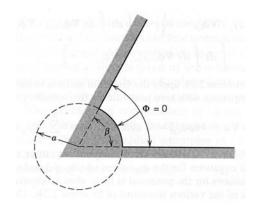
九十八學年第一學期 PHYS5310 Electrodynamics Midterm exam (2 pages) [Jackson Ch.1-4] 2009/11/17, 10:10am-12:00am, 教師:張存續

* 記得寫上學號,班別及姓名等。請依題號順序每頁答一題。

Formula:
$$\frac{1}{\left|\mathbf{x}-\mathbf{x}'\right|} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^{l} \frac{1}{2l+1} \frac{r_{<}^{l}}{r_{>}^{l+1}} Y_{lm}^{*}(\theta', \varphi') Y_{lm}(\theta, \varphi); \quad Y_{l0}\left(\theta, \phi\right) = \sqrt{\frac{2l+1}{4\pi}} P_{l}\left(\cos\theta\right)$$

- 1. (10%, 10%) Using Dirac delta functions in the appropriate coordinates, express the following charge distributions as three-dimensional charge densities $\rho(\mathbf{x})$.
- (a) In cylindrical coordinates, a charge Q per unit length uniformly distributed over a cylindrical surface of radius b.
- (b) In spherical coordinates, a charge Q spreads uniformly over a ring of radius a located on the x-y plane.
- 2. (10%, 10%) The two-dimensional region, $\rho \ge a$, $0 \le \phi \le \beta$, is bounded by conducting surfaces at $\phi = 0$, $\rho \ge a$, and $\phi = \beta$ held at zero potential, as indicated in the figure below. At large ρ the potential is determined by some configuration of charges and/or conductors at fixed potentials.
- (a) Write down a solution for the potential $\Phi(\rho,\phi)$ (that satisfies the boundary conditions for finite ρ).
- (b) Keeping only the lowest non-vanishing terms, calculate the electric field components E_{ρ} and E_{ϕ} and also the surface-charge densities $\sigma(\rho,0)$ and $\sigma(a,\phi)$ on the two boundary surfaces.

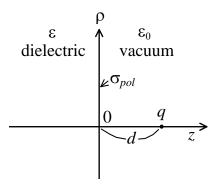


- 3. (10%, 10%) Two point charges +q and -q are located on the z axis at z=+a and z=-a, respectively.
- (a) Find the electrostatic potential as an expansion in spherical harmonics and powers of r for both r > a and r < a.
- (b) Keeping the product $qa \equiv \text{constant}$, take the limit of $a \rightarrow 0$ and find the potential for $r \neq 0$ (i.e. r > a). This is by definition a dipole along the z axis and its potential.

4. (10%, 10%) A point charge q is placed a distant d in front of a semi-infinite dielectrics ε .

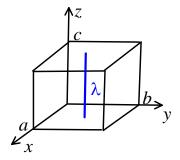
- (a) Find the potential ϕ in both regions (z>0 and z<0).
- (b) Find the polarization surface charge (σ_{pol}) at z=0.

[Hint: Image charge method.]



5.(10%, 10%) Consider a grounded rectangular box with a line-change density in the middle $(\lambda = \lambda_0 \sin(z\pi/c))$.

- (a) Find the Green function for the rectangular box $\nabla^2 G(\mathbf{x}, \mathbf{x}') = -4\pi\delta(\mathbf{x} \mathbf{x}')$ with $G(\mathbf{x}, \mathbf{x}') = 0$ at the boundaries. [Hint: Eigen function expansion.]
- (b) Find the potential inside the box.



1.

(a)
$$\rho(r) = a\delta(r-b)$$
 $\lambda = \int \rho(r)drd\theta = 2\pi ba$ $\Rightarrow \rho = \frac{\lambda}{2\pi b}\delta(r-b)$

(b) Lecture note p.48

The x-y plane is at $\theta = \pi/2$. The charge density $\rho(\mathbf{x})$ can be written as:

Let
$$\rho = c_0 \delta(r - a) \delta(\cos \theta)$$

$$Q = \iiint_{\mathcal{V}} \rho(\mathbf{x}) r^2 \sin \theta dr d\theta d\phi = c_0 2\pi a^2 \implies c_0 = \frac{Q}{2\pi a^2}$$

$$\rho(\mathbf{x}) = \frac{Q}{2\pi a^2} \delta(r - a) \delta(\cos \theta)$$

2. (a)

$$\Phi(\rho,\phi) = (a_0 + b_0 \ln \rho)(A_0 + B_0 \ln \phi) + \sum_{\nu} (a_{\nu}\rho^{\nu} + b_{\nu}\rho^{-\nu})(A_{\nu}\cos\nu\phi + B_{\nu}\sin\nu\phi) \quad (2.69 \text{ and } 2.70)$$

