



國家同步輻射研究中心  
National Synchrotron Radiation Research Center

# 同步輻射在X光繞射上的應用

## Applications of Synchrotron Radiation in X-ray Diffraction

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南一中演講

Nov., 12, 2010

NSRRC



# 大綱

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## Synchrotron radiation (同步輻射)

- History and general background (背景說明)
- Properties (特性)
- Sources (bend magnets, wigglers, undulators) (光源)

## X-ray diffraction (X光繞射)

- History and general background (背景說明)
- Basics of crystal structure (晶體結構)

## Applications (應用)

- Material science (材料科學)
- Protein structure (大分子結構)
- X-ray cavity (X光共振腔--X光繞射實驗)

## TPS project (台灣光子源計畫)

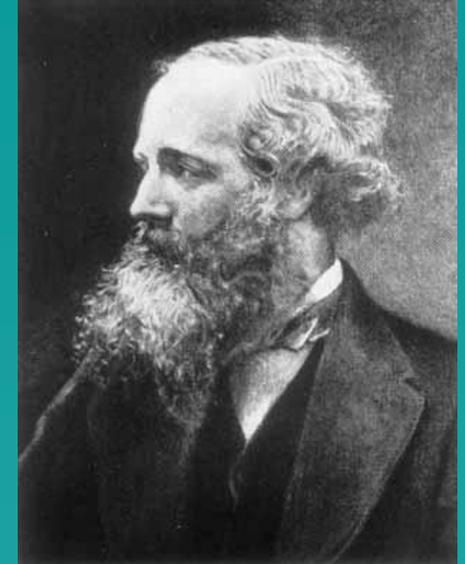
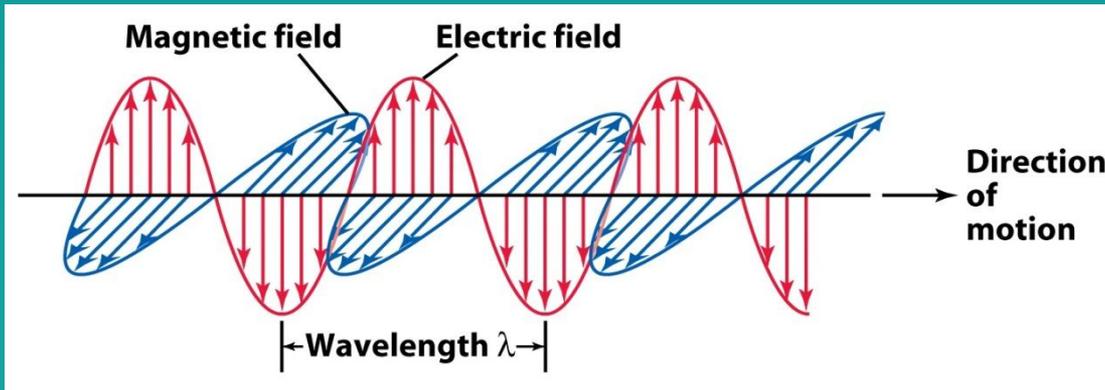
同步輻射

Synchrotron Radiation

# James Clerk Maxwell 馬克斯威爾

Light is a traveling electromagnetic wave (1862)

- Unified electromagnetism and optics
- Predicted the existence of invisible forms of light



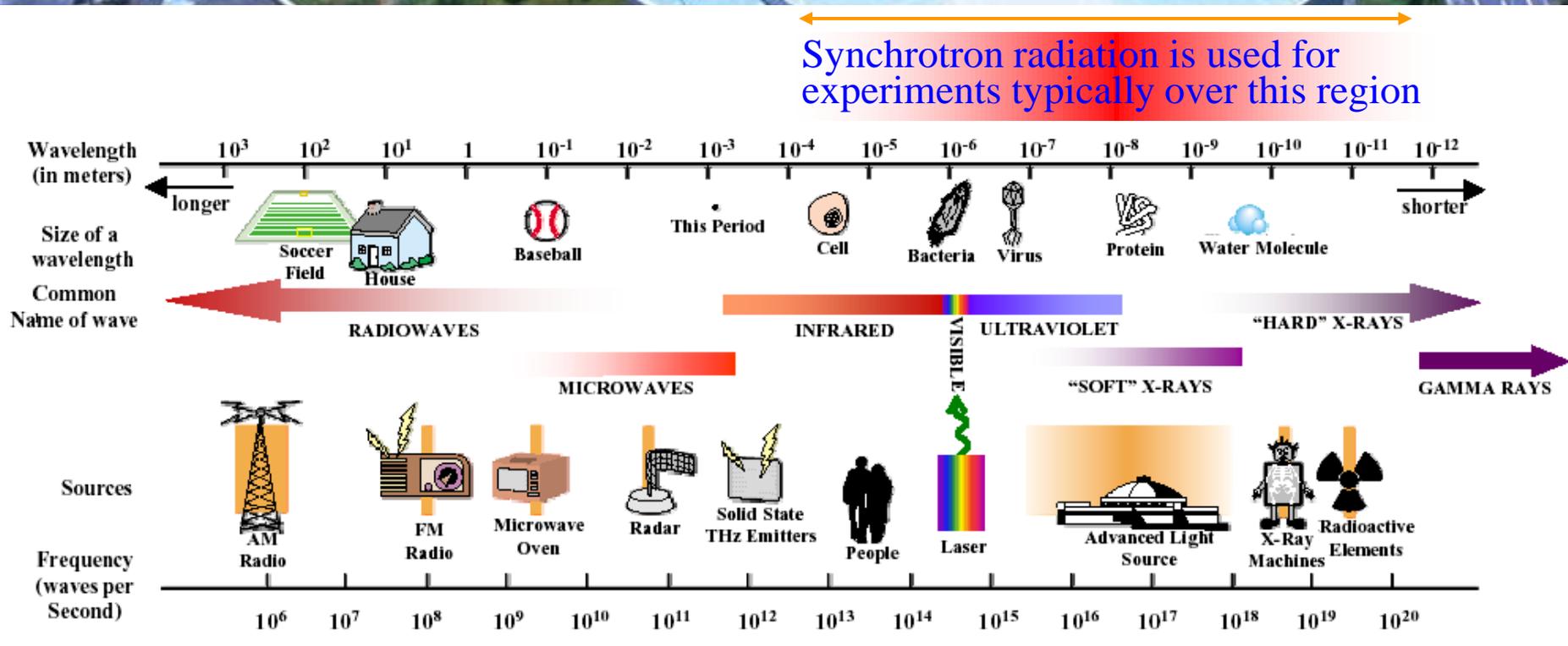
(1831-1879)  
Scottish

Maxwell equation

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon} \quad \nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad \nabla \times \mathbf{B} = \mu \mathbf{J} + \epsilon \mu \frac{\partial \mathbf{E}}{\partial t}$$

# The Electromagnetic Spectrum



# Discovery of X-rays



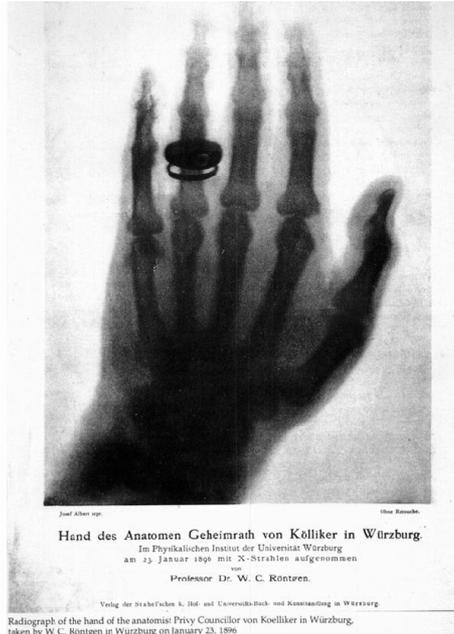
## The Nobel Prize in physics in 1901

“in recognition of the extraordinary services he has rendered by the discovery of the remarkable rays subsequently named after him”

### Wilhelm Conrad Röntgen

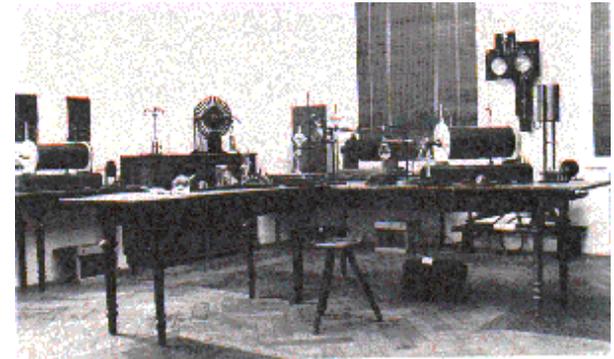


Germany (1845-1923)



Röntgen used photographs of his wife's hand to publicize his discovery.

They appeared in newspapers of the time and captured the public's imagination.



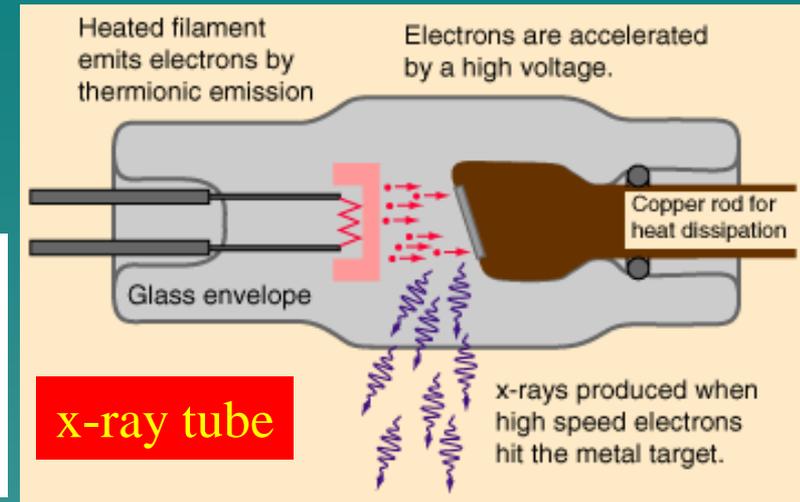
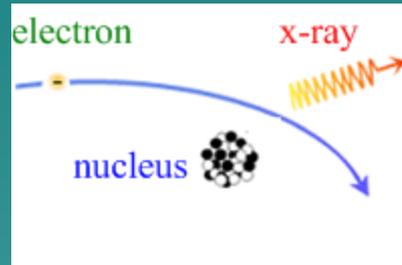
Roentgen's Laboratory  
in Würzburg, Germany

Röntgen performed experiments with cathode rays (electron beams) in 1895.

# How to Generate X-rays

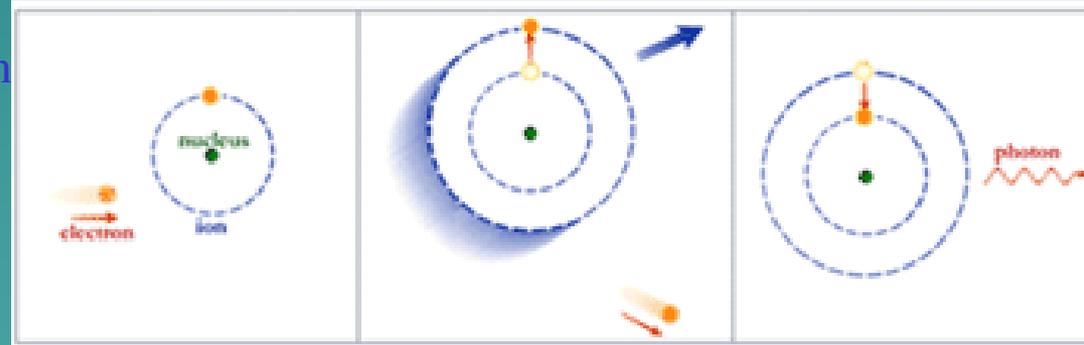
## 1. Bremsstrahlung (braking radiation)

- electrons from an external source are deflected around the nucleus of an atom
- various wavelengths



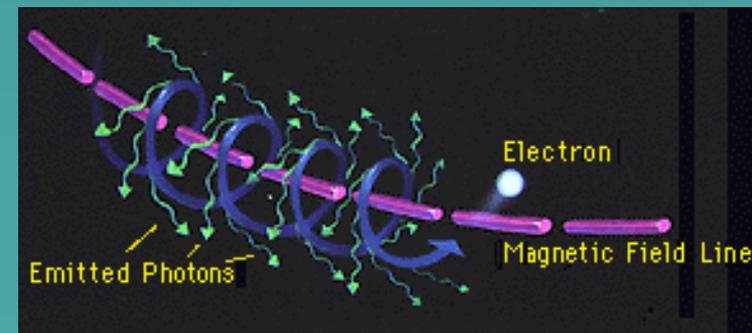
## 2. K-shell emission

- electrons change orbits within atom
- higher-intensity x-rays than bremsstrahlung, and a single wavelength.



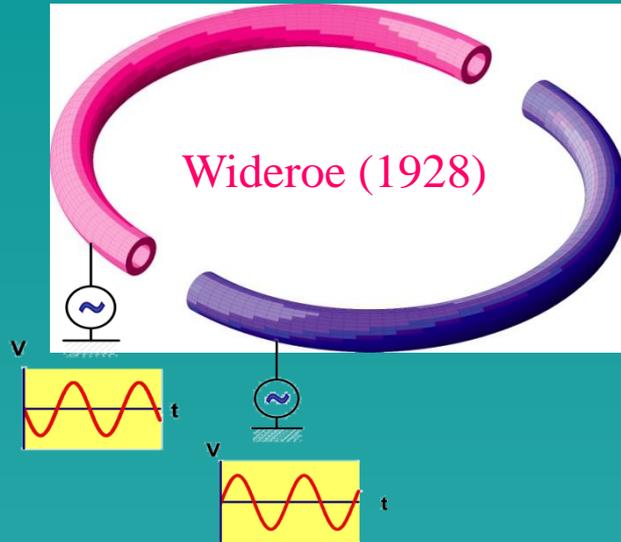
## 3. Accelerator: Synchrotron

Relativistic electron spirals around a magnetic field line

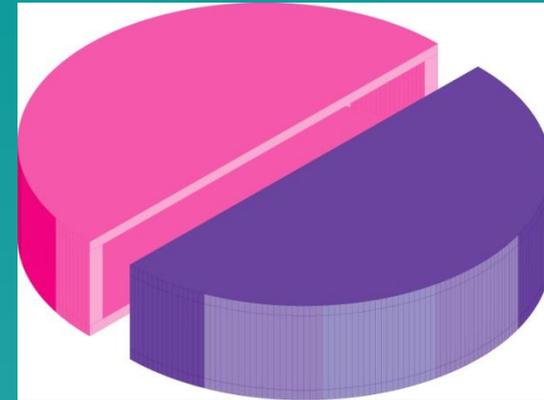


# From Wideroe linac to cyclotron and synchrotron

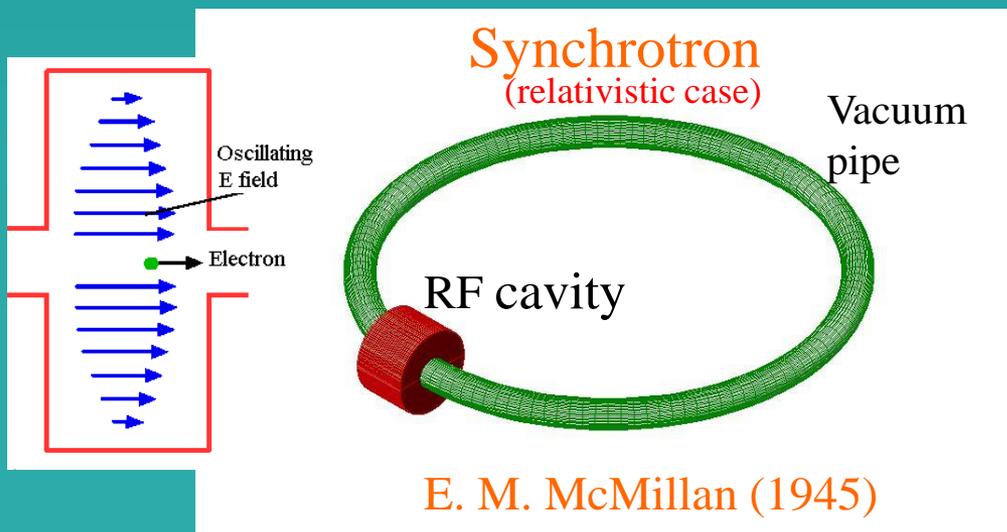
## Alternating current Accelerator



## Cyclotron (non-relativistic case)



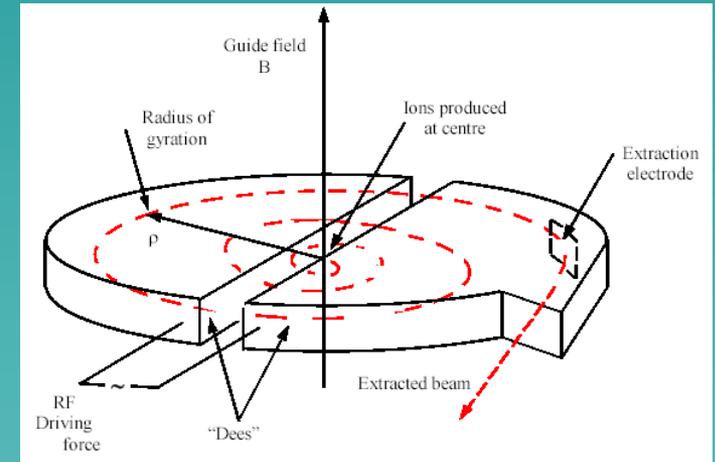
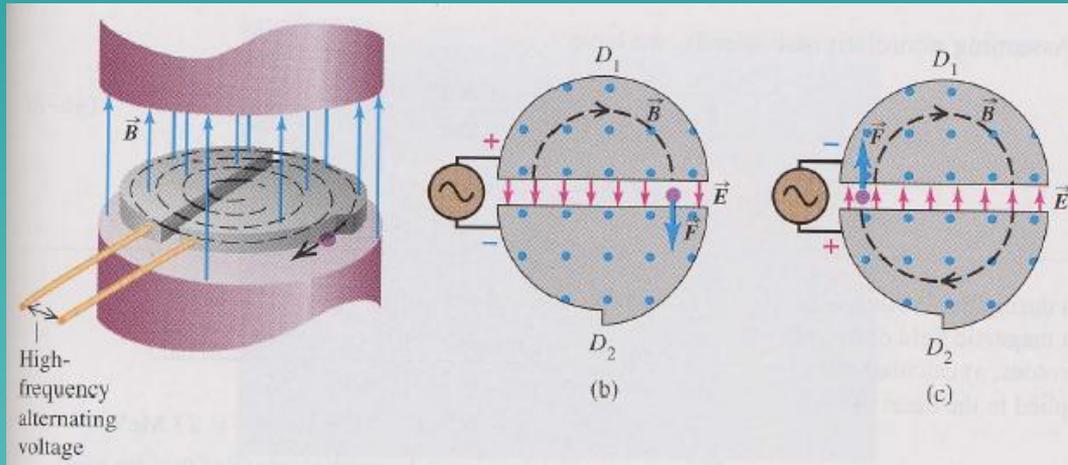
E. O. Lawrence (1932)



To keep the radius constant during acceleration, the strength of the magnetic field  $B$  has to increase **synchronously** with the **beam energy**  $E$ ,  $B \propto E$ . Therefore, this type of **accelerator** is called a **synchrotron**.

# Cyclotron (The earliest accelerators) 迴旋加速器

Cyclotron consists of two D-shaped objects (dees) with a potential difference between them. A stream of particles moved in a plane perpendicular to a uniform magnetic field, which bent the particle tracks so that they passed through an electric field in a gap between dees and were accelerated.

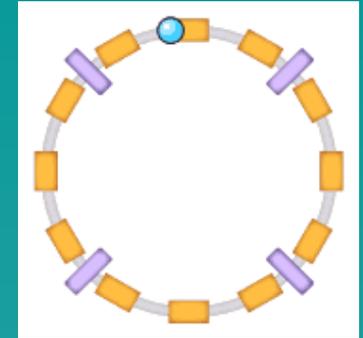


For **non-relativistic** particles the frequency of the machine was determined by the **Lorentz force law**,  $F = e v B$ , and the formula for centripetal acceleration,  $v^2 / r = F / m = e v B / m$ , so the angular frequency is given by:  $\omega = e B / m$ .

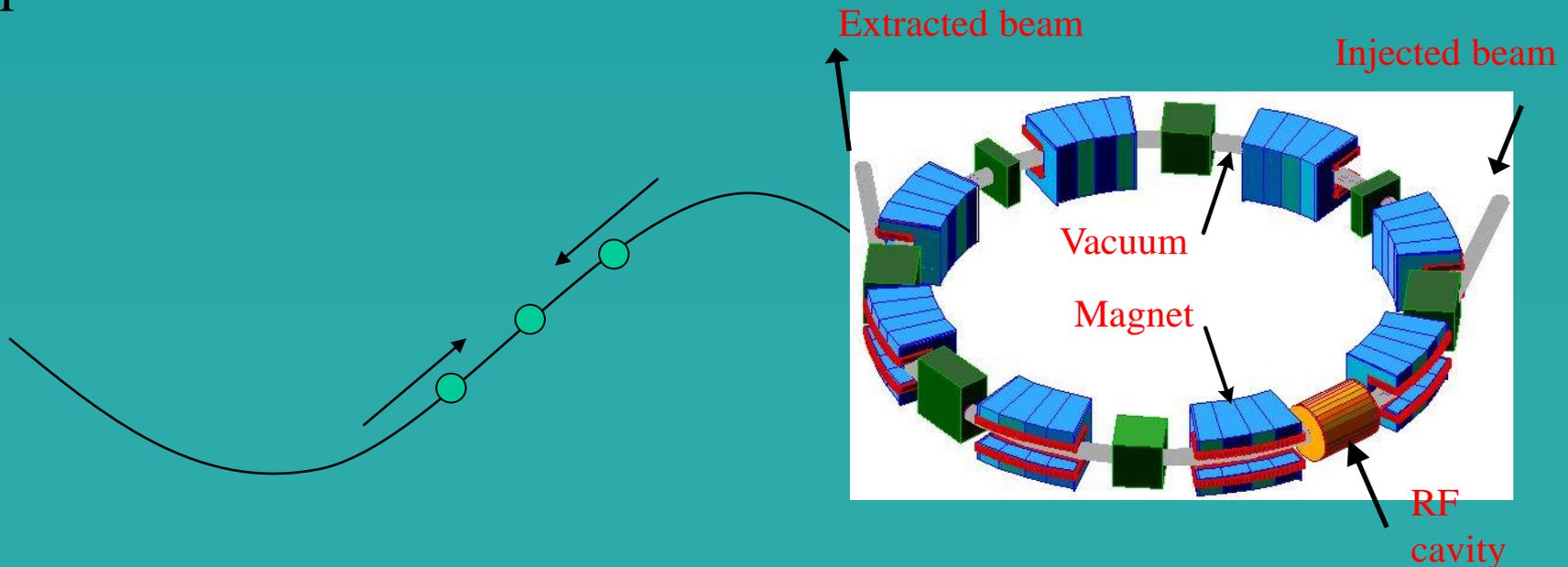
# Synchrotron 同步加速器

- Particles “ride” an RF wave

- Slower particles get pushed
- Faster get slowed
- Leads to synchrotron oscillations
  - However, “bunches” are formed
- Transverse instabilities lead to “betatron” oscillations

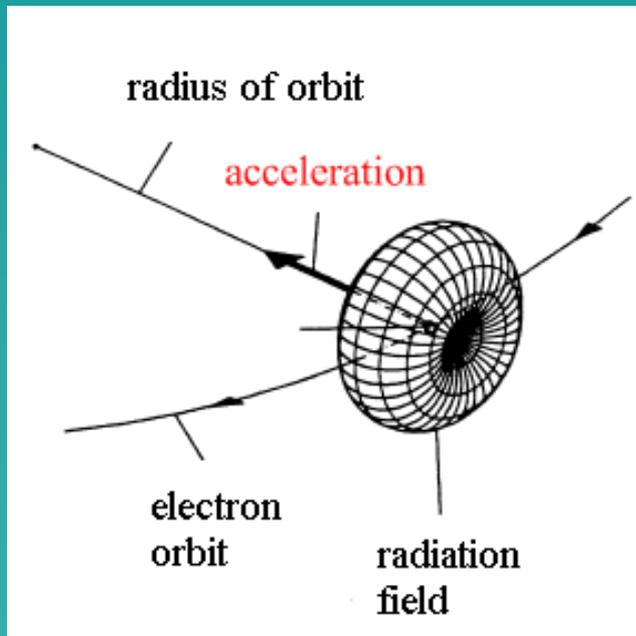


- Magnets are used to steer the beams during acceleration and during “fill”

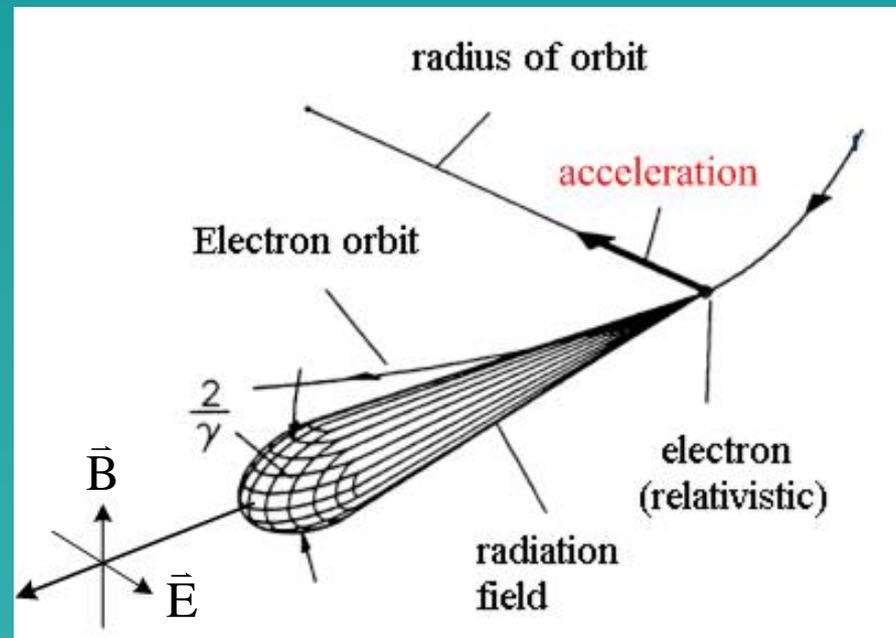


# Radiation Fundamental

- When electrons are accelerated (e.g. linear acceleration in a radio transmitter antenna) they emit electromagnetic radiation (*i.e.*, radio waves) in a rather non-directional pattern
- Electrons in circular motion are also undergoing acceleration (centripetal)



