0)\rangle = |1\rangle$ , find  $|\phi(t)\rangle$  at  $t > 0$ . **(5)**

**Bonus(+1)** Suppose that  $V(x \leq 0) = \infty$  and  $V(x \geq (N+1)a) = \infty$ ,  $V(x) = 0$  otherwise. Using the above correct expression to find the eigenvalues and eigenvectors of the discrete momentum operator. Find the corresponding expressions when  $a \rightarrow 0$ .