$$\Phi(\rho,\phi) = a_0 (1 + b_0' \ln \rho) (1 + B_0' \ln \phi) + \sum_{\nu} (a_{\nu} \rho^{\nu} + b_{\nu} \rho^{-\nu}) (A_{\nu} \cos \nu \phi + B_{\nu} \sin \nu \phi) \quad (2.69 \text{ and } 2.70)$$

$$\Phi(\rho,0) = 0$$
 $\Rightarrow B'_0 = 0, a_0 = 0, A_v = 0$

$$\Phi(\rho,\beta) = 0$$
 $\Rightarrow v = \frac{m\pi}{\beta}, m = 1,2,3...$

$$\Phi(\alpha,\phi) = 0 = \sum_{m=1}^{\infty} \left(A_m \alpha^{\frac{m\pi}{\beta}} + B_m \alpha^{\frac{-m\pi}{\beta}} \right) \sin \frac{m\pi\phi}{\beta} \quad \Rightarrow B_m = -A_m a^{\frac{2m\pi}{\beta}}$$

$$\Phi(\rho,\phi) = \sum_{m=1}^{\infty} A_m \left(\rho^{\frac{m\pi}{\beta}} - a^{\frac{2m\pi}{\beta}} \rho^{\frac{-m\pi}{\beta}} \right) \sin \frac{m\pi\phi}{\beta}$$

(b). The lowest non-vanishing term corresponding to m=1 is

$$\Phi(\rho,\phi) \approx A_{\rm l} \left(\rho^{\pi/\beta} - a^{\frac{2\pi}{\beta}} \rho^{\frac{-\pi}{\beta}} \right) \sin \frac{\pi \phi}{\beta}$$

$$E_{\rho}(\rho,\phi) = -\frac{\partial \Phi}{\partial \rho} = A_{1} \frac{\pi}{\beta} \left(\rho^{\frac{\pi}{\beta}-1} - a^{\frac{2\pi}{\beta}} \rho^{\frac{-\pi}{\beta}-1} \right) \sin \frac{\pi \phi}{\beta}$$

$$E_{\phi}(\rho,\phi) = -\frac{1}{\rho} \frac{\partial \Phi}{\partial \phi} = A_{1} \frac{\pi}{\beta} \left(\rho^{\pi/\beta - 1} - a^{\frac{2\pi}{\beta}} \rho^{\frac{-\pi}{\beta} - 1} \right) \cos \frac{\pi\phi}{\beta}$$

$$\Rightarrow E_{\rho}(a,\phi) = A_{1} \frac{\pi}{\beta} a^{\pi/\beta^{-1}} \sin \frac{\pi \phi}{\beta}, \ E_{\phi}(a,\phi) = 0, \ E_{\rho}(\rho,0) = 0$$

$$E_{\phi}(\rho,0) = A_{l} \frac{\pi}{\beta} \left(\rho^{\frac{\pi}{\beta}-1} - a^{\frac{2\pi}{\beta}} \rho^{\frac{-\pi}{\beta}-1} \right); E_{\phi}(\rho,\beta) = A_{l} \frac{\pi}{\beta} \left(\rho^{\frac{\pi}{\beta}-1} - a^{\frac{2\pi}{\beta}} \rho^{\frac{-\pi}{\beta}-1} \right)$$

$$\sigma = -\varepsilon_0 \frac{\partial \Phi}{\partial n} \implies \begin{cases} \sigma(a, \phi) = \varepsilon_0 E_{\rho} \Big|_{\rho = a} = -2C_1 \varepsilon_0 \frac{\pi}{\beta} a^{\frac{\pi}{\beta} - 1} \sin \frac{\pi \phi}{\beta} \\ \sigma(\rho, 0) = \varepsilon_0 E_{\phi} \Big|_{\phi = 0} = -C_1 \varepsilon_0 \frac{\pi}{\beta} \left(\frac{\pi}{\rho^{\frac{\pi}{\beta} - 1}} - a^{\frac{2\pi}{\beta}} \rho^{-\frac{\pi}{\beta} - 1} \right) \end{cases}$$

3. (a) Use the binominal expansion and expand in spherical coordinate.

$$\begin{split} \Phi(\vec{x}) &= \frac{1}{4\pi\varepsilon_{0}} \left(\frac{q}{|\vec{x} - \vec{x}_{1}|} - \frac{q}{|\vec{x} - \vec{x}_{2}|} \right) \\ &= \frac{q}{4\pi\varepsilon_{0}} \left[4\pi \sum_{l,m} \frac{1}{2l+1} \frac{r_{<}^{l}}{r_{>}^{l+1}} (Y_{lm}^{*}(0,\phi_{1})Y_{lm}(\theta,\phi) - Y_{lm}^{*}(\pi,\phi_{2})Y_{lm}(\theta,\phi)) \right] \end{split}$$

The problem is azimuthally symmetric, only m=0 terms survive.