At low electron velocity (non-relativistic case) the radiation is emitted in a non-directional pattern



When the electron velocity approaches the velocity of light, the emission pattern is folded sharply forward. Also the radiated power goes up dramatically

# Fields of a Accelerated Charge

$$\vec{E}(\vec{r}, t) = \frac{q}{4\pi\epsilon_0} \left[ \frac{\vec{n} - \vec{\beta}}{(1 - \vec{n} \cdot \vec{\beta})^3 \gamma^2} \cdot \frac{1}{r^2} \right]$$

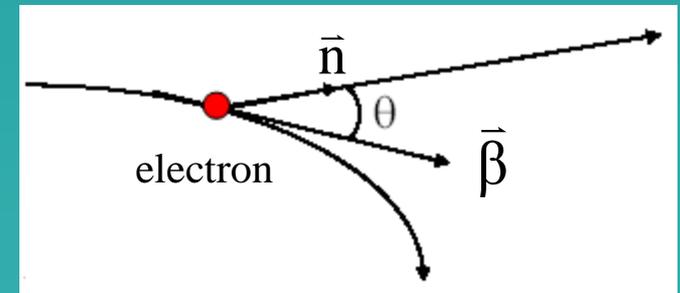
Coulomb field

$$+ \frac{q}{4\pi\epsilon_0 c} \left\{ \frac{\vec{n} \times [(\vec{n} - \vec{\beta}) \times \dot{\vec{\beta}}]}{(1 - \vec{n} \cdot \vec{\beta})^3 \gamma^2} \cdot \frac{1}{r} \right\}$$

Radiation field

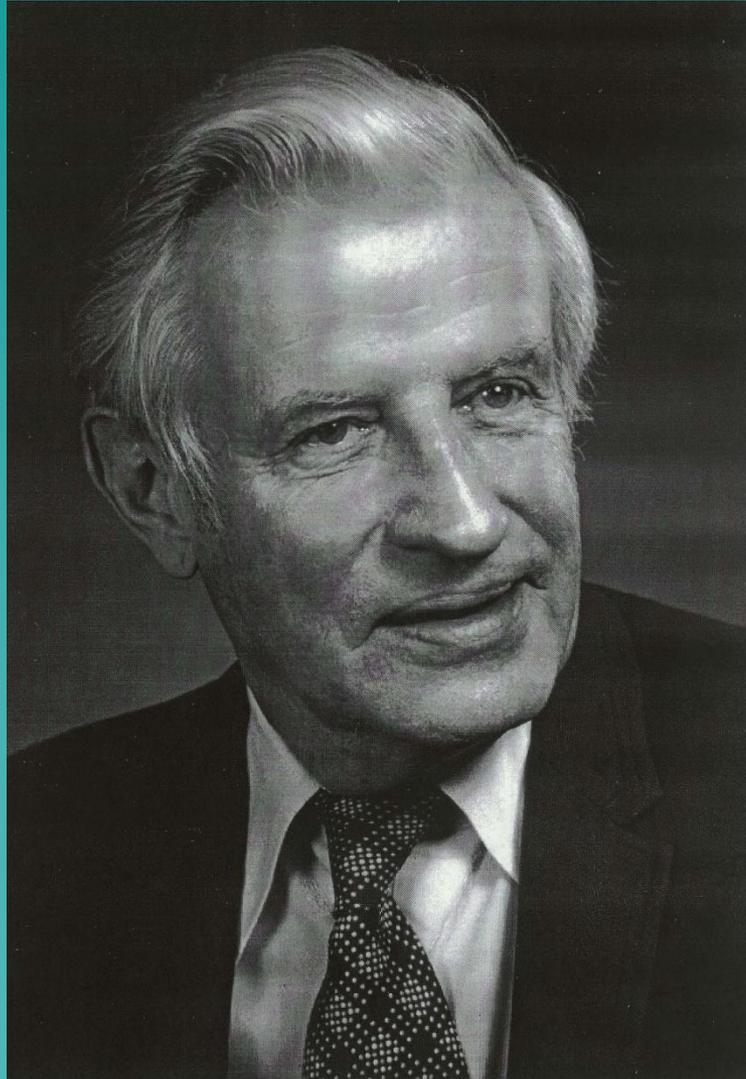
$$\vec{B}(\vec{r}, t) = \frac{1}{c} [\vec{n} \times \vec{E}(\vec{r}, t)]$$

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}}$$



$$\vec{n} = \frac{\vec{r}}{r}, \quad \vec{\beta} = \frac{\vec{v}}{c}$$

John Blewett Observed Effects of Synchrotron Radiation in  
General Electric 100 MeV Betatron - 1945



## Radiation Losses in the Induction Electron Accelerator

JOHN P. BLEWETT

*Research Laboratory, General Electric Company, Schenectady, New York*

(Received September 13, 1945)

This paper discusses the possibility that radiation losses because of the high radial accelerations experienced by the electrons in an induction electron accelerator may introduce limitations in the design of accelerators for energies above 100 million electron volts. The effects of radiation losses on the electron orbits are calculated, and it is shown that not only should the orbit shift pulse necessary to bring electrons to a target inside the equilibrium orbit fall below the value expected in the absence of radiation, but also electrons should eventually arrive at the target with no orbit shift pulse whatever, at a phase of the field wave predictable from the theory. Both effects have been observed in the General Electric 100-Mev unit in a manner consistent with the predictions of the theory. The radiation itself has not yet been detected.

### 1. INTRODUCTION\*

**I**N the induction electron accelerator, the electrons are subjected continually to radial accelerations of the order of  $10^{17}$  meters per

second per second. It has been pointed out by

$B_r$  and  $B_z$  are components of magnetic flux density  
(webers per sq. m)

$c$  = velocity of light =  $3.00 \times 10^8$  m per sec.

$e$  = charge on the electron =  $1.602 \times 10^{-19}$  Coulomb

$E_r$  and  $E_z$  are components of electric field (volts per m)

$f_n$  and  $f_t$  are normal and tangential components of the  
acceleration vector  $\mathbf{f}$  (m per sec. per sec.)

$F(\omega t) = (\omega t / \sin \omega t) - \cos \omega t - (2/3) \sin^2 \omega t \cos \omega t$

$h$  = Planck's constant =  $6.624 \times 10^{-34}$  joule sec.

$H_r$  and  $H_z$  are components of magnetic field

$I$  = beam current (amperes)

$m_0$  = rest mass of the electron =  $9.107 \times 10^{-31}$  kg

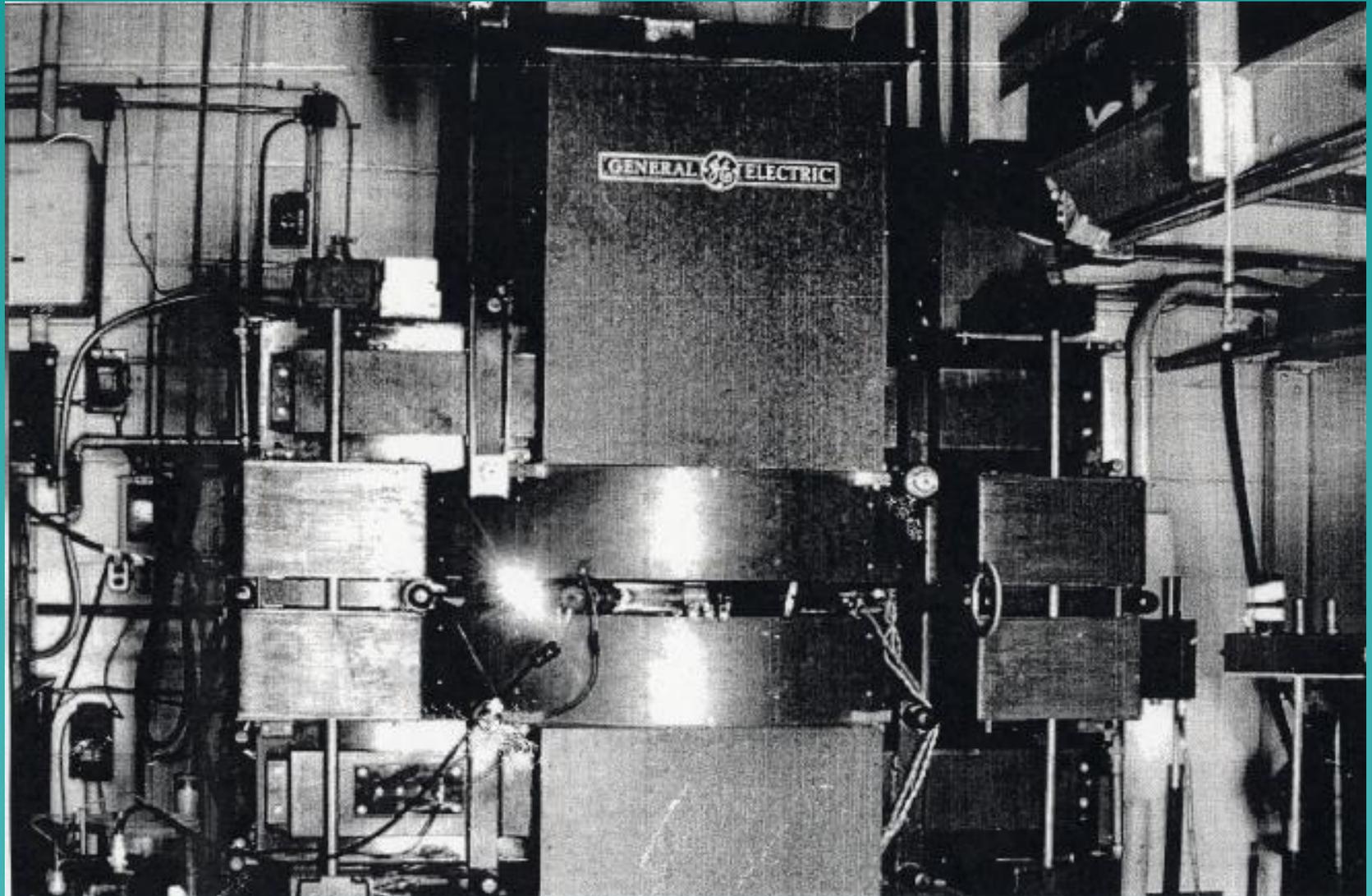
\* *Symbols*:—Unrationalized m.k.s. units will be used throughout: The following symbols will be employed:

$A$  = peak value of applied magnetic flux density at the  
equilibrium orbit (webers per sq. m)

$A'$  = peak value of magnetic flux in orbit shrinking pulse  
at the equilibrium orbit (webers per sq. m)

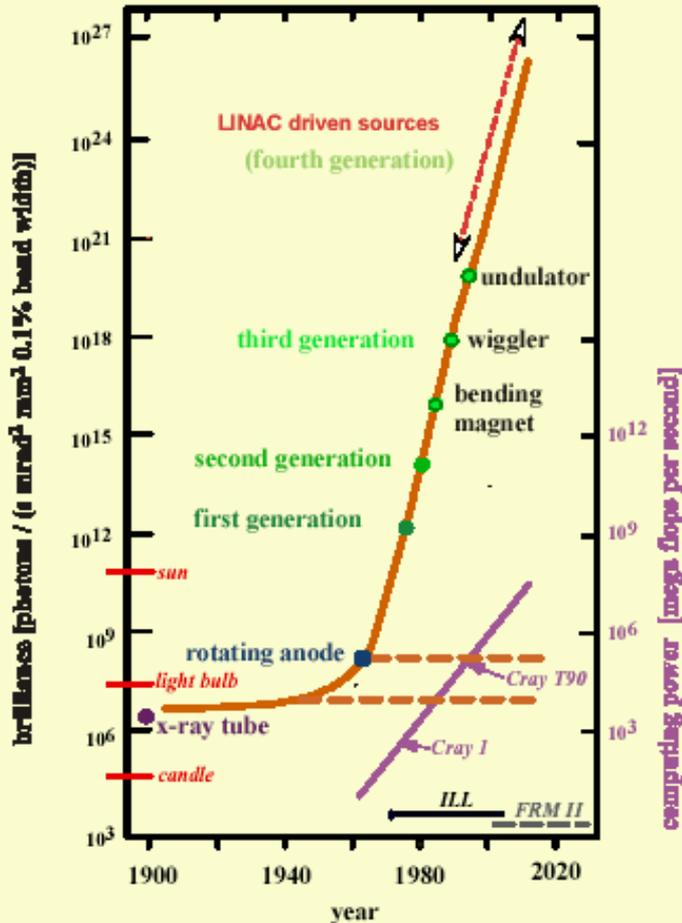
$B_0$  = applied magnetic flux density at the equilibrium  
orbit (webers per sq. m)

**Synchrotron Radiation** was first observed in 1947 by Elder, Gurewitsch, Langmuir and Pollock in a 70MeV synchrotron.



# History of Synchrotron Radiation

The Brilliance of X-ray Sources since their Discovery in 1895



## Zeroth generation sources

1950's-60's: Electron synchrotrons (cyclic accelerators)

## First generation sources (storage rings)

1970's: e<sup>+</sup>/e<sup>-</sup> colliders (Mostly parasitic on high energy physics programs)

## Second generation sources

1980's: New rings and fully dedicated use of e<sup>+</sup>/e<sup>-</sup> colliders, use of wigglers & undulators

## Third generation sources

1990's: Low emittance ring with many straight sections for insertion devices

## Fourth generation sources

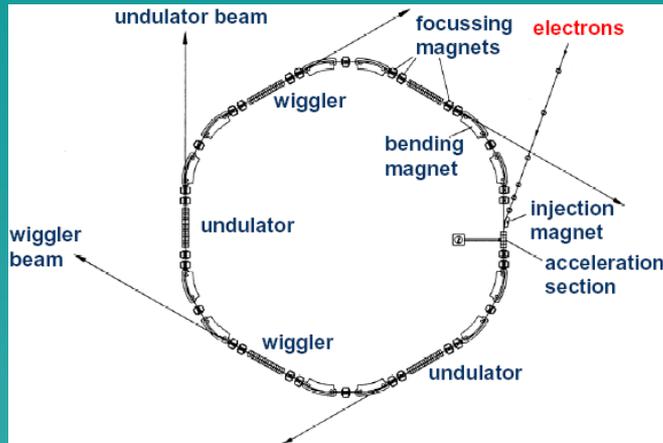
2000's: Linac-based sources

- Free-electron laser (FEL)
- Energy Recovery Linac (ERL)

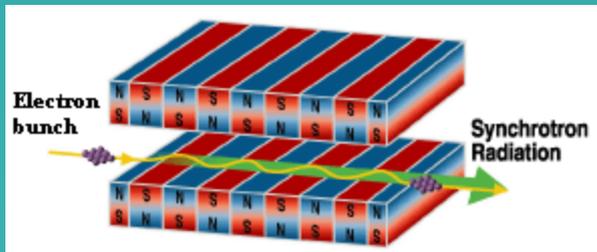
Diffraction-limited rings; Ultra-short bunches; New ideas

# Bending magnet(偏轉磁鐵) & insertion device(插件)

## Storing Ring

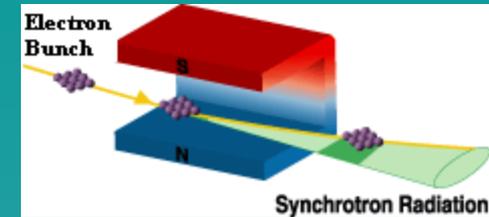


## Undulator / Wiggler



## • Bending Magnet

- White X-rays
- Wide horizontal divergence
- $1/\gamma$  limited vertical divergence
- Moderate power
- Moderate power density



## • Wiggler

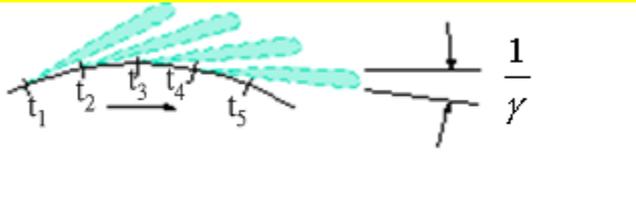
- White X-rays
- Moderate horizontal divergence
- $1/\gamma$  Limited vertical divergence
- High power
- High power density
- Elliptically polarized/linearly polarized

## • Undulator

- Quasi-monochromatic X-rays
- Small vertical and horizontal divergence (Central Cone)
- High power
- Extremely high power density
- Circularly polarized/ linearly polarized

# Three forms of Synchrotron Radiation

## Bending Magnet: a sweeping searchlight

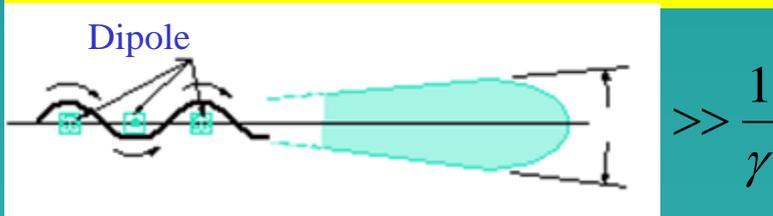


## Continuous spectrum

Critical energy  $E_c = \hbar\omega_c = \frac{3e\hbar B\gamma^2}{2m}$

$E_c(\text{keV}) = 0.665 B(\text{T})E^2(\text{GeV})$

## Wiggler: incoherent superposition



eg: for  $B = 1.35\text{T}$ ,  $E = 2\text{GeV}$

$E_c = 3.6\text{keV}$

增頻磁鐵

$B$  increases  $\rightarrow E_c$  increases

## Undulator: coherent superposition



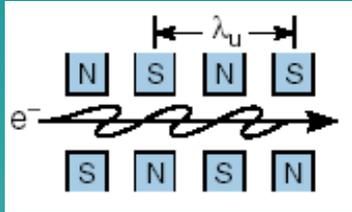
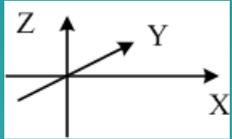
Quasi-monochromatic spectrum  
with peaks at lower energy than a  
wiggler

聚頻磁鐵

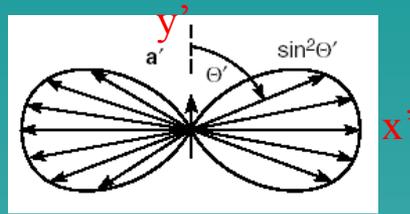
$$\lambda = \frac{\lambda_u}{2\gamma^2} \left( 1 + \frac{K^2}{2} + \gamma^2\theta^2 \right)$$

# Undulator radiation

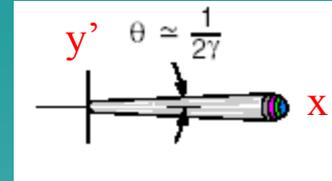
Laboratory  
Frame of Reference



Frame  
of Moving Electron



Laboratory  
Frame of Reference



Lorentz  
transformation

$$E = \gamma mc^2$$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Electron radiates at the  
Lorentz contracted  
wavelength

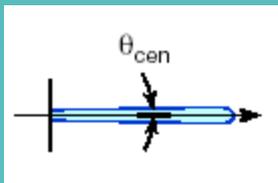
$$\lambda' = \frac{\lambda_u}{\gamma}$$

Doppler shortened  
Wavelength on axis

$$\lambda = \lambda' \gamma (1 - \beta \cos \theta)$$

$$= \frac{\lambda_u}{2\gamma^2} (1 + \gamma^2 \theta^2)$$

Following  
Monochromator



$$\theta_{cen} \cong \frac{1}{\gamma \sqrt{N}}$$

Bandwidth  $\frac{\lambda'}{\Delta \lambda'} \cong N$

Accounting for transverse  
Motion due to the periodic  
Magnetic field

$$\lambda = \frac{\lambda_u}{2\gamma^2} \left( 1 + \frac{K^2}{2} + \gamma^2 \theta^2 \right)$$

$$K = \frac{eB_0 \lambda_u}{2\pi mc}$$

# Synchrotron radiation - basic properties

1. High flux and brightness

2. High stability

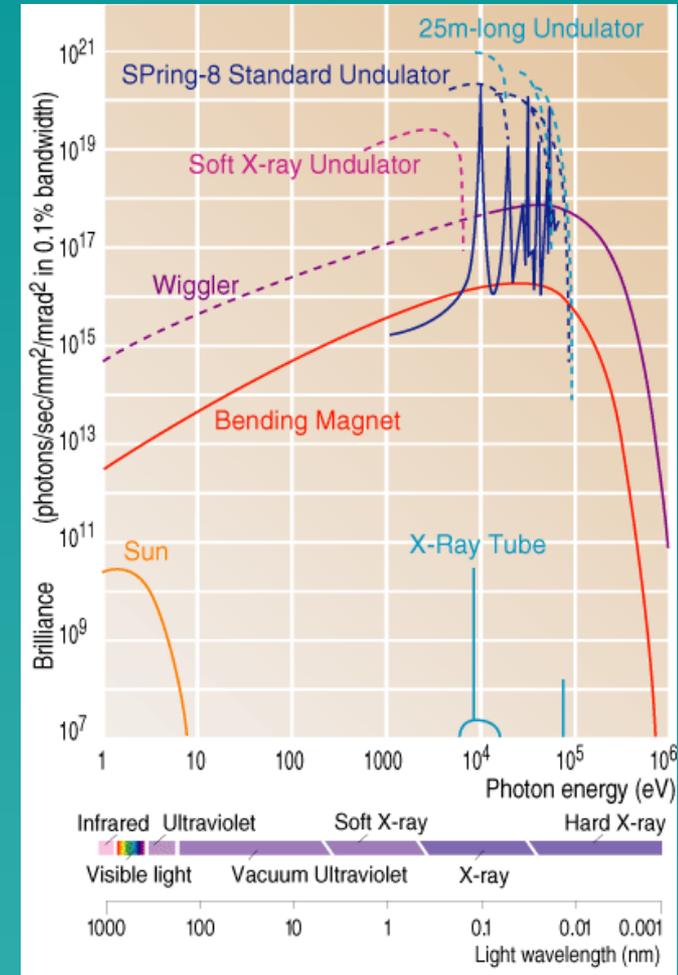
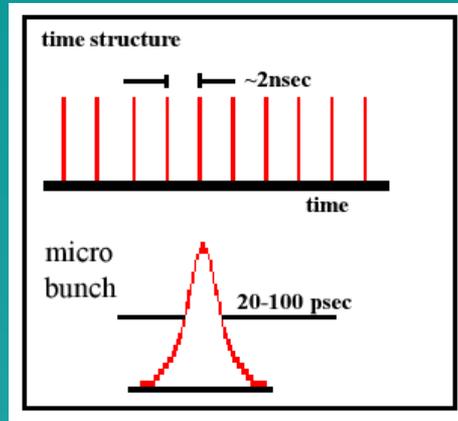
3. Broad spectral range

4. Small source size

5. Partial coherence

6. Polarized (linear, elliptical, circular)

7. Pulsed time structure



$$\text{Flux} = \frac{\# \text{ of photons in given } \Delta\lambda/\lambda}{\text{sec, mrad } \theta}$$

$$\text{Brightness} = \frac{\# \text{ of photons in given } \Delta\lambda/\lambda}{\text{sec, mrad } \theta, \text{ mrad } \varphi, \text{ mm}^2}$$

# Temporal and Spatial coherence 同調性

Mutual coherence factor

$$\Gamma_{12}(\tau) \equiv \langle E_1(t)E_2^*(t+\tau) \rangle$$

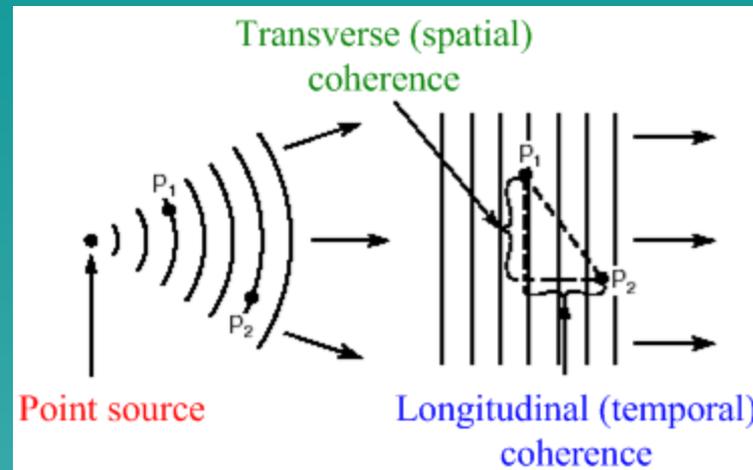
Normalized degree of spatial coherence

$$\gamma_{12}(\tau) = \frac{\langle E_1(t)E_2^*(t+\tau) \rangle}{\sqrt{\langle |E_1(t)|^2 \rangle} \cdot \sqrt{\langle |E_2(t)|^2 \rangle}}$$

Incoherent  $\gamma_{12} = 0$

Partially coherent  $\gamma_{12} < 1$

Fully coherent  $\gamma_{12} = 1$



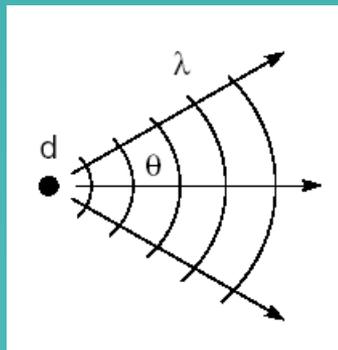
Longitudinal coherence length

$$l_{coh} = \frac{\lambda^2}{2\Delta\lambda}$$

Transverse coherence length ( $\theta$ )

$$d \cdot \theta = \frac{\lambda}{2\pi} \quad (\text{The uncertainty principle})$$

**d: source size**



# NSRRC

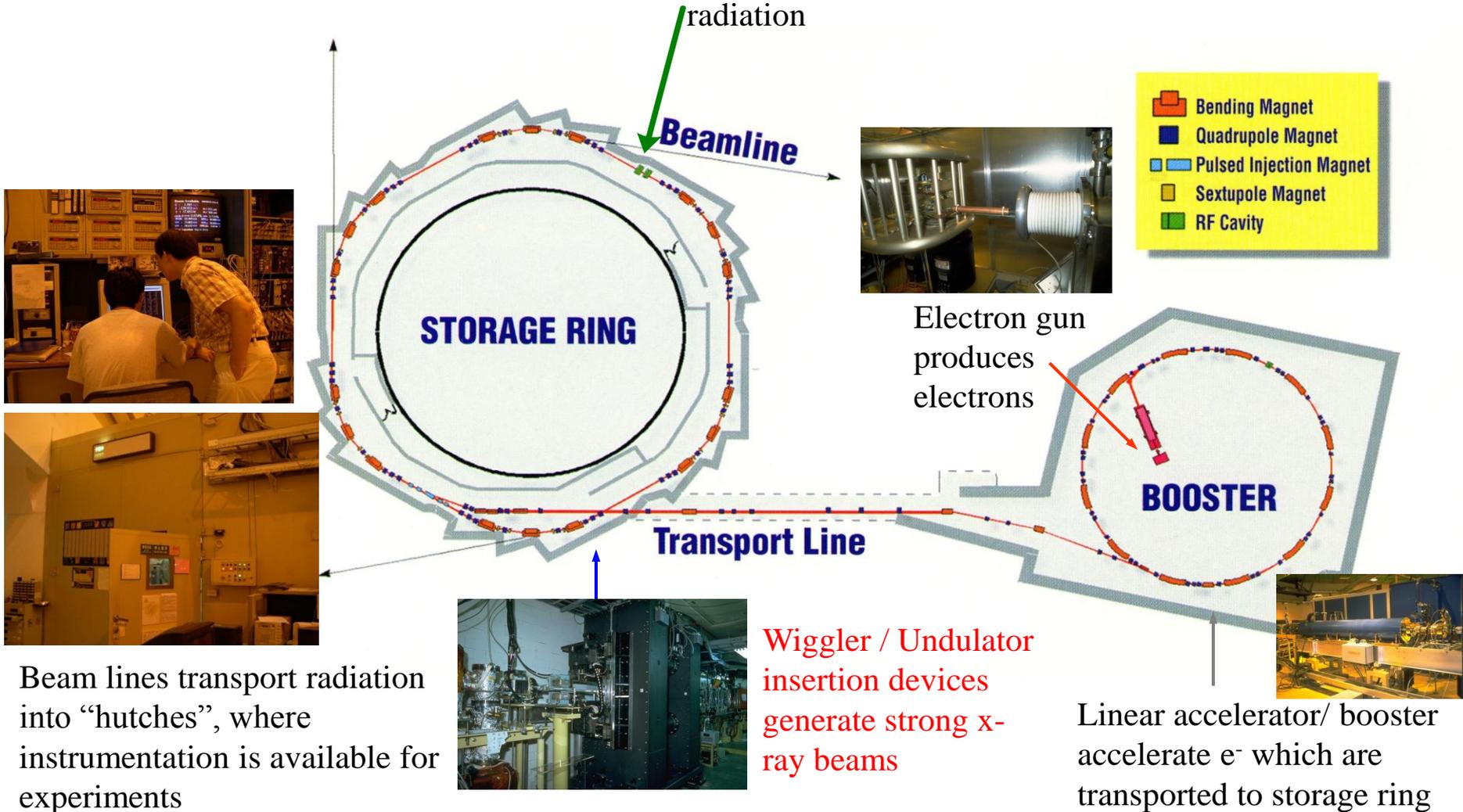


1.5 GeV;  $\gamma = 3000$ ; 120 m circumference

# How is it Practically Produced and Used for Research?

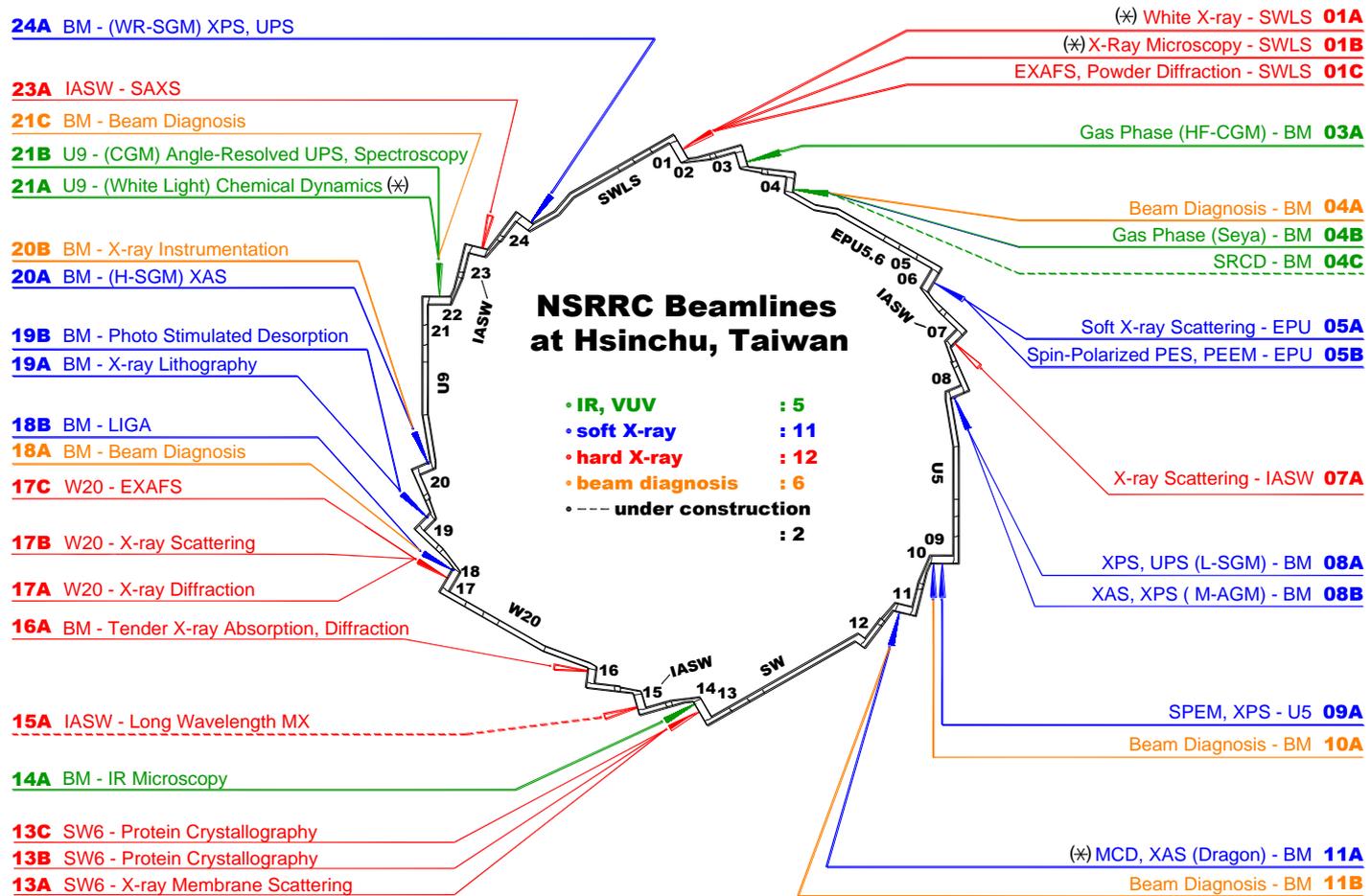
The storage ring circulates electrons, where they are bent, synchrotron radiation is produced

Klystrons generate high power radio wave to sustain electron acceleration, replenishing energy lost to synchrotron radiation

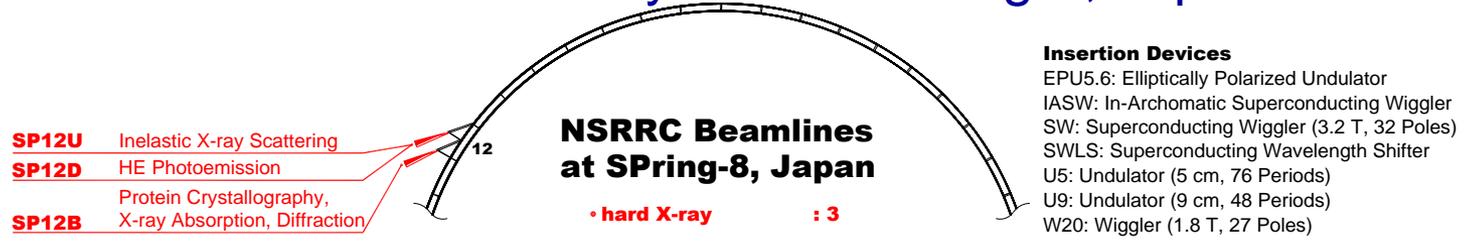


# NSRRC TLS Beamlines Layout

30 BLs, over 50 End stations



## 3 additional hard x-ray BLs at SPring-8, Japan

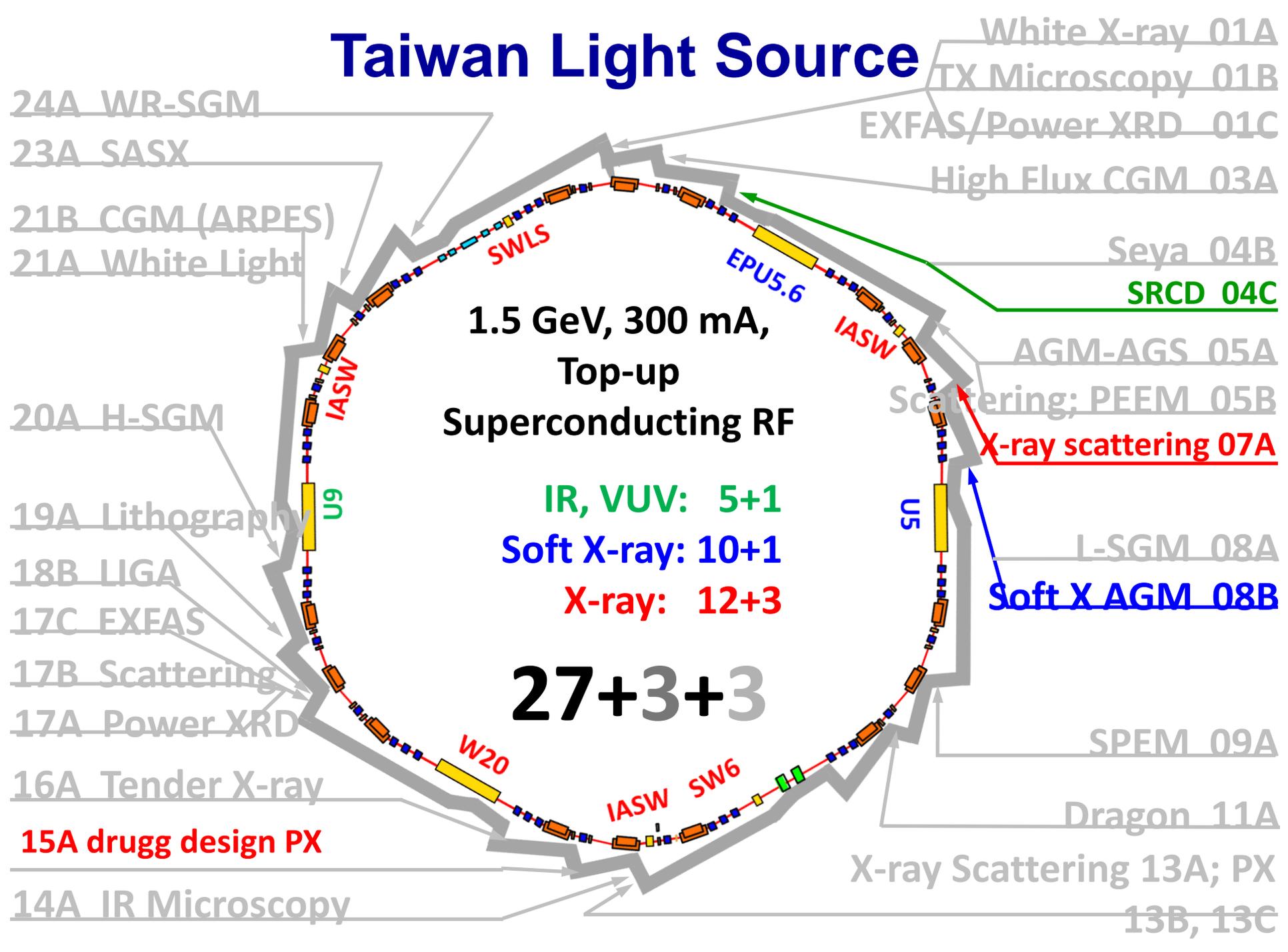


**Insertion Devices**

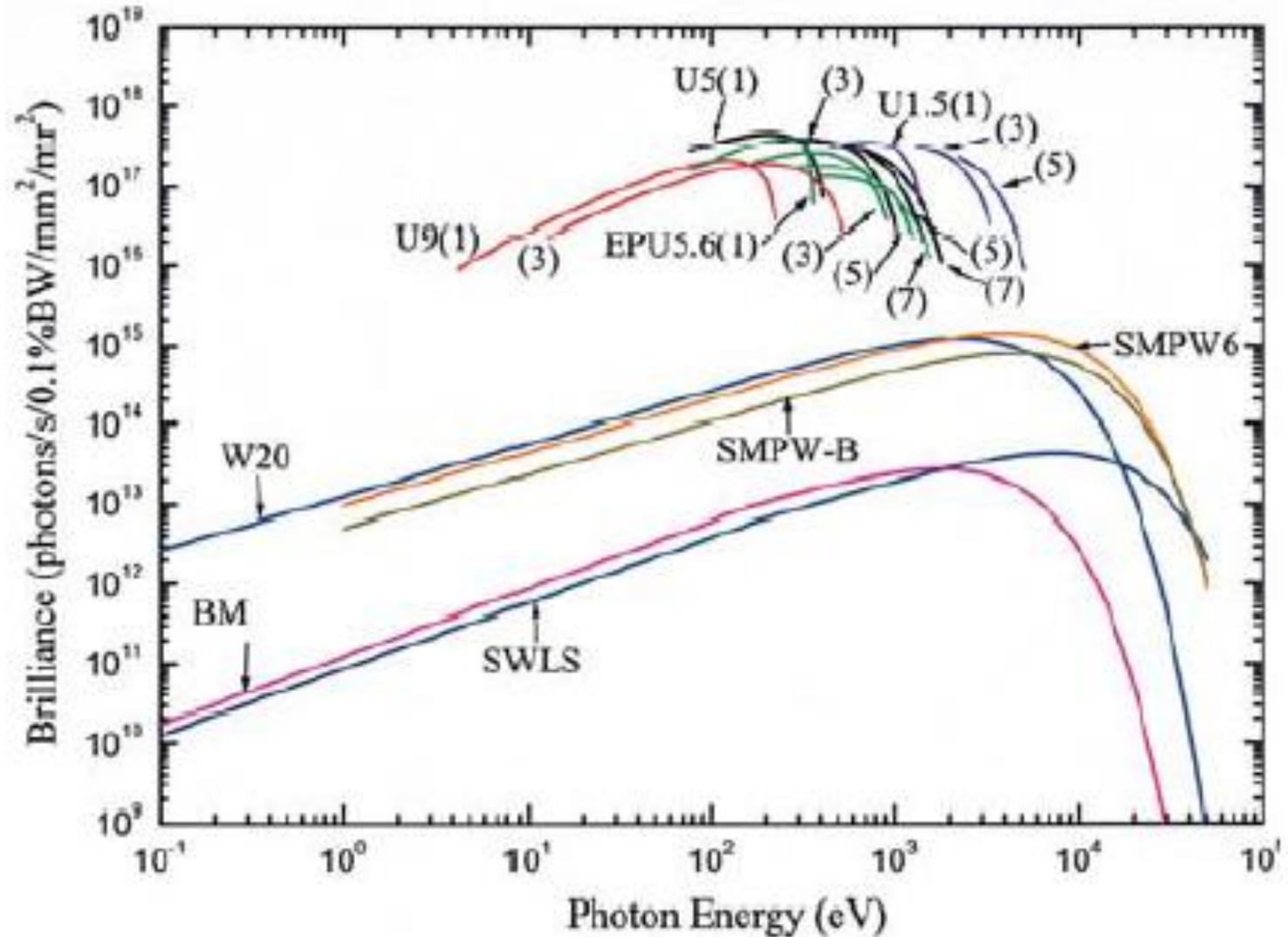
- EPU5.6: Elliptically Polarized Undulator
- IASW: In-Archromatic Superconducting Wiggler
- SW: Superconducting Wiggler (3.2 T, 32 Poles)
- SWLS: Superconducting Wavelength Shifter
- U5: Undulator (5 cm, 76 Periods)
- U9: Undulator (9 cm, 48 Periods)
- W20: Wiggler (1.8 T, 27 Poles)

(\*) : Participating Research Group

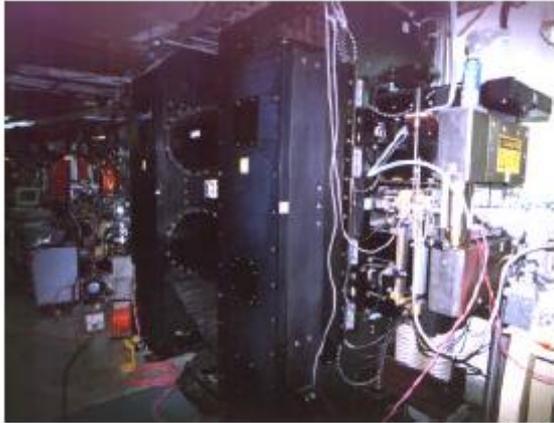
# Taiwan Light Source



# NSRRC Photon Spectrum from Insertion Devices



# NSRRC Insertion Devices



**W20 wiggler**



**U10 undulator**



**U5 undulator**



**U9 undulator**



**EPU5.6 undulator**

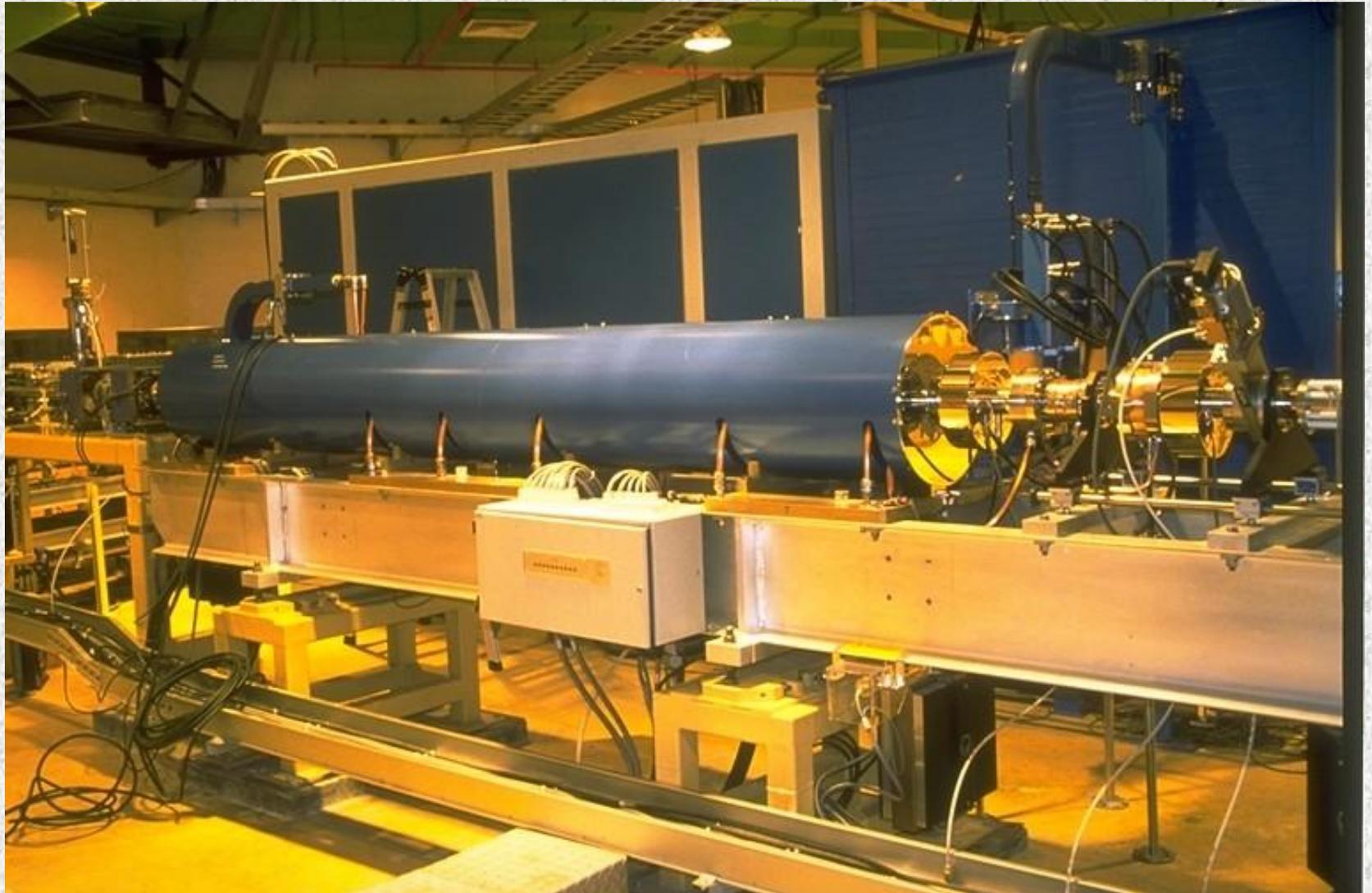


**IVXU3.2 undulator**

# Parameters of the Insertion Devices

	W20	U10	U5	U9	EPU5.6	SWLS	SW6
Type	Hybrid	Hybrid	Hybrid	Hybrid	Pure	SC	SC
$\lambda$ [cm]	20	10	5	9	5.6	32.56	6
Gap [mm]	22	22	18	18	18	56	18
N	13	20	76	48	66	1.5	32
Photon energy (keV)	4-15	----	0.06-1.5	0.004-0.5	0.06-1.5	4-33	6.5-19
Deflection K	33.6	9.34	2.99	10.5	3.5 (2.35)	----	17.9
$B_{\max}$ [T]	1.8	1.0	0.64	1.25	0.67 (0.45)	6	3.2
Phase error $\Theta$	N/A	2.5°	3°	3.7°	5.5°	N/A	N/A
$\theta_x(\theta_y)$ [ $\mu$ rad]	N/A	15 (6)	6 (11)	20 (10)	10 (15)	N/A	N/A
$\delta_x(\delta_y)$ [ $\mu$ m]	N/A	5 (2)	7 (10)	8 (2)	2 (1)	N/A	N/A
Beam duct aperture(mm <sup>2</sup> )	17x80	17x80	13x80	13x80	13x80	20x100	11x80
Installation	Dec.1994	Oct.1995	Mar.1997	Apr.1999	Sep.1999	Apr. 2002	Jan. 2004

# NSRRC Linear Accelerator



50 Mev

# NSRRC Booster Synchrotron



# Focusing Magnet



# Synchrotron Radiation Centres around the world



SR sources around the world

[http://www.spring8.or.jp/ENGLISH/general\\_info/overview/sr.html](http://www.spring8.or.jp/ENGLISH/general_info/overview/sr.html)

# Synchrotron Radiation Centers in USA



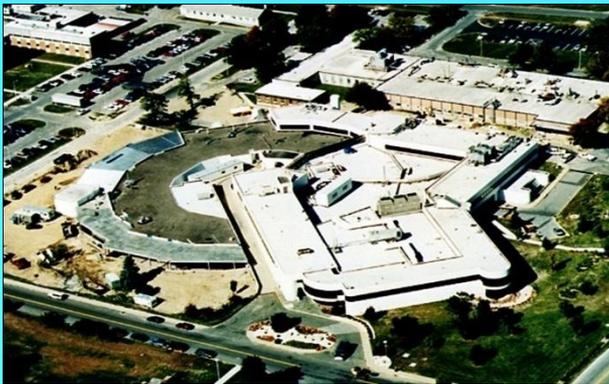
1.9 GeV,  $\gamma = 3720$ , 197 m circumference

Advanced Light Source (ALS),  
Lawrence Berkeley National Laboratory (1993)

Two US 3rd Generation  
facilities - among the  
world's first



Advanced Photon Source (APS),  
Argonne National Laboratory (1996)



National Synchrotron Light Source (NSLS),  
Brookhaven National Laboratory (1982)

First US 2nd Generation facility



Stanford Synchrotron Radiation Laboratory (SSRL),  
Stanford Linear Accelerator Center (1974)

World's first SR storage ring x-ray user facility

# European Synchrotron Radiation Facility (ESRF)



6 GeV;  $\gamma = 11800$ ; 884m circumference

# Spring-8 (Super Photon ring-8 GeV) Japan

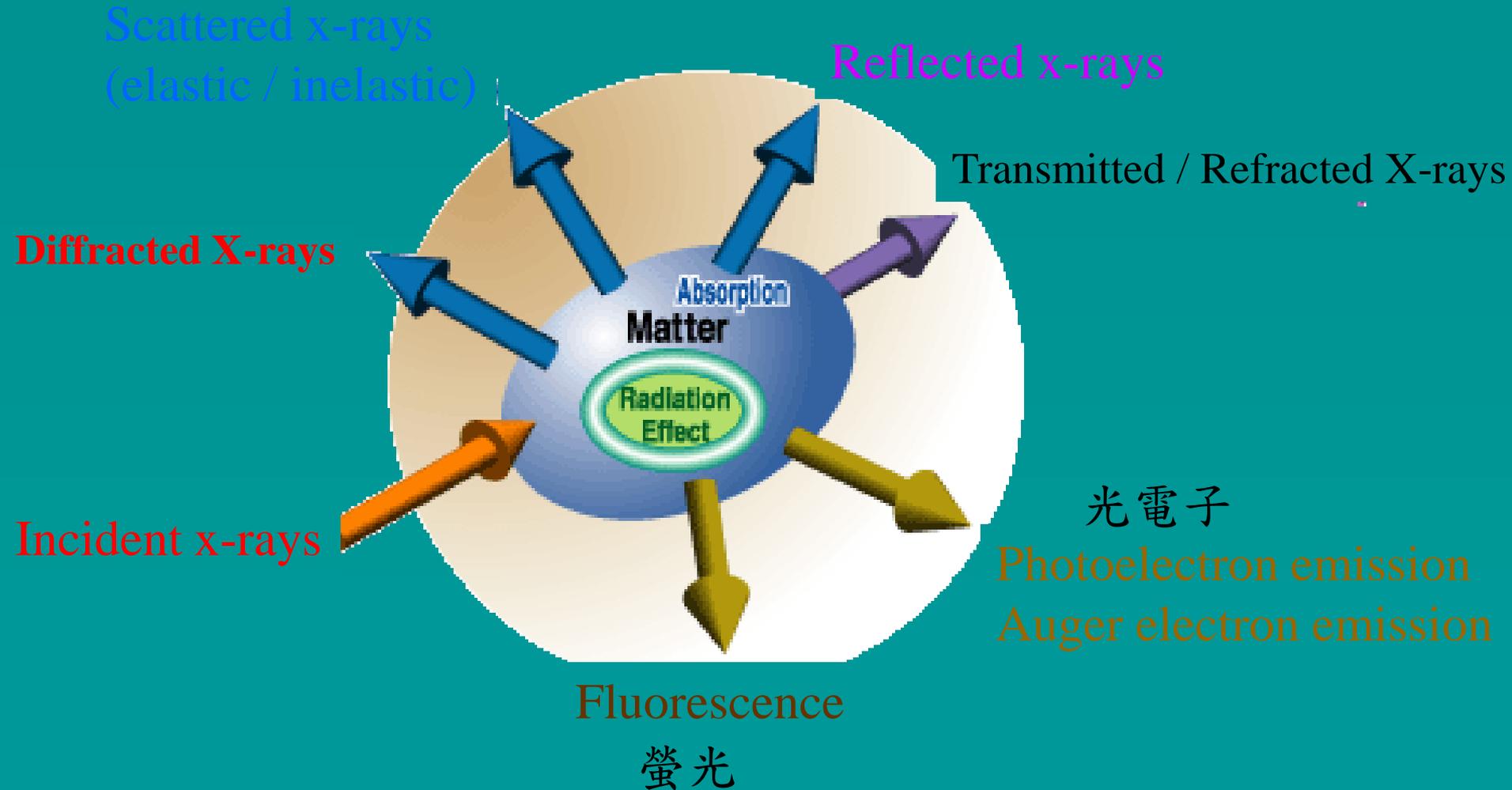


8 GeV;  $\gamma = 15,700$ ; 1.44 km circumference

**X光繞射**

**X-ray Diffraction**

# Interaction between x-rays and matter



# Discovery of X-ray Diffraction



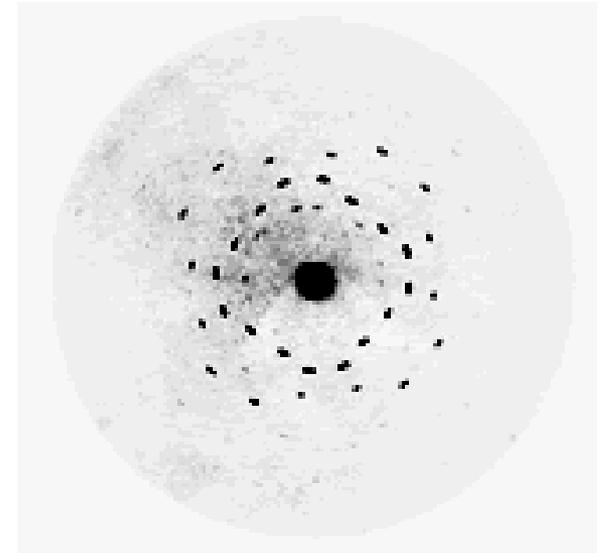
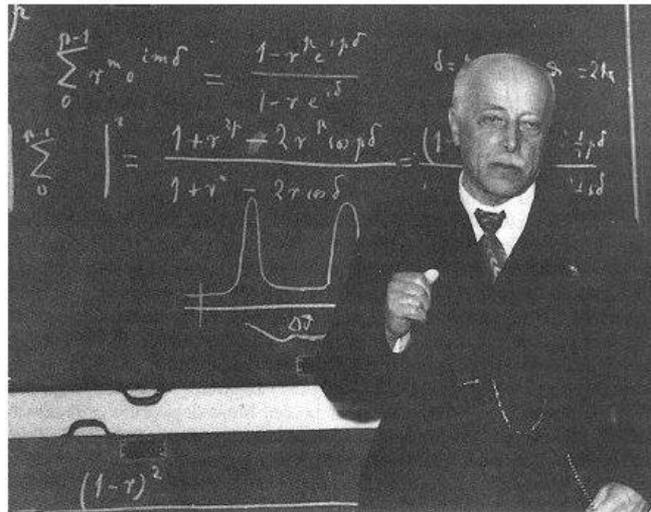
The Nobel Prize in physics 1914

“for his discovery of the diffraction of x-rays by crystals”

Max Von Laue



(1879-1960)  
Germany



Interference pattern observed by von Laue and collaborators using a photographic plate in 1912. The large central spot is due to the unscattered X-ray beam. The dark spots correspond to directions where x-rays scattered from crystal (ZnS) layers interfere constructively.

# Analysis of Crystal Structure

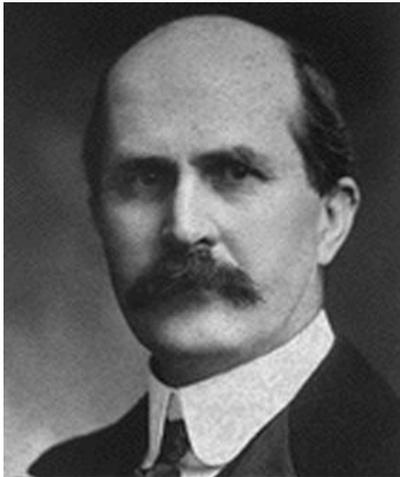


The Nobel Prize in physics 1915

“for their service in analysis of crystal structure by means of x-rays”

Sir William Henry Bragg

Sir William Lawrence Bragg



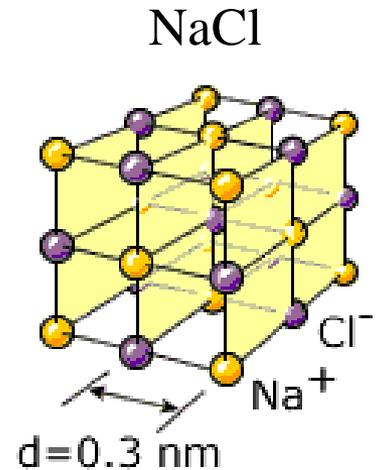
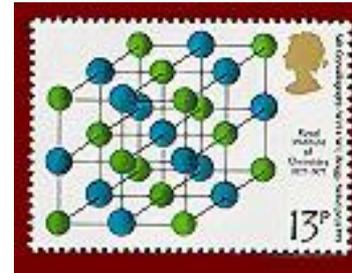
(1862-1942)

United Kingdom



(1890-1971)

United Kingdom



Determining the crystal structure NaCl, ZnS, Diamond,...

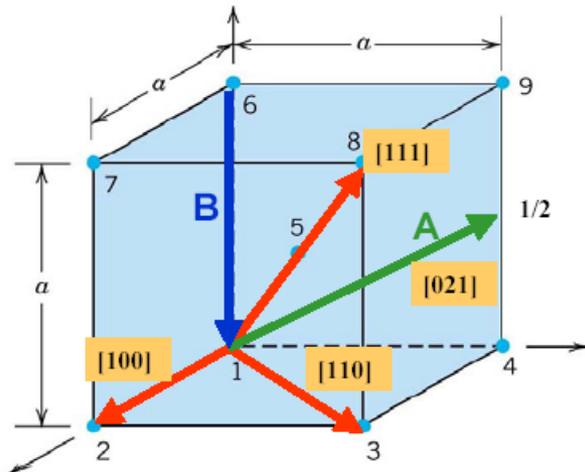
<http://www.nobel.se/physics/laureates/1915/>

<http://www.nobel.se/physics/educational/x-rays/what-6.html>

# Crystal Systems- Body-Centered Cubic (BCC)

## Indexing crystallographic directions

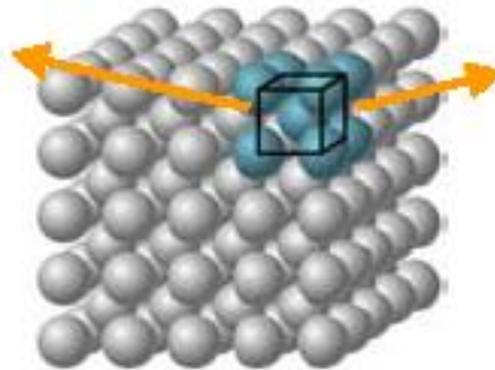
### Miller indices



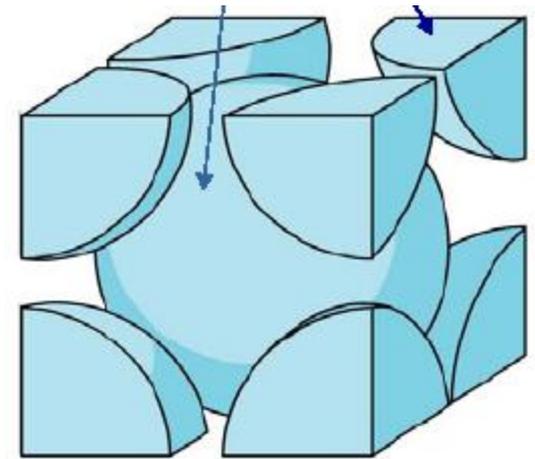
direction A =  $[021]$

direction B =  $[00\bar{1}]$

## Crystal



## Unit cell



$(111)$  plane

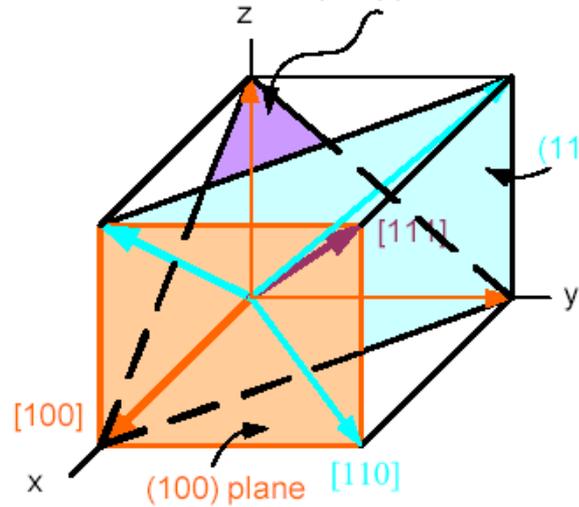
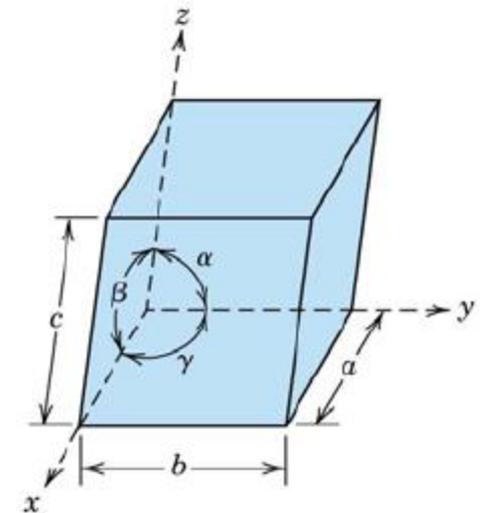
$(110)$  plane

$[111]$

$[110]$

## Lattice constant

$a, b, c, \alpha, \beta, \gamma$

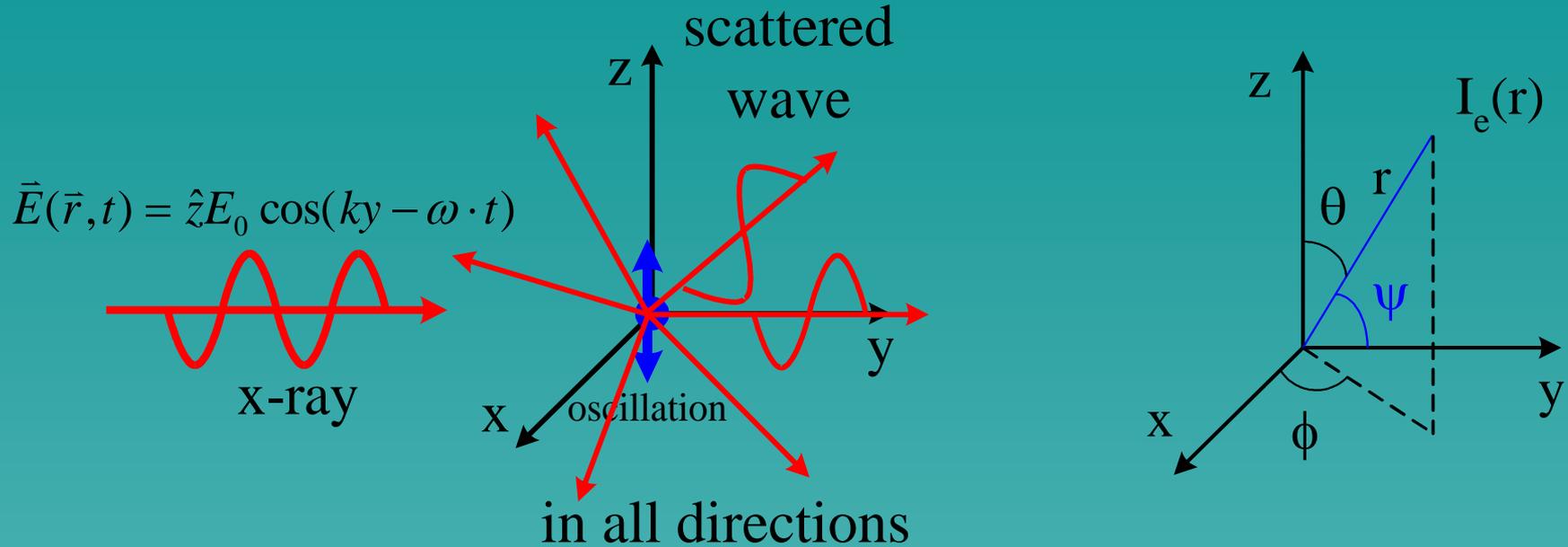


# 3D: 14 Bravais Lattices, 7 Crystal Systems

*Bravais Lattice*: an infinite array of discrete points with an arrangement and orientation that appears exactly the same from whichever of the points the array is viewed.

Name	Number of Bravais lattices	Conditions
Triclinic	1 (P)	$a_1 \neq a_2 \neq a_3$ $\alpha \neq \beta \neq \gamma$
Monoclinic	2 (P, C)	$a_1 \neq a_2 \neq a_3$ $\alpha = \beta = 90^\circ \neq \gamma$
Orthorhombic	4 (P, F, I, A)	$a_1 \neq a_2 \neq a_3$ $\alpha = \beta = \gamma = 90^\circ$
Tetragonal	2 (P, I)	$a_1 = a_2 \neq a_3$ $\alpha = \beta = \gamma = 90^\circ$
Cubic	3 (P, F, I)	$a_1 = a_2 = a_3$ $\alpha = \beta = \gamma = 90^\circ$
Trigonal	1 (P)	$a_1 = a_2 = a_3$ $\alpha = \beta = \gamma < 120^\circ \neq 90^\circ$
Hexagonal	1 (P)	$a_1 = a_2 \neq a_3$ $\alpha = \beta = 90^\circ$ $\gamma = 120^\circ$

# Scattering from one electron



J. J. Thomson (1906) analyzed scattering in detail. Scattered intensity at a distance,  $r$ , from a single electron for an **unpolarized** source is

$$I_e(r) = I_0 \frac{K}{r^2} \left( \frac{1 + \cos^2 \psi}{2} \right)$$

$$K = 7.94 \times 10^{-30} m^2$$

# Thomson scattering

## Electric field of incident beam

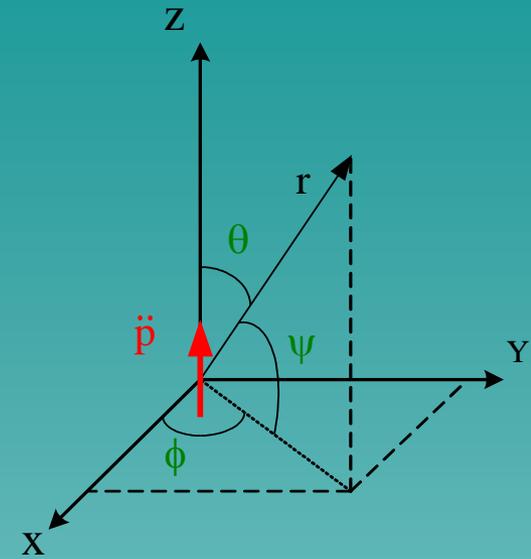
$$\vec{E}_i = \hat{e}_0 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

## Radiation field scattered by the electron

$$\begin{aligned} \vec{E}(\vec{r}, t) &= \frac{1}{4\pi\epsilon_0} \frac{e^2}{c^2 r^3} [\vec{r} \times (\vec{r} \times \ddot{\vec{r}})] \\ &= -\frac{1}{4\pi\epsilon_0} \frac{e^2 E_0}{mrc^2} [\hat{n} \times (\hat{n} \times \hat{e}_0)] e^{i(\vec{k} \cdot \vec{r} - \omega t)} \\ &= \vec{E}_s(\vec{r}, t) e^{i(\vec{k} \cdot \vec{r} - \omega t)}, \quad \hat{n} = \frac{\vec{r}}{r} \end{aligned}$$

$$\vec{B}(\vec{r}, t) = \frac{1}{c} \hat{n} \times \vec{E}_s(\vec{r}, t) e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\vec{E}_s(\vec{r}, t) = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{mc^2} \frac{E_0}{r} [\hat{n} \times (\hat{n} \times \hat{e}_0)]$$



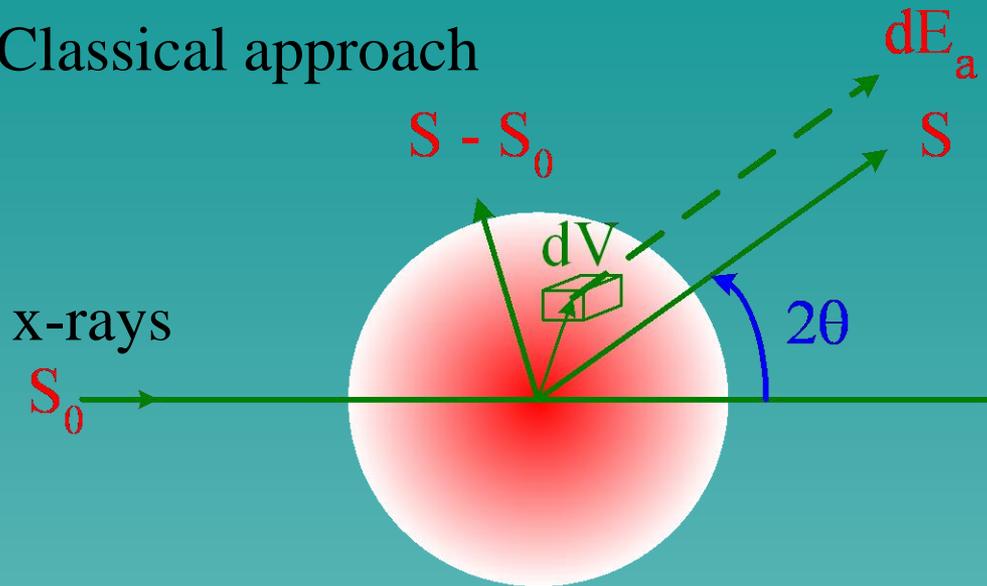
$$m\ddot{\vec{r}} = -e\vec{E}_i$$

$$\vec{r} = \hat{e}_0 \frac{eE_0}{m\omega^2} e^{-i\omega t}$$

$$\begin{aligned} r_c &= \frac{1}{4\pi\epsilon_0} \frac{e^2}{mc^2} \\ &= 2.82 \times 10^{-15} \text{ m} \end{aligned}$$

# Scattering from one atom

Classical approach



$\rho(r)$ : charge density of an atom

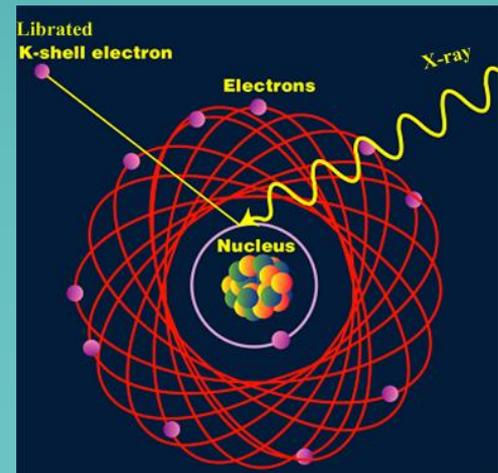
Assuming  $\rho(r)$  is spherically symmetric

$E_a(s)$ : Scattering amplitude from one atom

$E_e(s)$ : Scattering amplitude from one electron

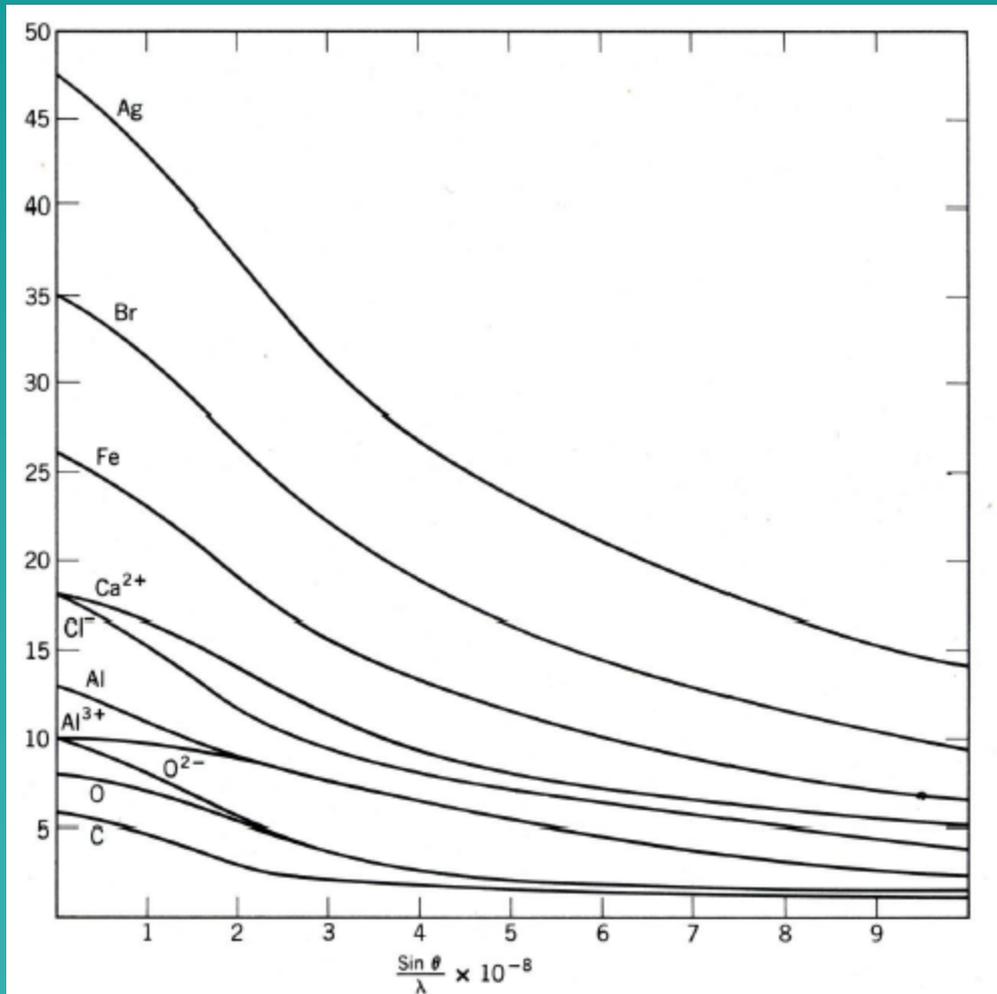
$$\begin{aligned}
 E_a(S) &= E_e(S) \int_{atom} \rho(r) e^{i\vec{S} \cdot \vec{r}} dV \\
 &= E_e(S) \int_{r=0}^{\infty} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \rho(r) r^2 e^{i2\pi r s \cos \theta} \sin \theta dr d\theta d\phi \\
 &= 4\pi \cdot E_e(S) \int_{r=0}^{\infty} \rho(r) r^2 \frac{\sin 2\pi r s}{2\pi r s} dr
 \end{aligned}$$

Atomic scattering factor:  $f \equiv \frac{E_a(s)}{E_e(s)}$



# Example: Atomic scattering factor

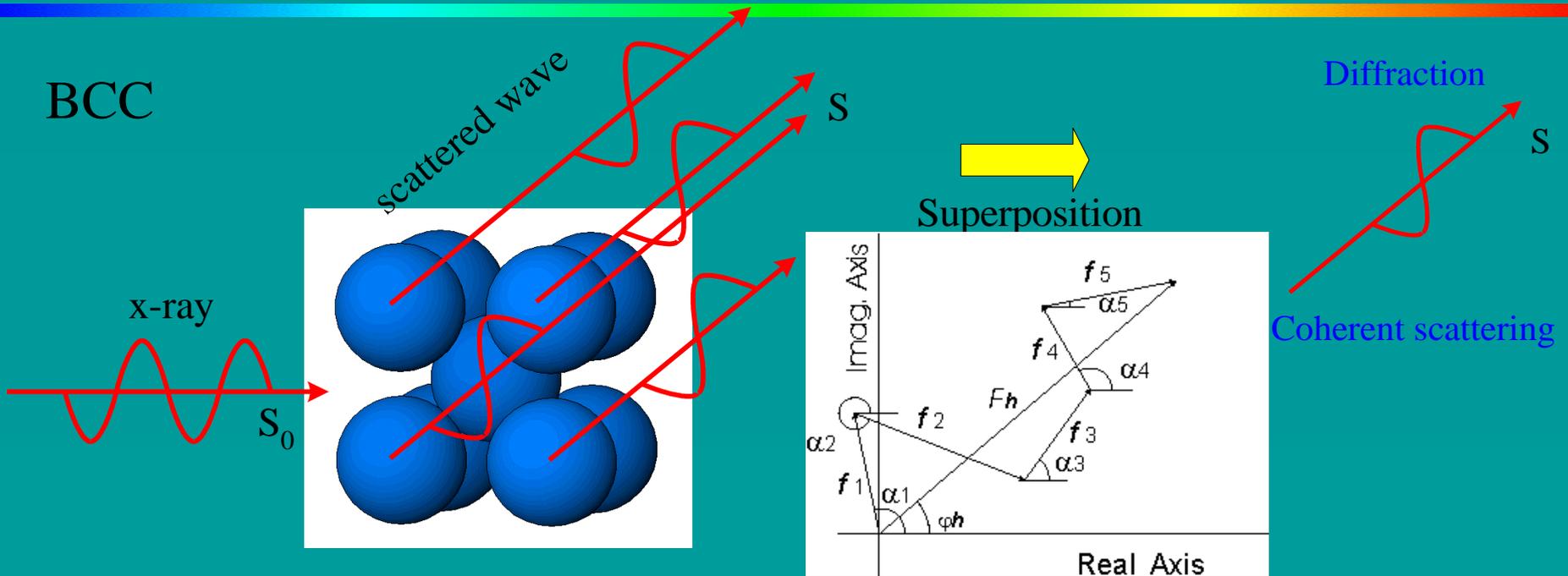
$$f = \frac{\text{amplitude of the wave scattered by an atom}}{\text{amplitude of the wave scattered by one electron}}$$



■ “ $f$ ” is expressed as a function of  $\sin\theta/\lambda$  as the interference depends on both  $\lambda$  and the scattering angle.

■ “ $f$ ” is equivalent to the **atomic number** at low angles, but it drops rapidly at higher  $\sin\theta/\lambda$ .

# Scattering from one unit cell



$$E_{cell}(s) = E_e(s) \cdot \int_{cell} \rho(\vec{r}) e^{2\pi i \vec{s} \cdot \vec{r}} d\vec{r}, \quad \rho(\vec{r}) = \sum_j \rho(\vec{r} - \vec{r}_j)$$

$$= E_e(s) \cdot \sum_j \left[ \int_{cell} \rho_j(\vec{r} - \vec{r}_j) e^{2\pi i \vec{s} \cdot (\vec{r} - \vec{r}_j)} d\vec{r} \right] \cdot e^{2\pi i \vec{s} \cdot \vec{r}_j}$$

$$= E_e(s) \cdot \sum_j f_j e^{2\pi i (hx_j + ky_j + lz_j)}$$

# Structure factor 結構因子

$|F_{hkl}| = \frac{\text{amplitude of the wave scattered by all the atoms in one unit cell}}{\text{amplitude of the wave scattered by one electron}}$

$$F_{hkl} = \frac{E_{cell}(s)}{E_e(s)} = \sum_{cell} \rho(r) e^{2\pi i \vec{s} \cdot \vec{r}} d\vec{r}$$
$$= \sum_j f_j e^{2\pi i (hx_j + ky_j + lz_j)}$$

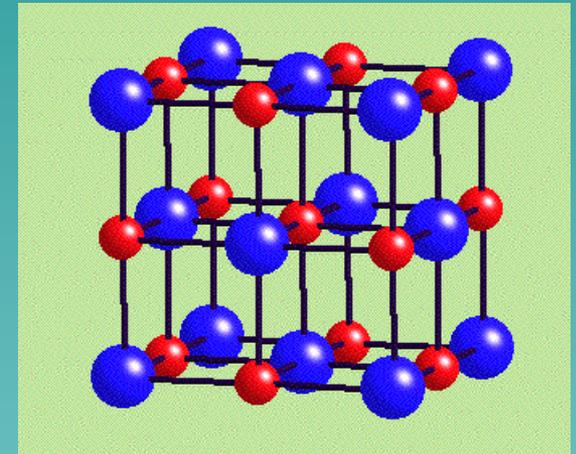
**Kinematical theory** 靜力

$$I = |F_{hkl}|^2 \cdot p \cdot \left( \frac{1 + \cos^2 2\theta}{\sin^2 \theta \cdot \cos \theta} \right) \cdot A(\theta) \cdot e^{-2M}$$

**Dynamical theory** 動力

$$I = \frac{8}{3\pi} \left( \frac{e^2}{mc^2} \right) \frac{N\lambda^2 |F_{hkl}|}{\pi \sin 2\theta} \left( \frac{1 + |\cos 2\theta|}{2} \right)$$

NaCl

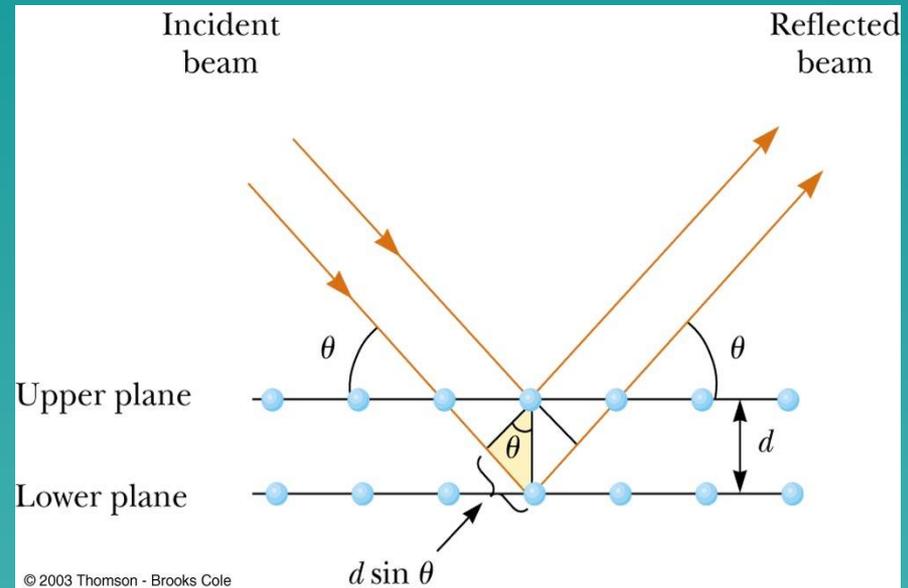


One unit cell contains 4(Na<sup>+</sup>Cl<sup>-</sup>)

# Bragg's Law 布拉格定律

- The beam reflected from the lower surface travels farther than the one reflected from the upper surface
- If the path difference equals some integral multiple of the wavelength, **constructive interference** occurs
- *Bragg's Law* gives the conditions for **constructive interference**

## Perfect Crystal = 3D Grating



$$2d \sin \theta = n\lambda$$

$$n=1, 2, 3, \dots$$

$d$ : Lattice spacing

$\lambda$ : Wavelength of x-ray

$\theta$ : Bragg angle

應用

Applications

# Applications of X-ray Diffraction

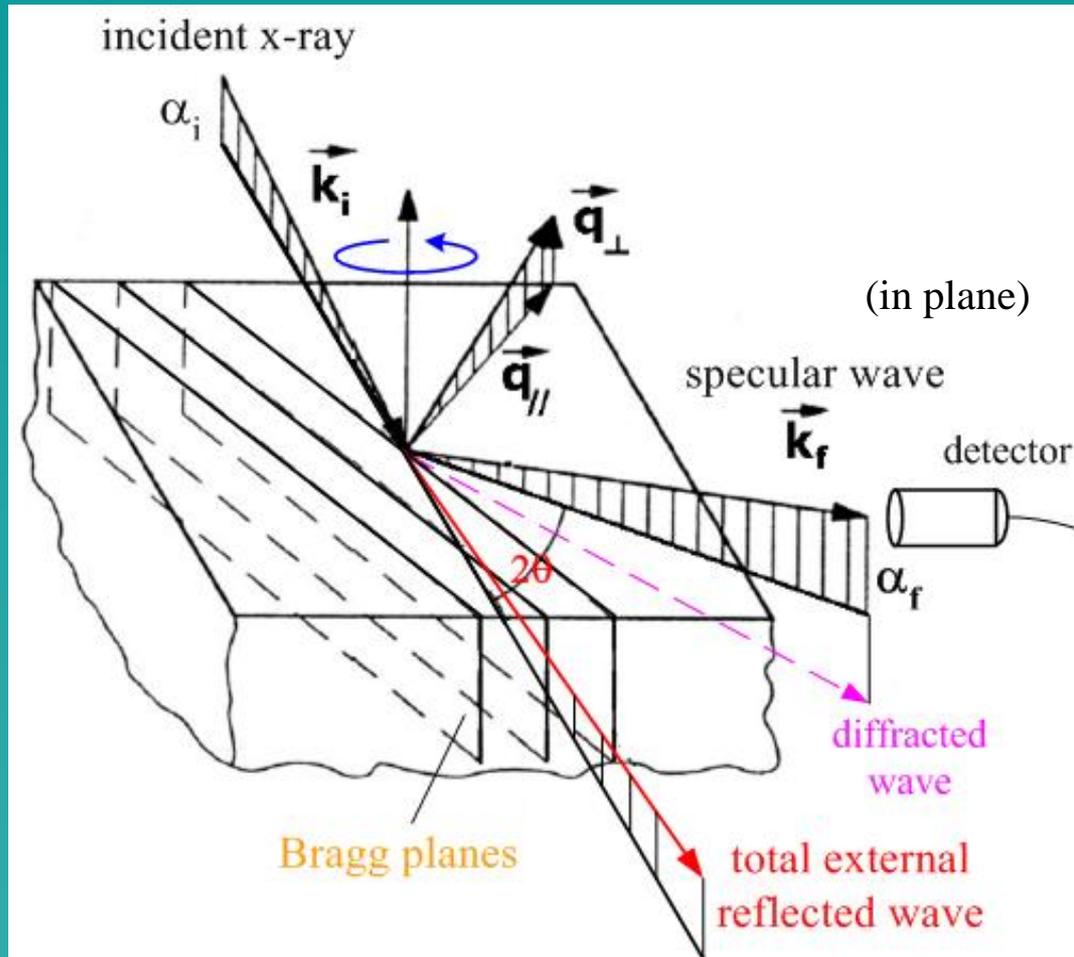
- **Biological and chemical 3D molecular structure determination (SCD)**
- **Qualitative and quantitative crystalline phase composition analysis (XRPD)**
- **Material structure analysis (XRD, HRXRD, XRR,  $\mu$ -XRD, SAXS)**
  - crystallite size / particle size in the nm range
  - **microstrain**
  - residual stress / fatigue stress
  - texture / preferred orientation
  - **thin films** and **multilayers**: thickness, layer sequence, density, surface and interface roughness
  - **Nanostructure** analysis (particle shape, size and orientation)

# Techniques used for X-ray Diffraction

---

- **X-ray powder diffraction (XRPD)**
- **High resolution X-ray diffraction (HRXRD)**
- **Single crystal diffraction (SCD)**
- **Small angle X-ray scattering (SAXS)**
  
- **X-ray reflectometry (XRR)**
- **Grazing incidence diffraction (GID)**
- **In-plane grazing incidence diffraction (IPGID)**
- **X-ray resonant diffraction**
- **X-ray multiple diffraction**

# Grazing (Glancing) Incidence X-Ray Diffraction (GIXD)



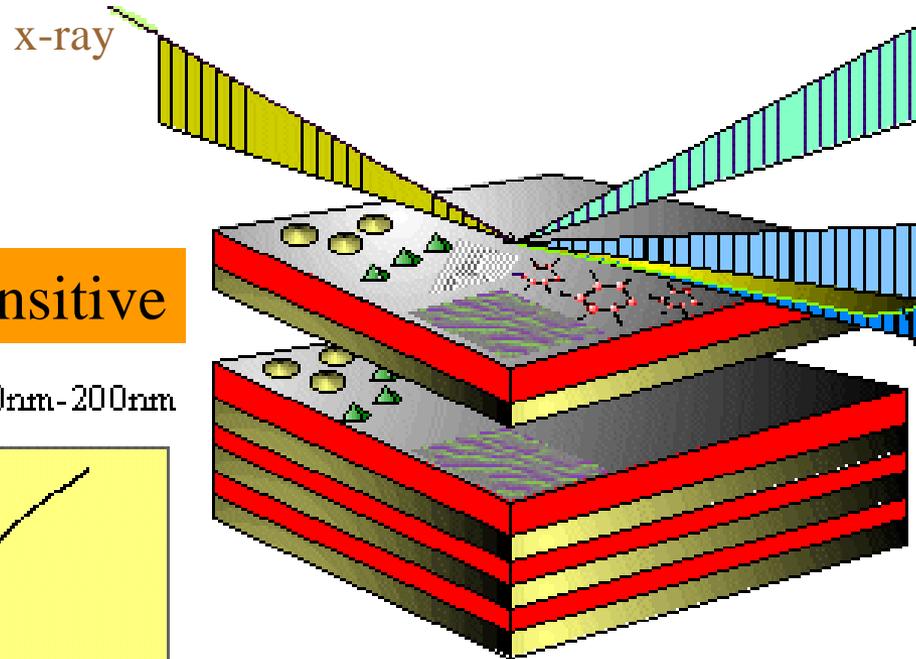
When the incident beam close to or below the **critical angle** for total external reflection, a Bragg reflection is excited from planes perpendicular to the surface.

**Greatly enhanced surface sensitivity**

# Information from Grazing Incidence X-rays

## X-ray methods at grazing incidence (GID, GISAXS,...)

### Structure of thin layers on substrate



diffracted beam (GID)

**Crystalline properties**  
Strain in thin layers  
**NANOSTRUCTURES**  
Phasetransitions

reflected beam

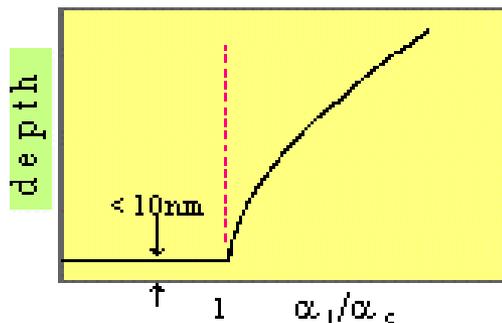
**density profile**  
thickness  
surface roughness  
interface properties  
buried layers

Diffuse scattering (GISAXS)

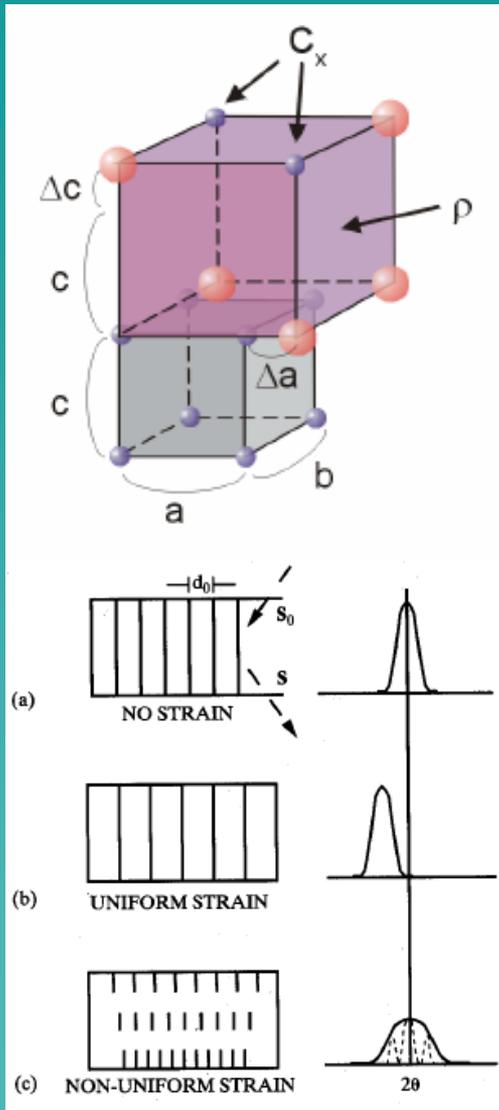
Morphology of nano-islands  
Interfaces roughness & growth  
Lateral surface structures  
Structure of macro-molecules

### Surface sensitive

depth resolution 10nm-200nm



# Strain



■  $\Delta a/a$

Mismatch (or strain) in lateral (in-plane, parallel) direction.

■  $\Delta c/c$

Mismatch (or strain) in normal (vertical, perpendicular) direction.

■  $R$

Lattice relaxation degree.

■  $\rho$

Mass density

■  $C_x$

Concentration for ternary solid solution

■  $C_y$

Concentration for quaternary solid solution

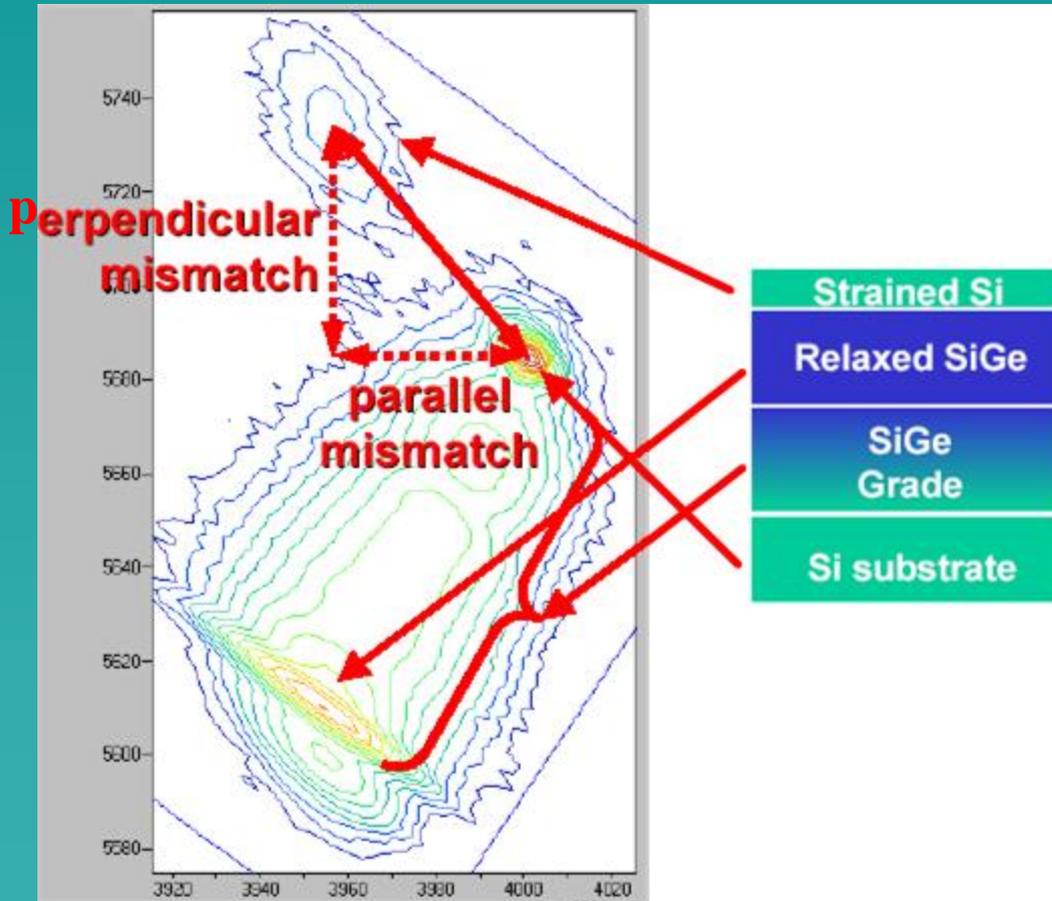
In a material there may be strains that vary from grain to grain or within a grain (microstrains) due to the local environment and there may also be a uniform strain due to an external load (macrostrain)

## Differential Bragg's law

$$\Delta 2\theta = -\frac{\Delta d}{d} \tan \theta$$

crystalline diffraction line

# How to measure strain



Horizontal offset gives parallel lattice constant

Vertical offset gives perpendicular lattice constant

Combining the lattice constant information gives the parallel and perpendicular strain

# Basic concepts of reflectivity

Specular x-ray reflectometry - characterize  $e^-$  density profile normal to the interface

Refractive index for x-ray:

$\mathbf{n} = \mathbf{1} - \delta + i\beta$ , less than 1, for X-ray!

$\delta = r_0 \rho_0 \lambda^2 / 2\pi$  ( $10^{-5}$  to  $10^{-6}$ )

and  $\beta$  - absorption,  $\beta = \mu \lambda / 4\pi$

$r_0$  - classical electron radius

$\rho_0$  - electron density

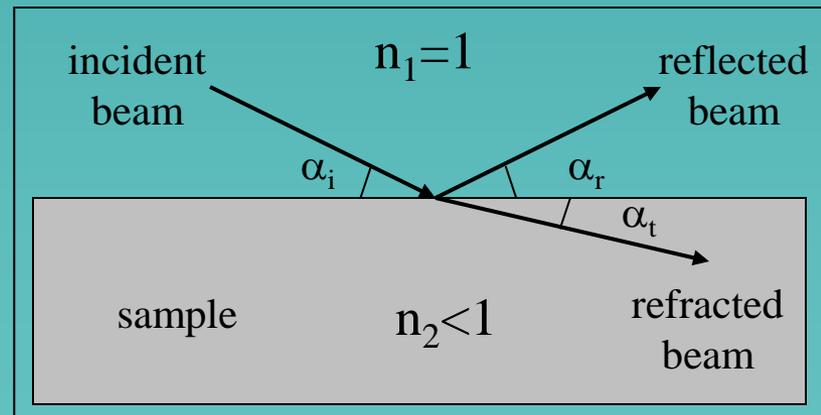
$\mu$  - mass absorption coefficient

Snell's law:

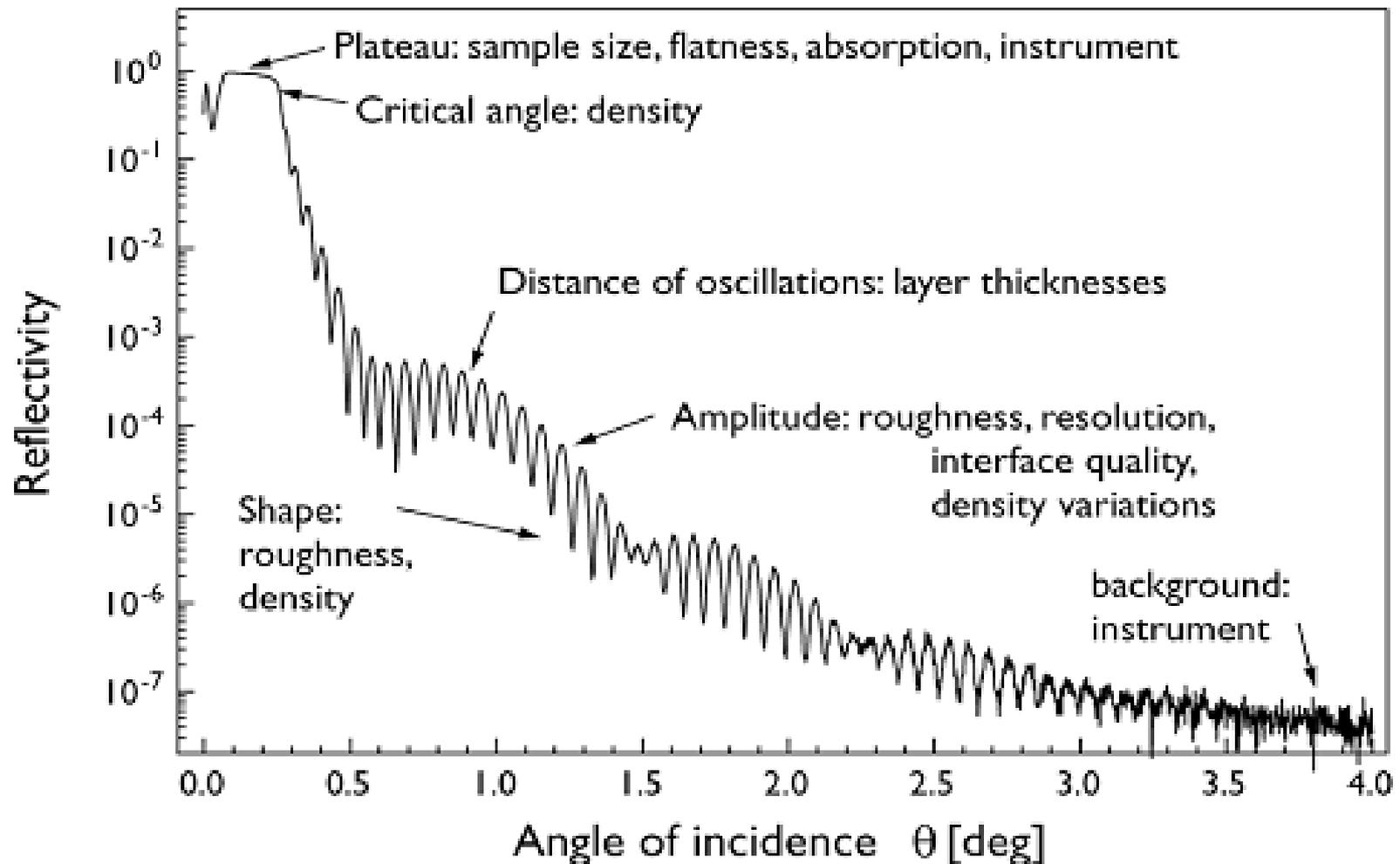
$$n_1 \cos \alpha_i = n_2 \cos \alpha_t$$

$$\alpha_c (\alpha_i \text{ for } \alpha_t = 0) = (2\delta)^{1/2}$$

$$\alpha_c \sim 0.1 - 0.6^\circ \text{ for } \lambda = 1.5 \text{ \AA}$$



# The Information Derived from a Reflectivity Curve



# X-ray crystallography X光結晶學

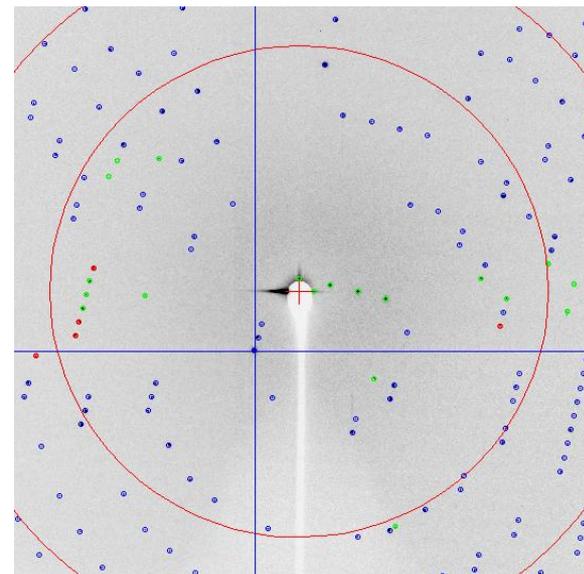
## 1. Crystallization



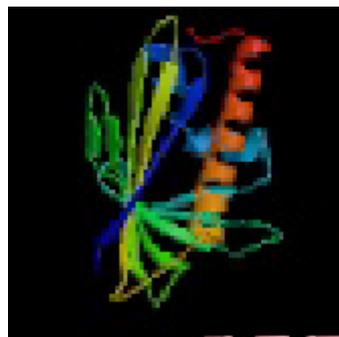
## 2. Data collection



## 3. Phase determination



## 6. Structure



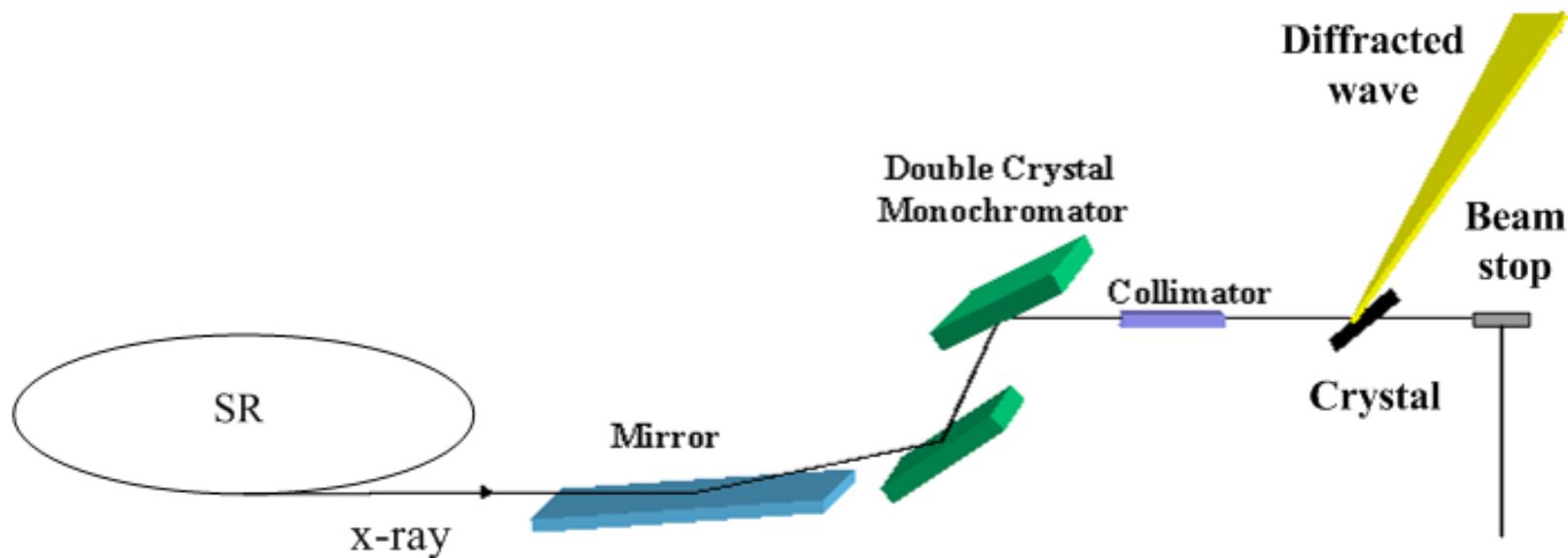
$$I \propto |F_{hkl}|^2$$

$$F_{hkl} = |F_{hkl}| e^{i\alpha_{hkl}}$$

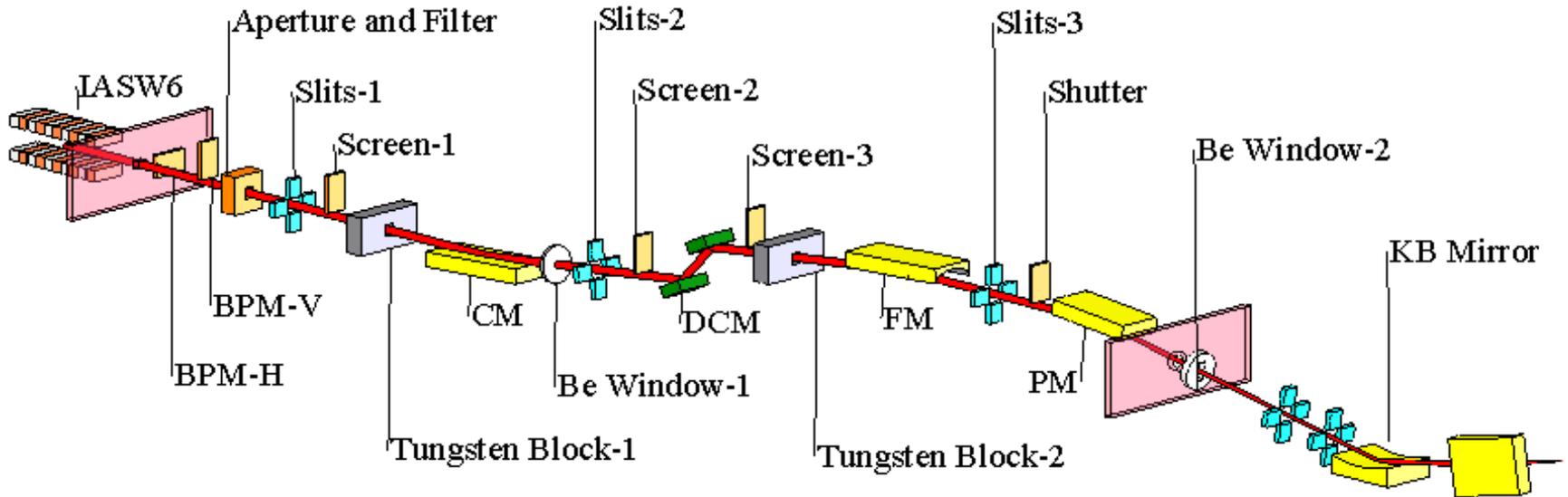
4. Calculation & interpretation of electron density map

5. Refinement of the molecular model

# Beamline Setup



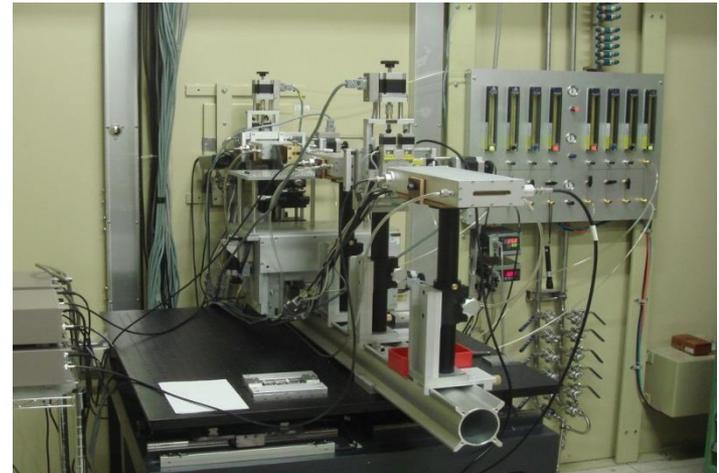
# X-ray Scattering BL07A (NTHU, Tamkang U., T3)



X-ray scattering/diffraction station

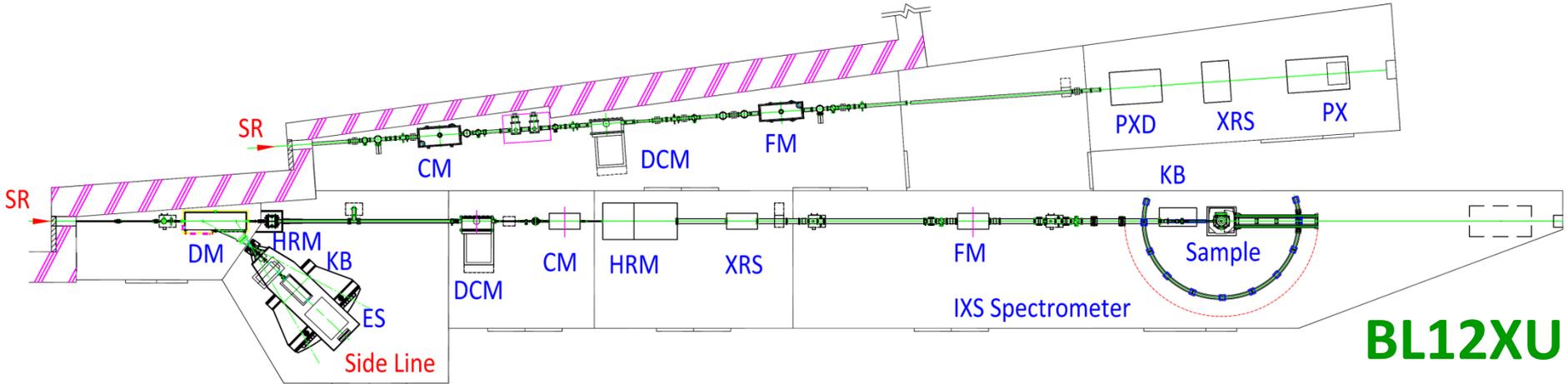


X-ray absorption station

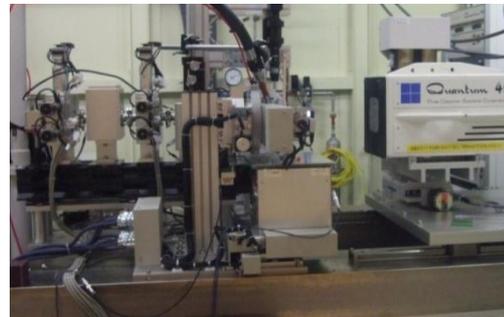
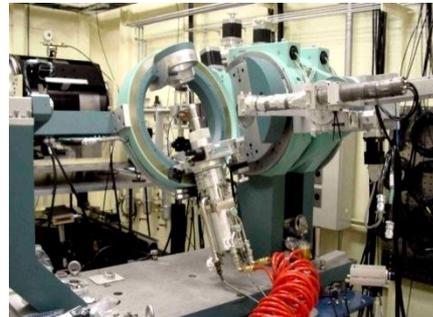


# SPring-8 Taiwan Beamlines

BL12B2



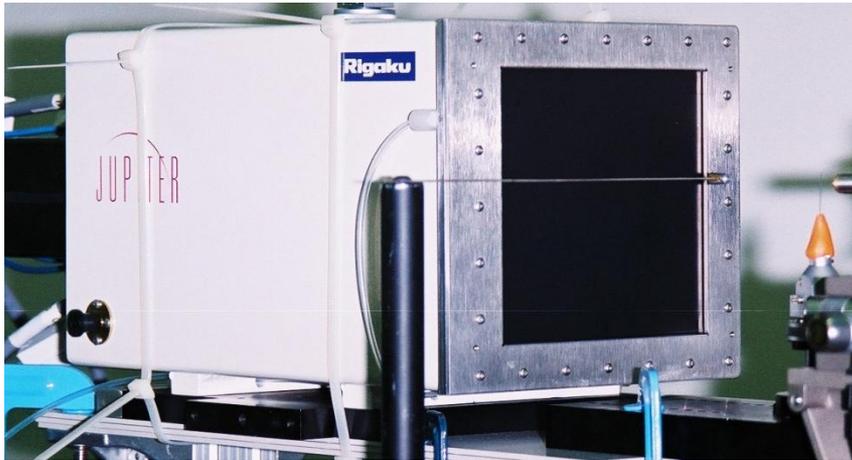
BL12XU



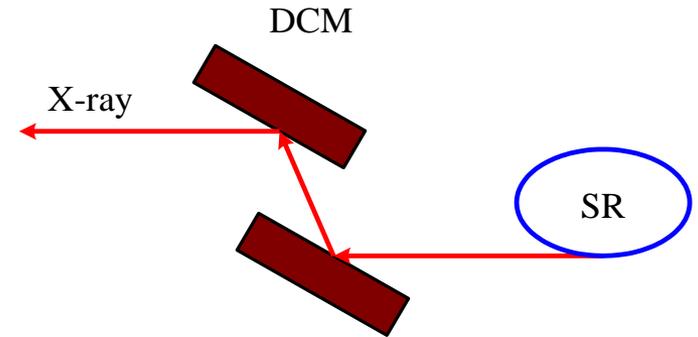
10Y review & 5Y contract renewal (2010/2/24)

# Data Collection

Charge Couple Device (CCD)



Lysozyme



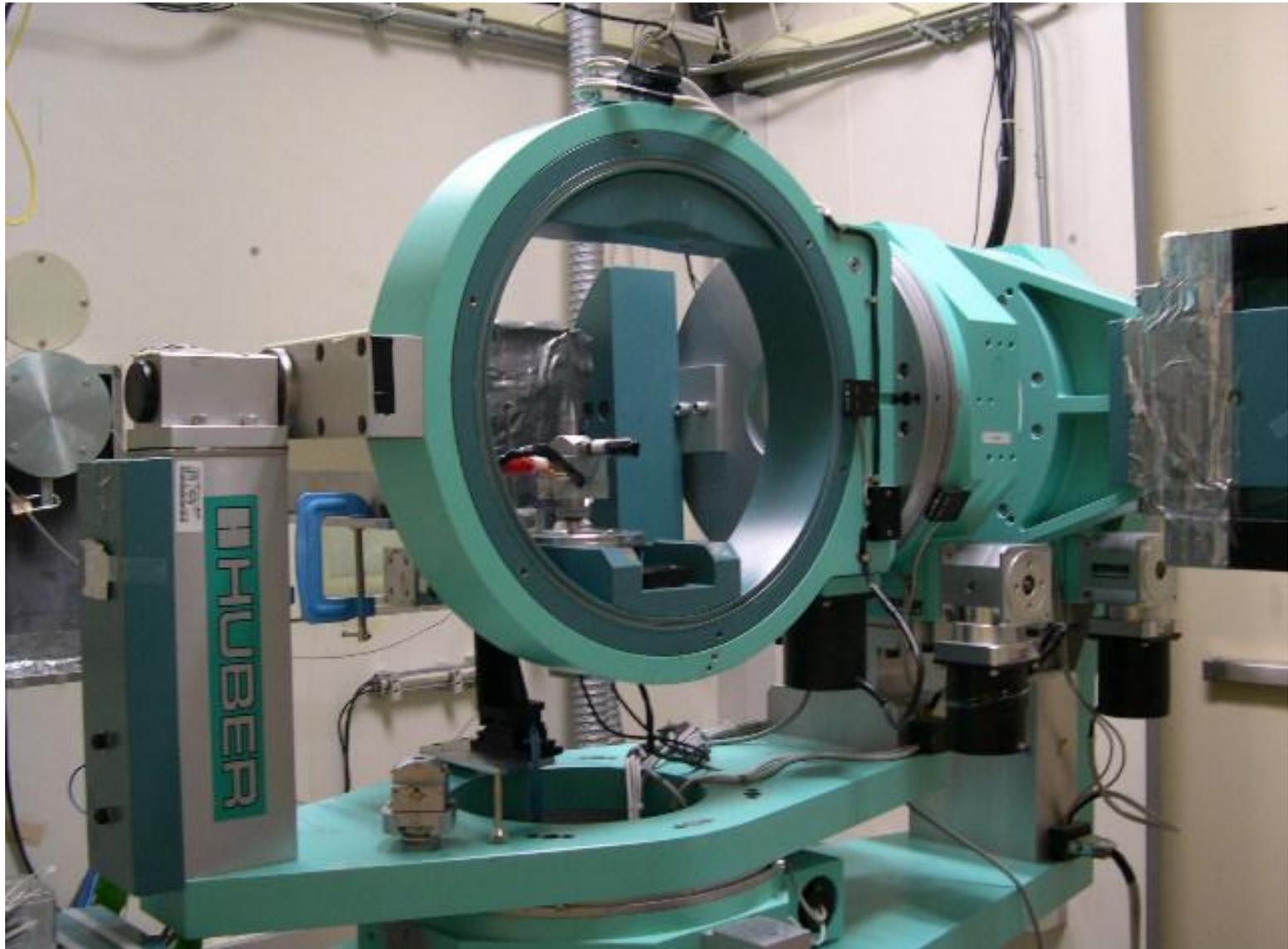
Control PC system



8C diffractometer



# 8C diffractometer 八環繞射儀器





# Diffraction Image of Lsozyme

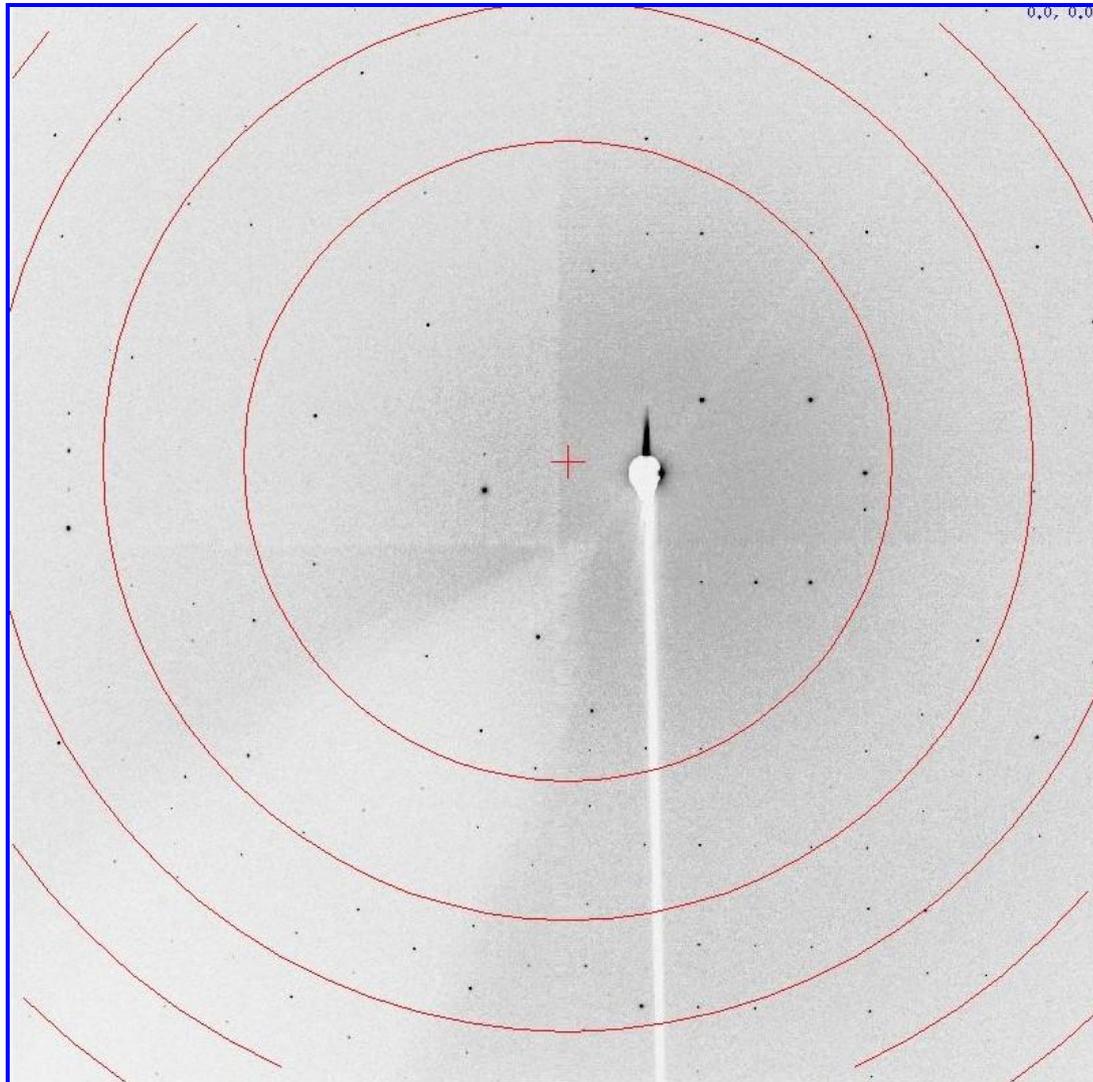




Photo: MRC Laboratory of  
Molecular Biology

**Venkatraman  
Ramakrishnan**

1/3 of the prize

United Kingdom

MRC Laboratory of  
Molecular Biology  
Cambridge, United  
Kingdom

b. 1952

(in Chidambaram,  
Tamil Nadu, India)



Credits: Michael  
Marsland/Yale University

**Thomas A.  
Steitz**

1/3 of the prize

USA

Yale University  
New Haven, CT, USA;  
Howard Hughes  
Medical Institute

b. 1940



Credits: Micheline  
Pelletier/Corbis

**Ada E. Yonath**

1/3 of the prize

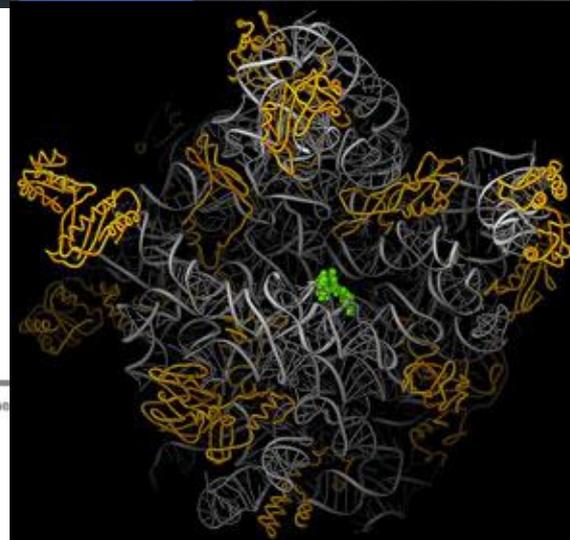
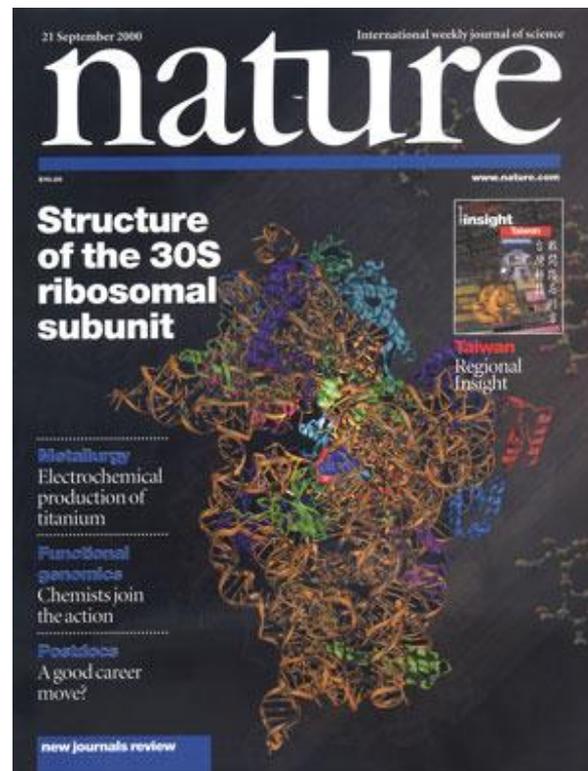
Israel

Weizmann Institute of  
Science  
Rehovot, Israel

b. 1939

Titles, data and places given above refer to the time of the award.

Source: Nobel Prize website



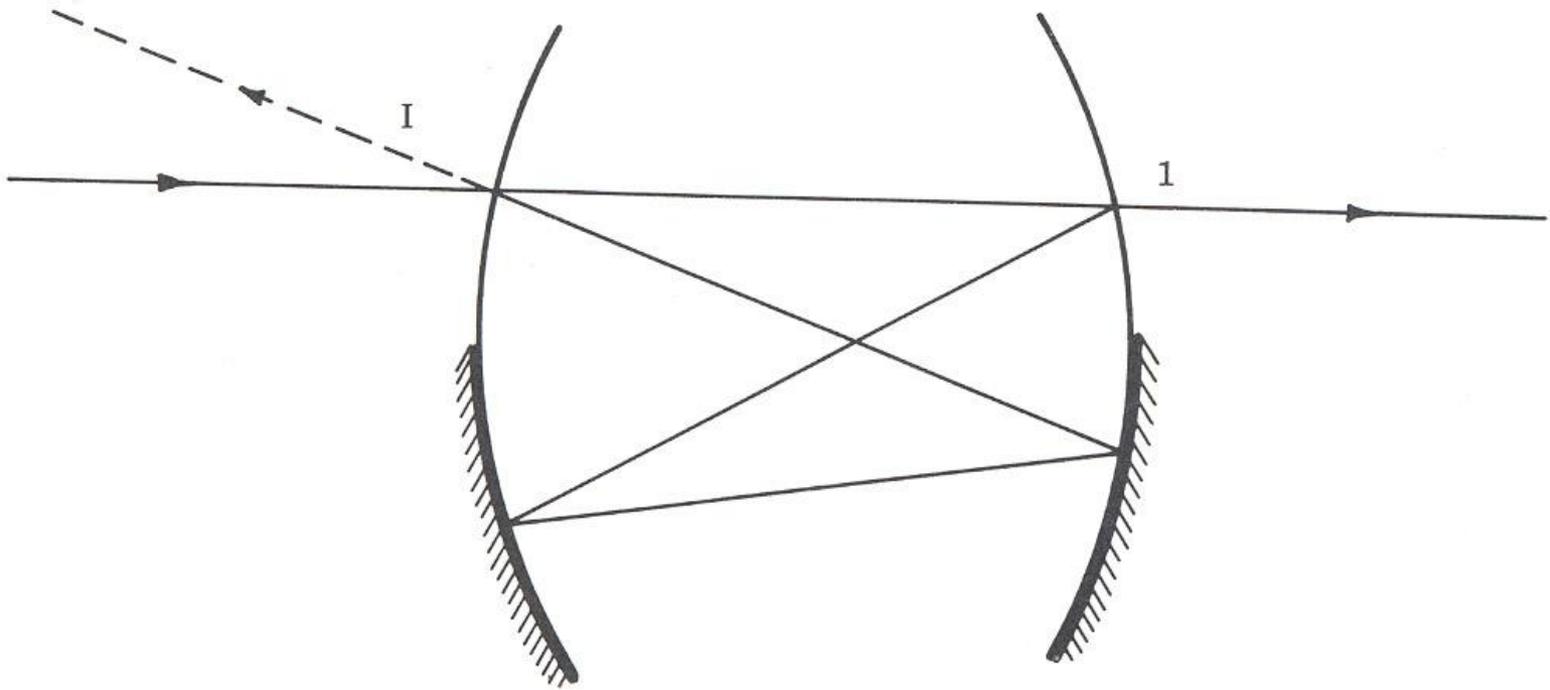
Source: NSLS website

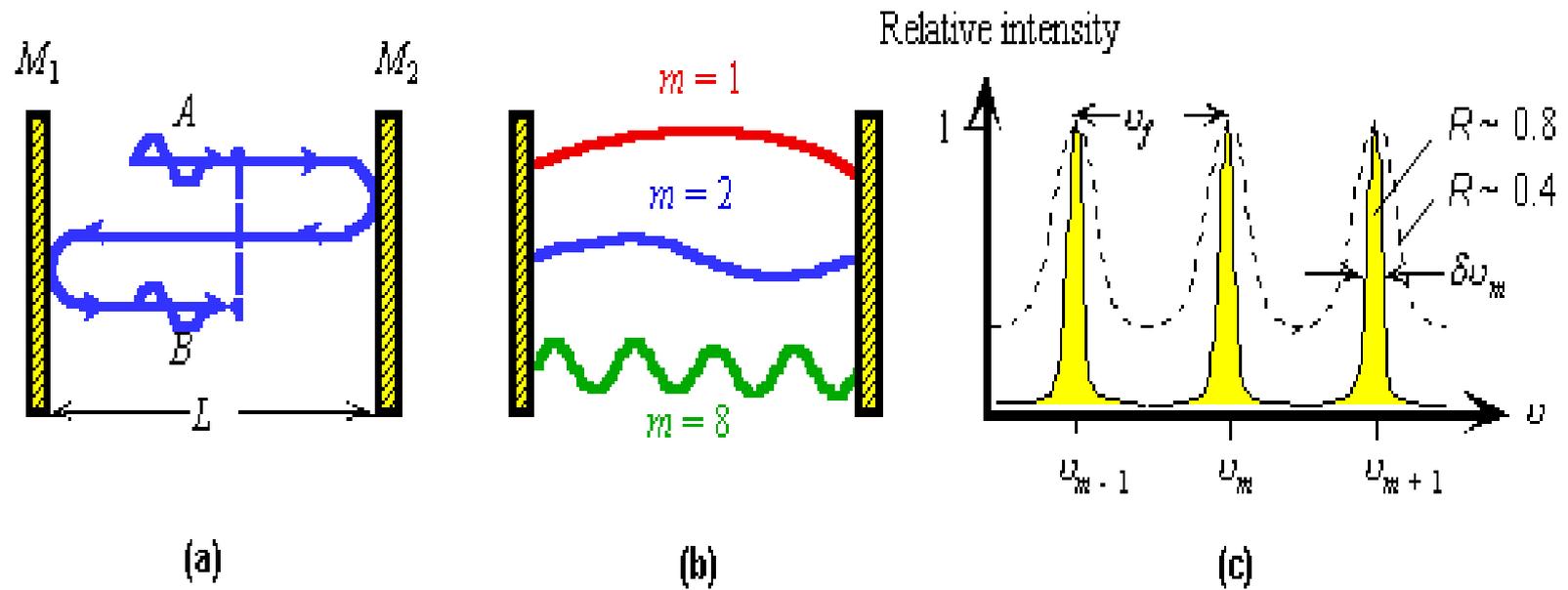
# **Realization of X-Ray Fabry-Perot Cavity: A Synchrotron Diffraction Experiment**

**X光共振腔--X光繞射實驗**

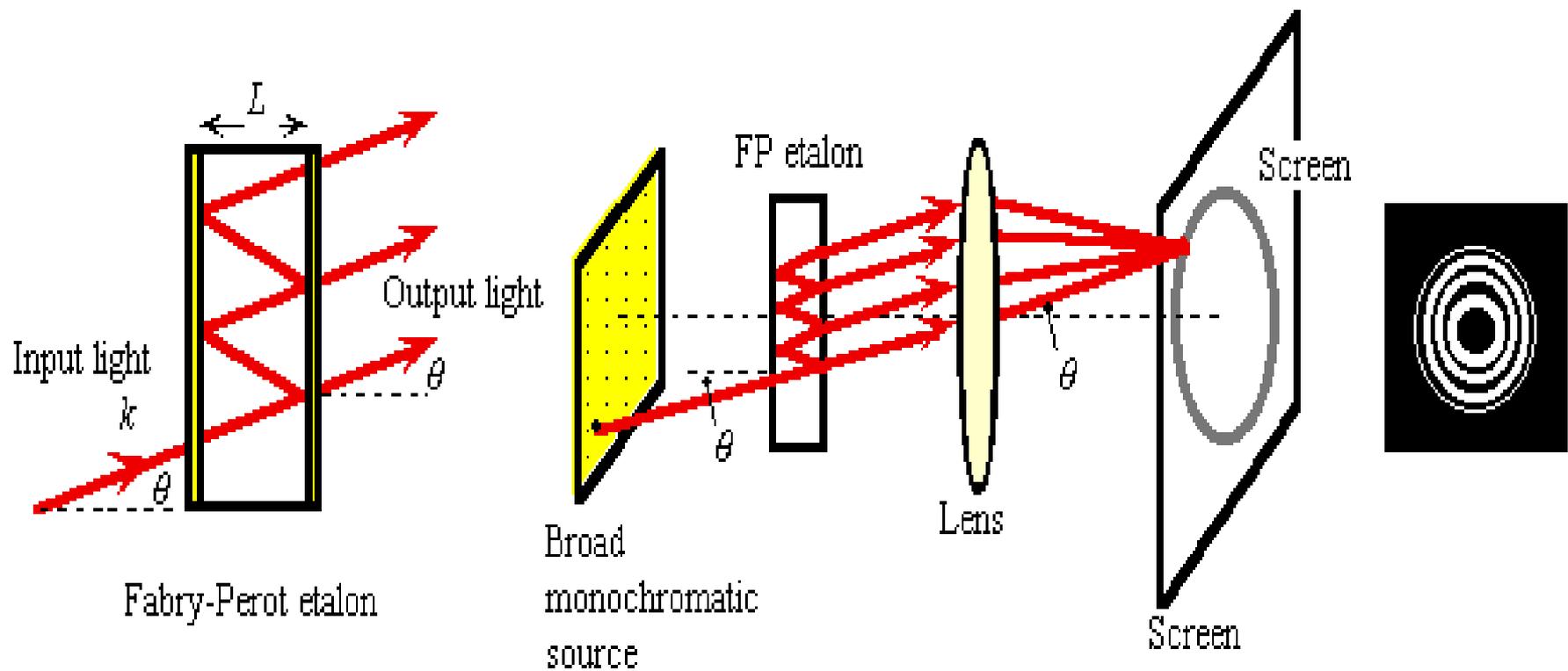
**NTHU, NSRRC, SPring8/Riken**

# Optical Fabry-Parot Interferometer



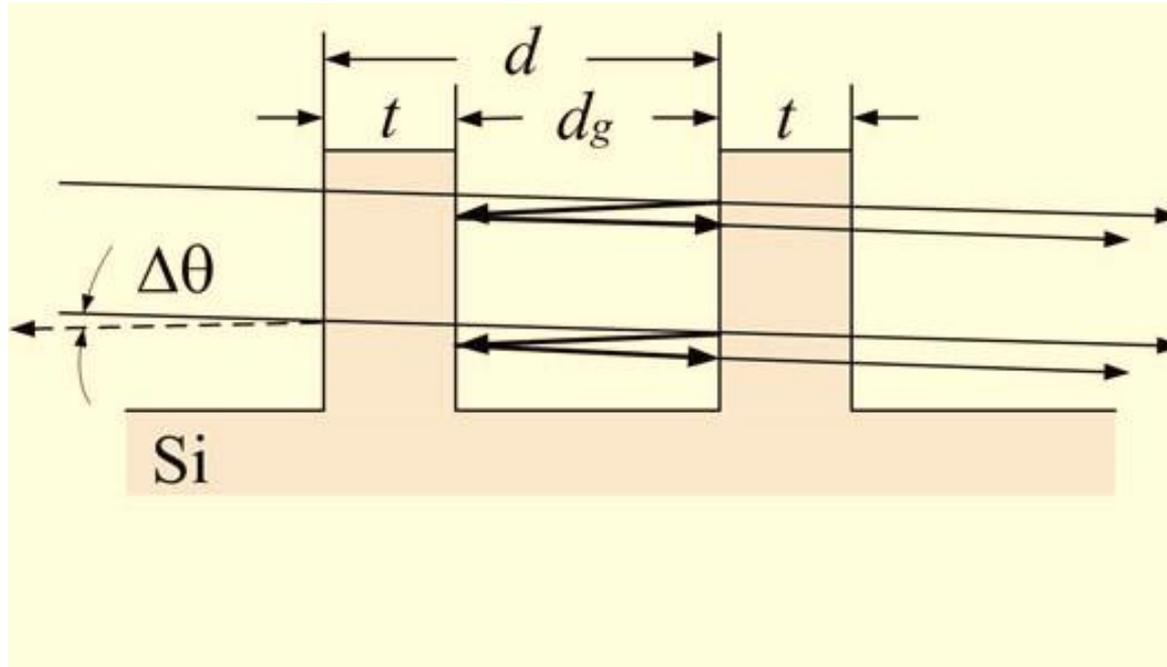


Schematic illustration of the Fabry-Perot optical cavity and its properties. (a) Reflected waves interfere. (b) Only standing EM waves, *modes*, of certain wavelengths are allowed in the cavity. (c) Intensity vs. frequency for various modes  $R$  is mirror reflectance and lower  $R$  means higher loss from the cavity.



Fabry-Perot optical resonator and the Fabry-Perot interferometer (schematic)

# X-ray Fabry-Perot resonator



背向反射(back reflection)

**Bragg angle=90 deg.**

# 微電子蝕刻術 (lithography)

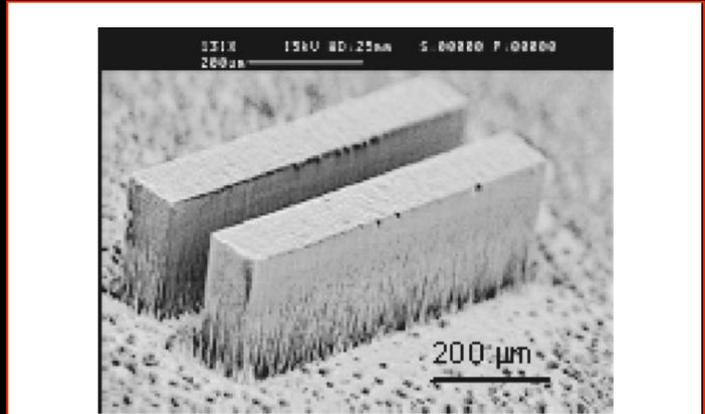
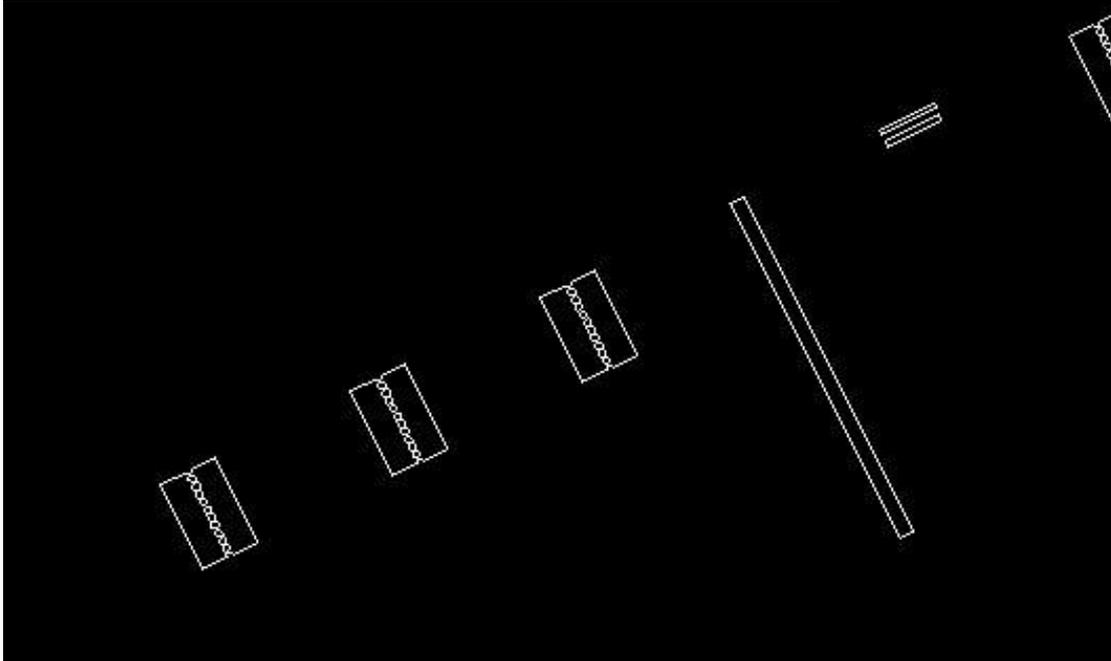
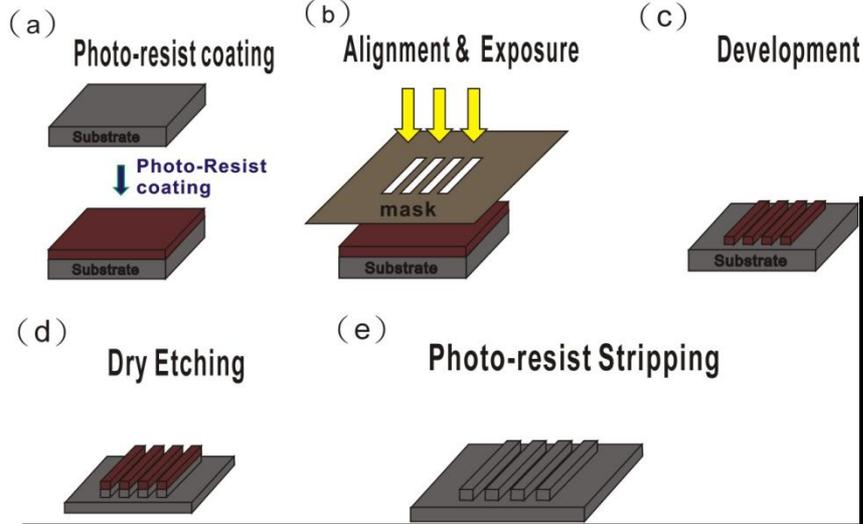
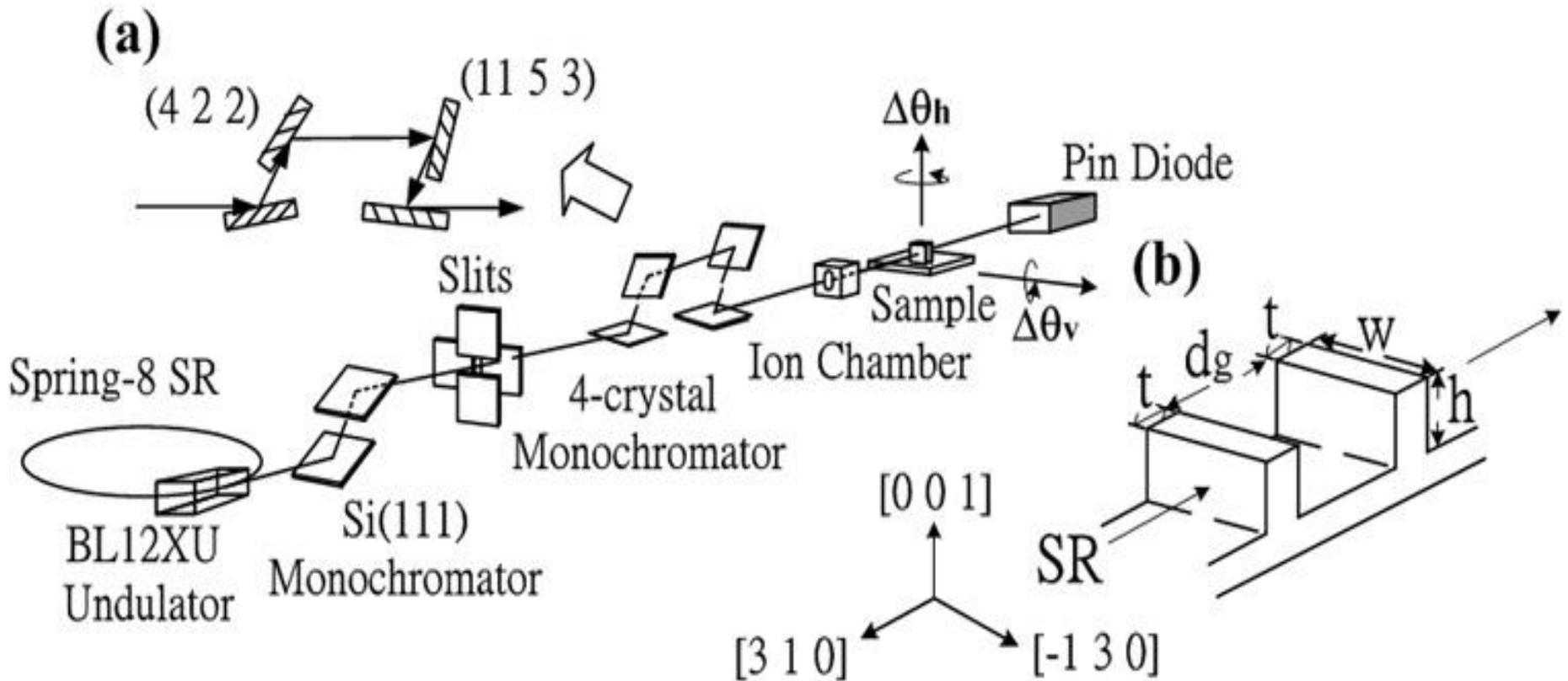


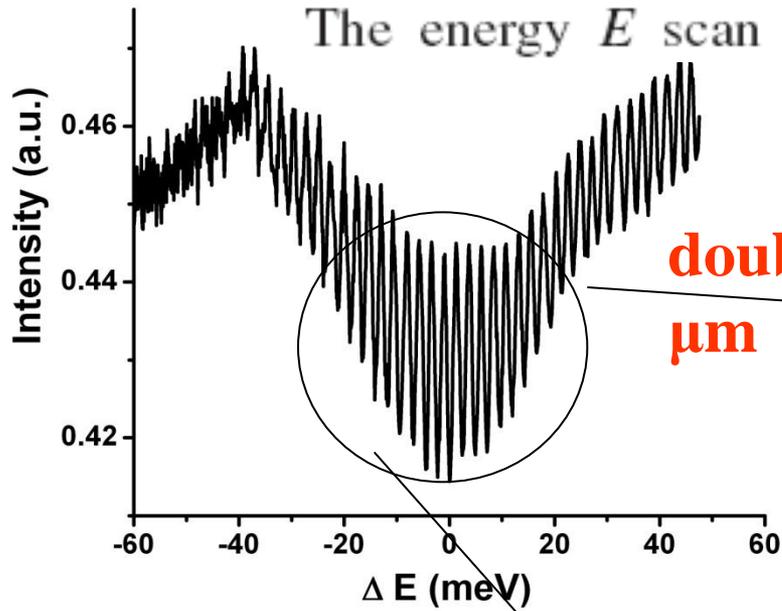
FIG. 1. A SEM photograph of a thick two-plate (100/100) x-ray Fabry-Pérot cavity of silicon. The width of the crystal plate is 800 μm, the height is 200 μm, and the thickness is 100 μm. The direction normal to the substrate plane is [001] and that perpendicular to the two crystal plates is [310].

# Experimental set-up

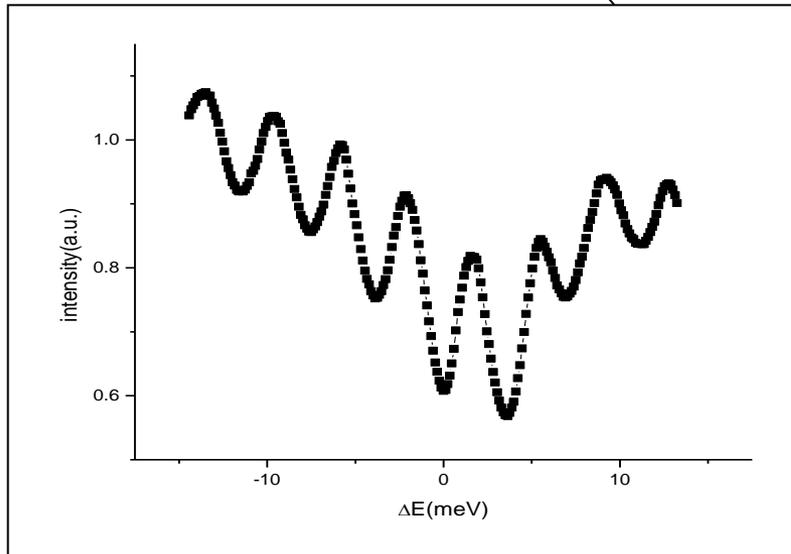
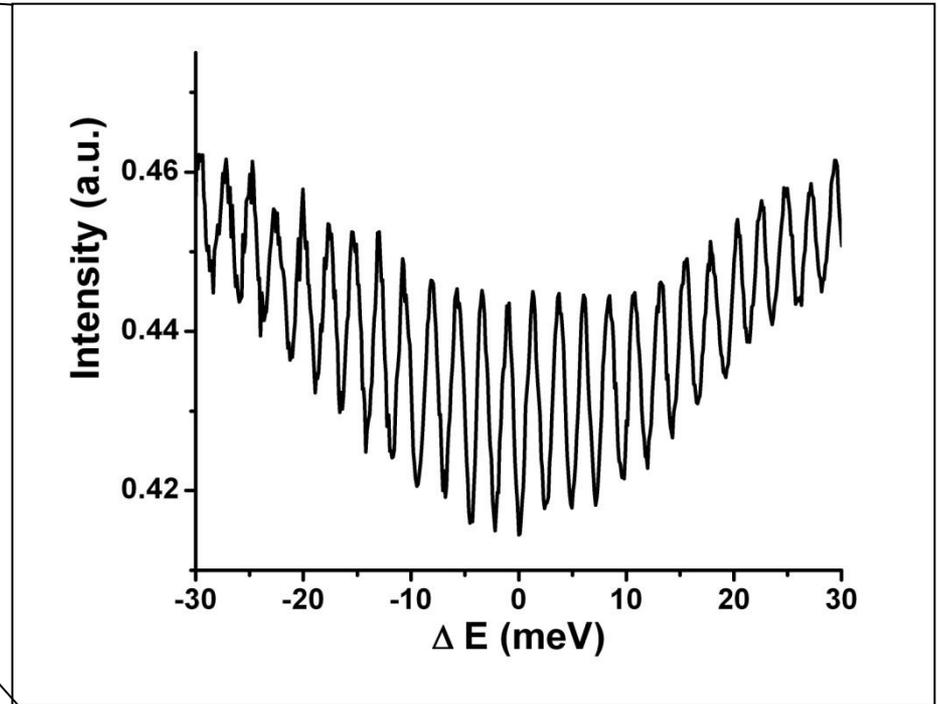


Si (12 4 0) back reflection  
for **E=14.4388 keV** (0.8588 Å)

The energy  $E$  scan of the forward-transmitted (000)



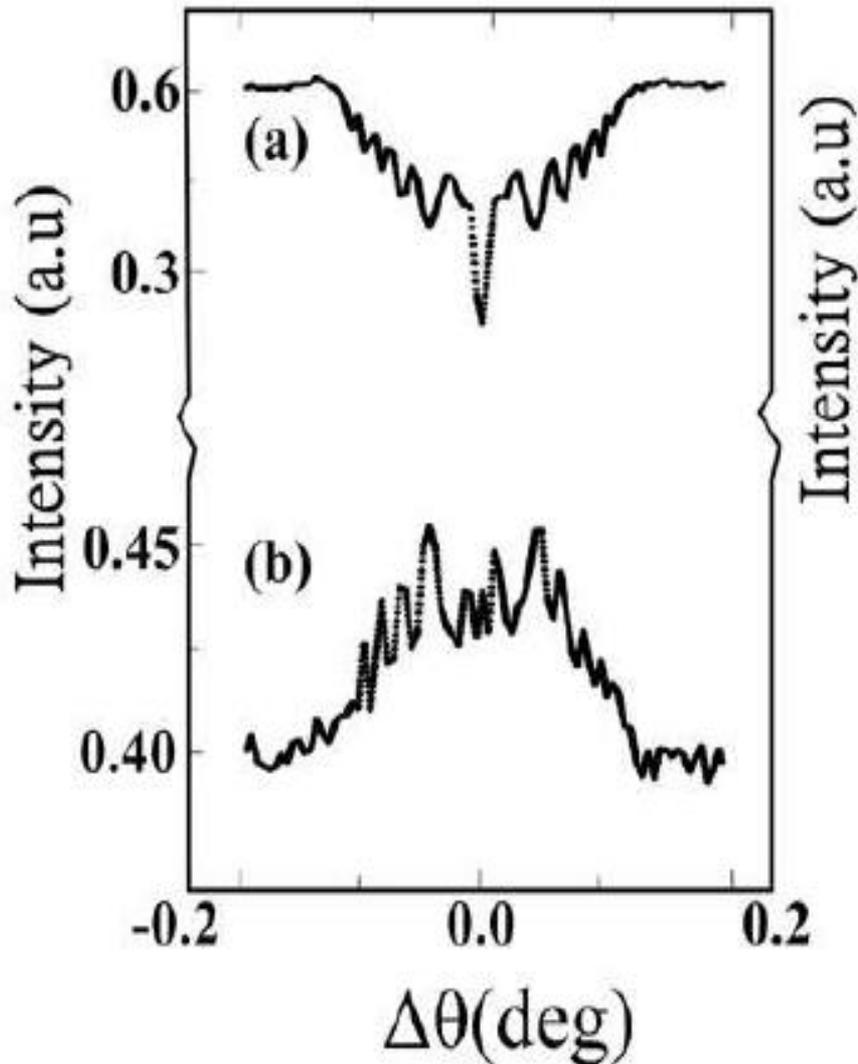
**double thin plates,  $t = 100\mu\text{m}$   $dg = 150\mu\text{m}$**



**double thinner plates,  $t = 70\mu\text{m}$   $dg = 100\mu\text{m}$**

$\Delta\theta$ - scans at  $\Delta E = 9$  meV

The  $\Delta\theta$ -scan at 0.002 deg./step

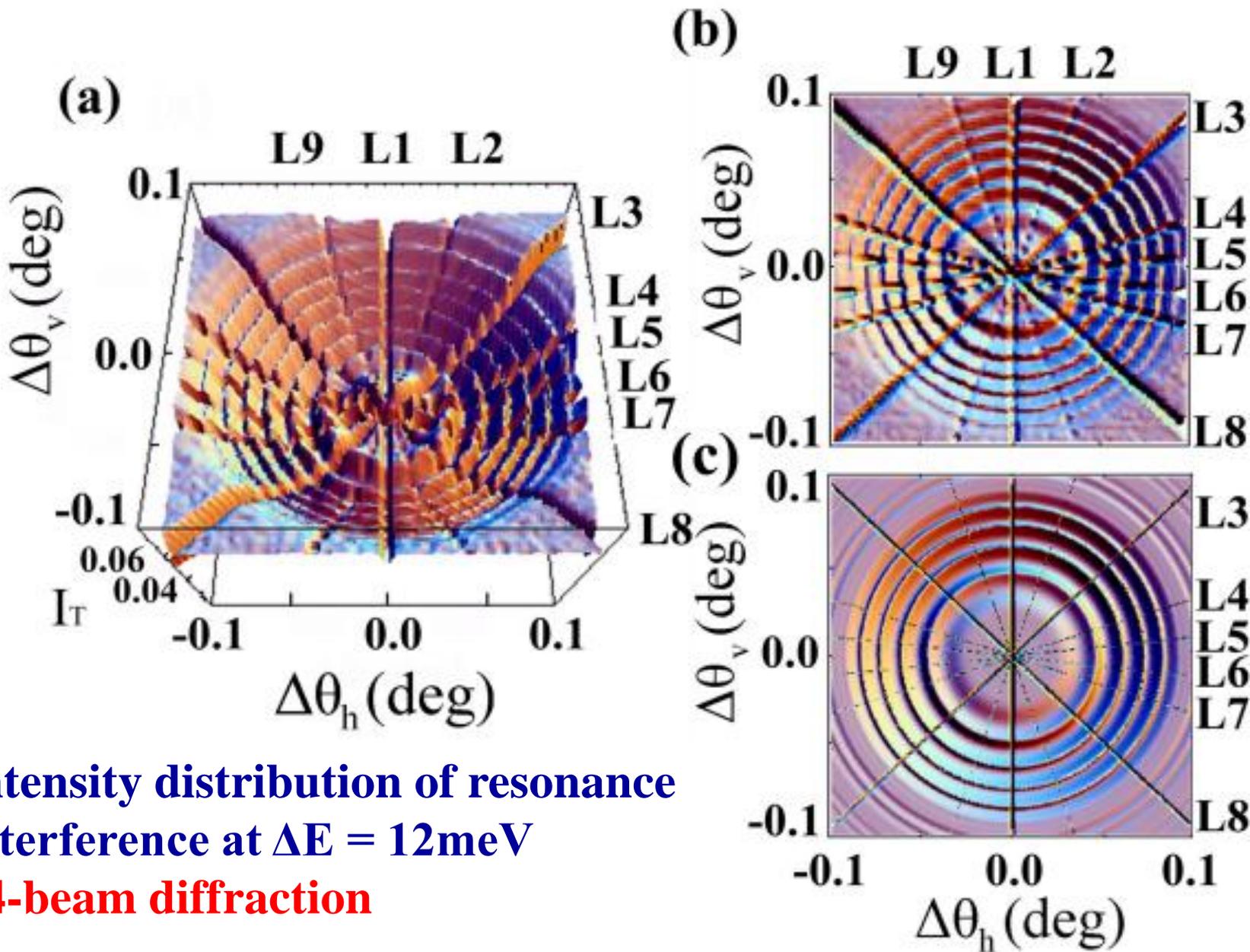


$\Delta E = 0.36$  meV

$\Delta E/E = 2.5 \cdot 10^{-8}$  at  
14.4388 keV ( 0.8588 Å)

- (a) Forward-transmitted  
(0 0 0) beam
- (b) Back-reflected  
(12 4 0) beam

Silicon crystal



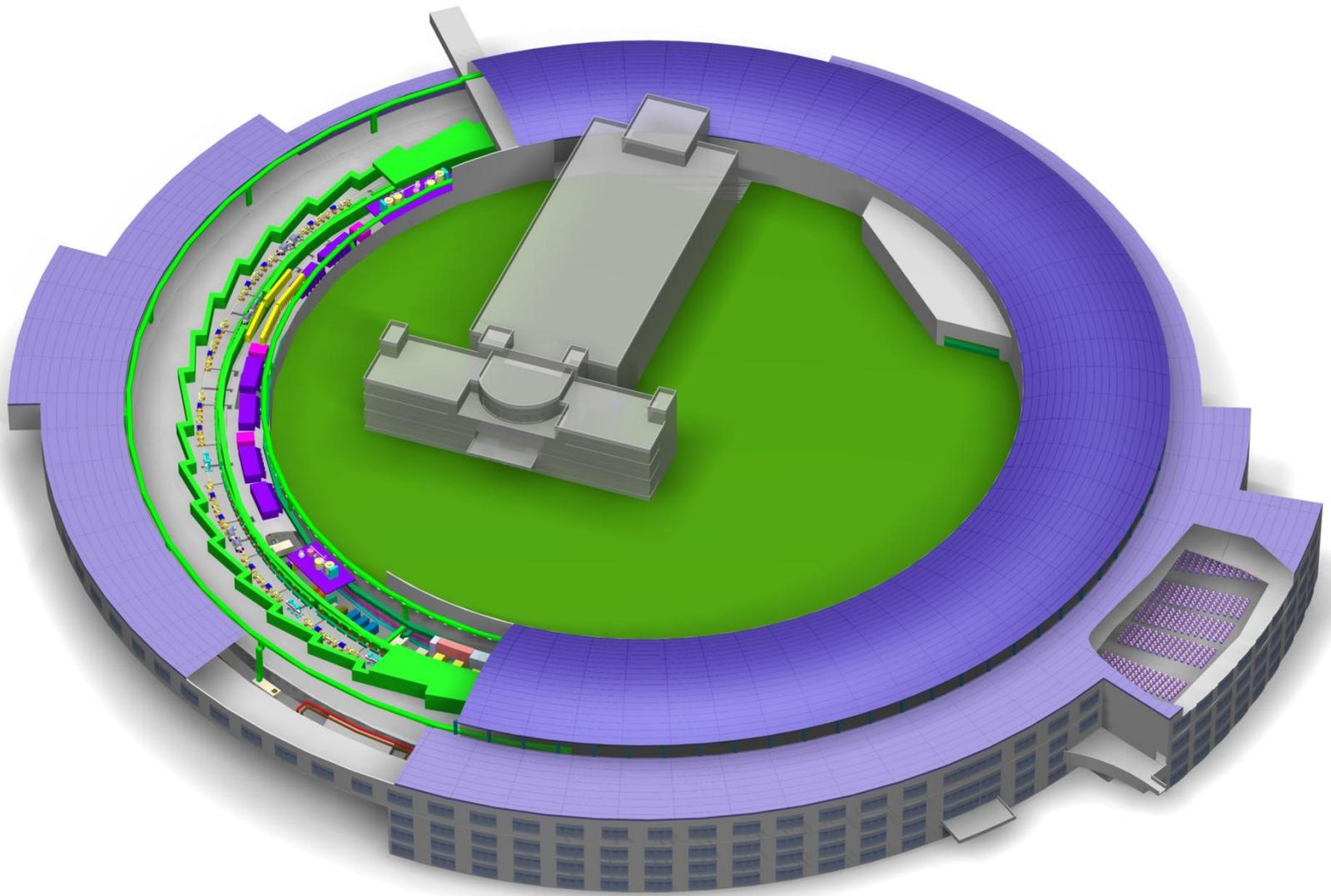
**Intensity distribution of resonance interference at  $\Delta E = 12\text{meV}$**

**24-beam diffraction**

Collaboration: NTHU, NSRRC, SP-8/RIKEN  
 S.-L. Chang, et al *Phys. Rev. Lett.* **94**, 174801 (2005).

台灣光子源計畫

**Taiwan Photon Source (TPS) Project**



**台灣光子源 Taiwan Photon Source (TPS)**

# Taiwan Photon Source (TPS)

3 GeV, 518.4 m, 500 mA

Taiwan Light Source  
(TLS)

Administration and Operation Center

Academic Activity Center

Natural emittance: 1.6 nm-rad  
Straight sections: 7 m (x 18); 12 m (x 6)

**3D Aerial View of NSRRC**

Full capacity: 48 ports

# Progress of Civil Construction



TLS交界處拆除及開挖



深開挖處植入土釘

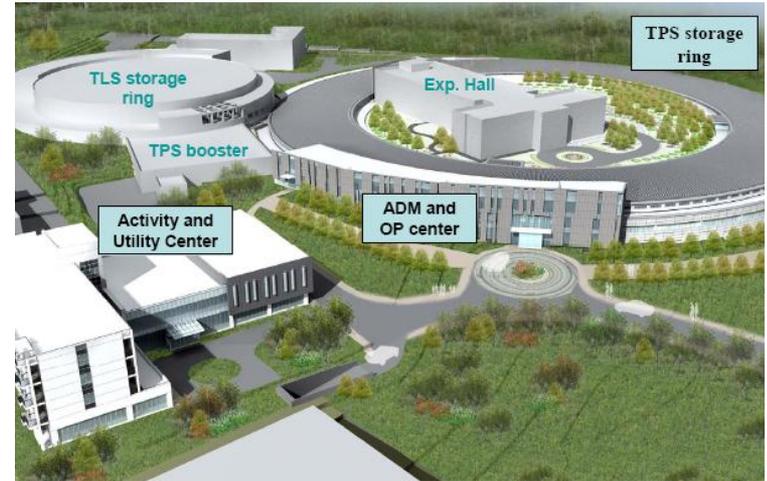


儲存環館第一、二區完成土層開挖



儲存環館鋼筋綁紮準備灌漿

## 台灣光子源興建工程動土典禮 2010.2.7



D棟地下二層開挖整地



地下二層佈植鋼筋灌漿



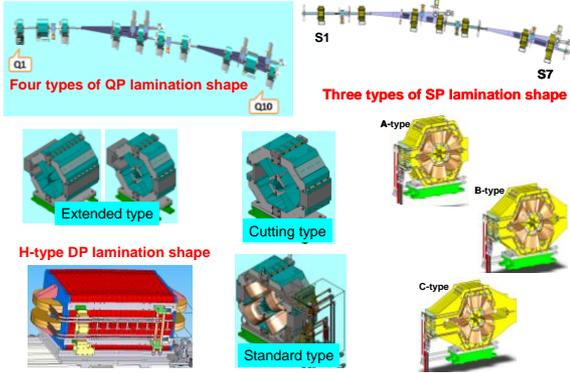
地下二層植筋模板準備灌漿



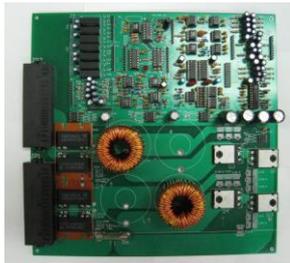
地下一層佈植鋼筋準備灌漿

# TPS Accelerator Design and Prototype

## 二極、四極與六極磁鐵



## 修正磁鐵電源供應器原型 (與工研院合作開發)



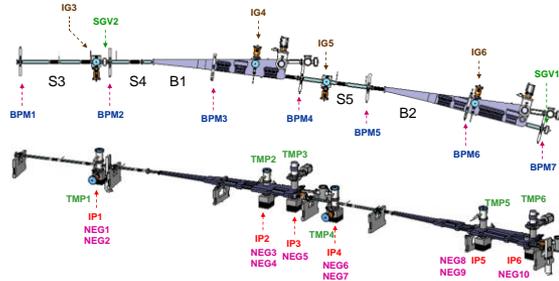
**Specification**  
 Max. volt/cur.:  $\pm 50V/\pm 10A$   
 Current ripple: 10 ppm  
 Short term stability: 5 ppm  
 Long-term stability: 10 ppm

Total 750 units to be fabricated by local company

## TPS儲存環1/24段實體照片



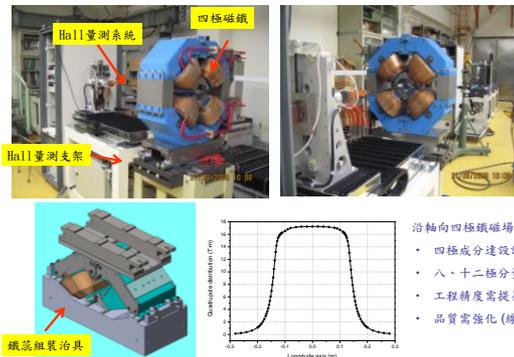
## 真空系統設計與射束診斷安排



## 屏蔽牆內儲存環與增能環結構



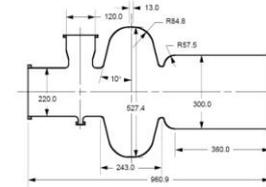
## 四極磁鐵原型及量測平台



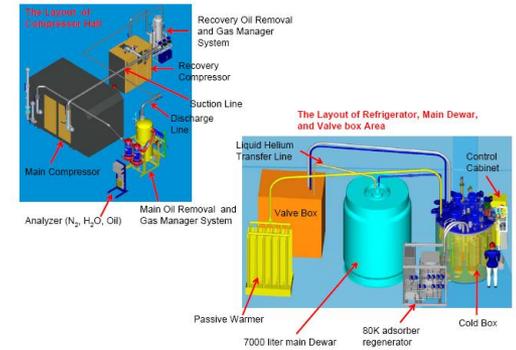
## 超導高頻共振腔

### KEKB Type SRF Module

- Installed at KEKB (508 MHz) and BECP-II/IHEP (500 MHz)



## 700 W液氮低溫系統配置



## 潔淨室無油加工鋁質二極真空腔

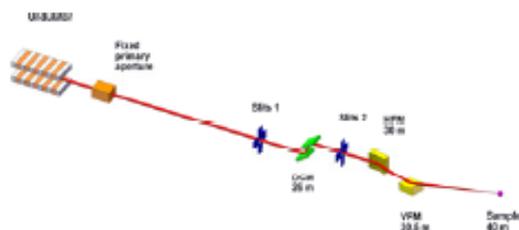


# TPS Phase I Beamlines

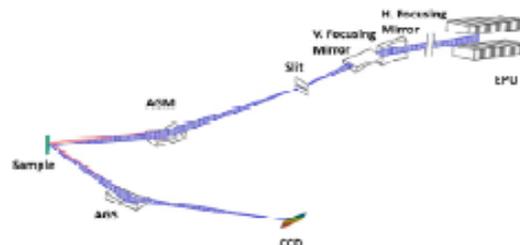
- *$\mu$ -focus macromolecular crystallography* (2013)  
(微聚焦巨分子結晶學光束線)
- *High resolution inelastic soft-x-ray scattering* (2013)  
(高解析非彈性軟X光散射學光束線)
- *Sub- $\mu$  soft x-ray photoelectron & fluorescence emission* (2013)  
(次微米軟X光能譜學光束線)
- *Coherent x-ray scattering (SAXS/XPCS)* (2014)  
(軟物質小角度散射學光束線)
- *Sub- $\mu$  x-ray diffraction* (2014)  
(次微米繞射光束線)
- *Nano-probe* (2014)  
(奈米探針光束線)
- *Temporal coherent x-ray scattering* (2014)  
(時間同調性散射光束線)

# TPS首期7條光束線之概念設計圖

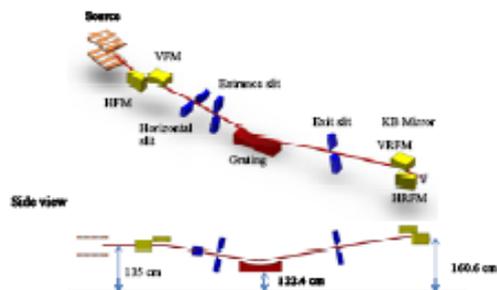
微聚焦巨分子結晶學光束線  
( *$\mu$ -focus macromolecular crystallography*)



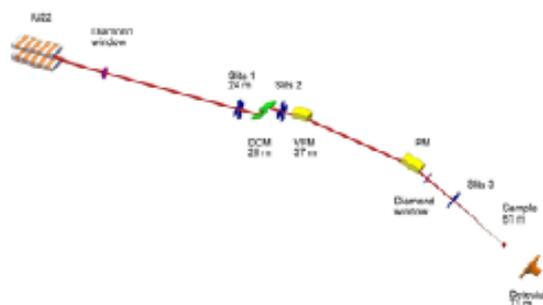
高解析非彈性軟X光散射學光束線  
(*High resolution Inelastic soft-x-ray scattering*)



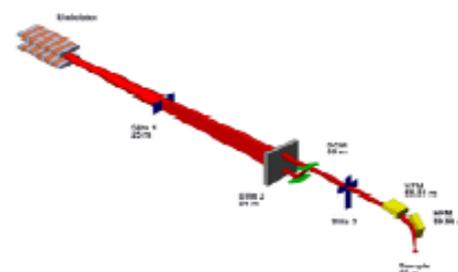
次微米軟X光能譜學光束線  
(*Sub- $\mu$  soft x-ray photoelectron & fluorescence emission*)



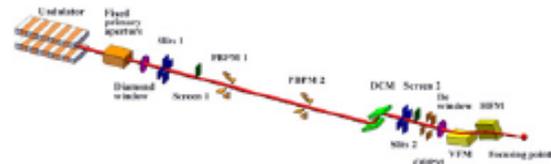
軟物質小角度散射學光束線  
(*Soft matter small angle scattering*)



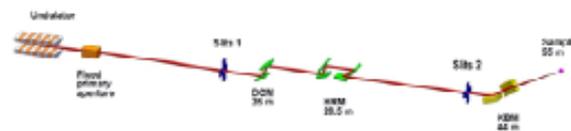
次微米繞射光束線  
(*Sub- $\mu$  x-ray diffraction*)



奈米探針光束線  
(*Nano-probe x-ray diffraction*)



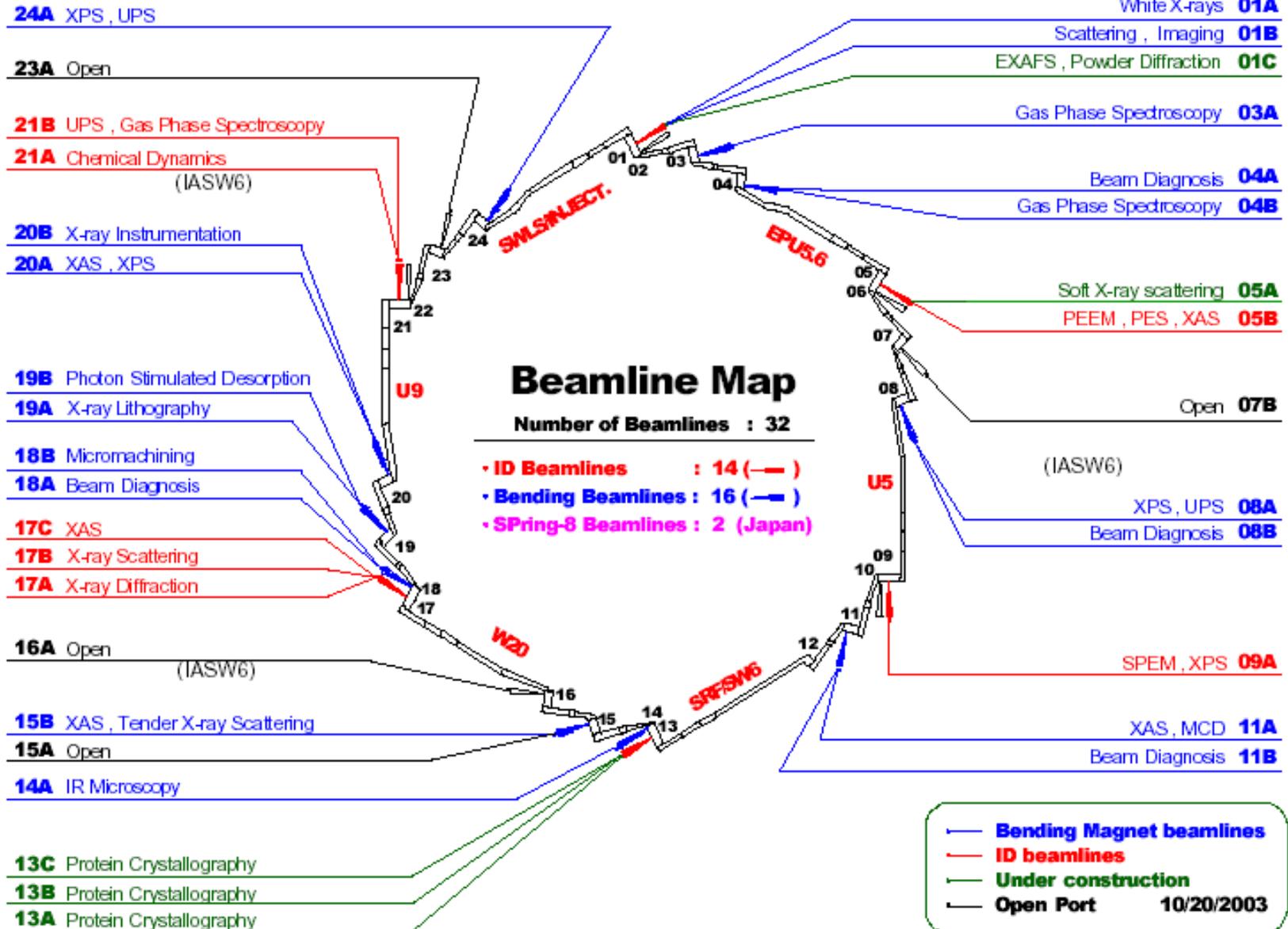
多用途同調性散射學光束線  
(*Multi-purpose coherence x-ray scattering*)



謝謝聆聽 !!

Thank you for your attention !!

# NSRRC Beamline



# Energy scans