$$Y_{lm}^{*}(0,\phi_{1}) = \sqrt{\frac{2l+1}{4\pi}} P_{l}(\cos 0) = \sqrt{\frac{2l+1}{4\pi}}$$

$$Y_{lm}^{*}(\pi,\phi_{2}) = \sqrt{\frac{2l+1}{4\pi}} P_{l}(\cos \pi) = (-1)^{l} \sqrt{\frac{2l+1}{4\pi}}$$

$$\Rightarrow \Phi(\bar{x}) = \frac{q}{\varepsilon_{0}} \sum_{l=1}^{\infty} \frac{1}{2l+1} \frac{2l+1}{4\pi} \frac{r_{<}^{l}}{r_{>}^{l+1}} P_{l}(\cos \theta) \Big[1 - (-1)^{l} \Big]$$

$$\Phi(\bar{x}) = \frac{q}{2\pi\varepsilon_{0}} \sum_{l=odd} \frac{r_{<}^{l}}{r_{>}^{l+1}} P_{l}(\cos \theta)$$

(a) $qa = \frac{p}{2}$, p is the dipole moment

$$\Phi(r > a, \theta, \phi) = \frac{q}{2\pi\varepsilon_0} \sum_{odd} \left(\frac{p}{2q}\right)^l \frac{1}{r^{l+1}} P_l(\cos\theta)$$

$$= \frac{1}{2\pi\varepsilon_0} \sum_{odd} \left(\frac{p}{2q}\right)^l \frac{p^l}{2^l q^{l-1}} P_l(\cos\theta)$$

$$\therefore \Phi \to \frac{p}{4\pi\varepsilon_0 r^2} \cos\theta$$

4. (a) To find Φ in the region $z \ge 0$, we put an image charge q' at z = -d. To find Φ in the region $z \le 0$, we put an image charge q'' at z = +d.

$$\begin{cases} q - q' = q'' \\ \frac{1}{\varepsilon_0} (q + q') = \frac{1}{\varepsilon} q'' \end{cases} \Rightarrow q' = -\frac{\varepsilon - \varepsilon_0}{\varepsilon + \varepsilon_0} q \& q'' = \frac{2\varepsilon}{\varepsilon + \varepsilon_0} q$$

$$\Phi(\rho, \phi, z > 0) = \frac{1}{4\pi\varepsilon_0} \left(\frac{q}{\sqrt{(z-d)^2 + \rho^2}} + \frac{q'}{\sqrt{(z+d)^2 + \rho^2}} \right)$$

$$\Phi(\rho, \phi, z < 0) = \frac{1}{4\pi\varepsilon_0} \left(\frac{q''}{\sqrt{(z-d)^2 + \rho^2}} \right)$$

(b)
$$\nabla \cdot \mathbf{P} = -\rho_{pol} \implies \sigma_{pol} = -\mathbf{P} \cdot \mathbf{n} = -\frac{q}{2\pi} \frac{\varepsilon_0(\varepsilon - \varepsilon_0)}{\varepsilon_0(\varepsilon + \varepsilon_0)} \frac{d}{(\rho^2 + d^2)^{3/2}}$$
 where
$$\mathbf{P} = (\varepsilon - \varepsilon_0) \mathbf{E}(z < 0) = -(\varepsilon - \varepsilon_0) \nabla \Phi(z < 0)$$

5. (a) Consider the corresponding eigenvalue problem with $f(\mathbf{x})=0$ and $\lambda \rightarrow k^2$ $\nabla^2 \psi(\mathbf{x}) + k^2 \psi(\mathbf{x}) = 0$ with the same b.c.

Let
$$\psi(\mathbf{x}) = X(x)Y(y)Z(z) \Rightarrow \underbrace{\frac{1}{X}\frac{d^2X}{dx^2}}_{L^2} + \underbrace{\frac{1}{Y}\frac{d^2Y}{dy^2}}_{L^2} + \underbrace{\frac{1}{Z}\frac{d^2Z}{dz^2}}_{L^2} + k^2 = 0$$

$$\Rightarrow \begin{cases} X(x) = Ae^{ik}l^{x} + Be^{-ik}l^{x} \\ Y(x) = Be^{ik}l^{y} + Ce^{-ik}l^{y} & \text{with } k^{2} = k_{l}^{2} + k_{m}^{2} + k_{n}^{2} \\ Z(x) = De^{ik}l^{z} + Ee^{-ik}l^{z} \end{cases}$$

b.c.
$$\begin{cases} X(x) = 0 \text{ at } x = 0 \text{ and } a \\ Y(x) = 0 \text{ at } y = 0 \text{ and } b \Rightarrow \begin{cases} k_l = \frac{l\pi}{a} \\ k_m = \frac{m\pi}{b} \end{cases} \text{ and } \begin{cases} X = \sin\frac{l\pi x}{a}, \\ Y = \sin\frac{m\pi y}{b}, \\ Z = \sin\frac{m\pi z}{c}, \end{cases}$$

$$\Rightarrow k^2 = k_{lmn}^2 = \pi^2 \left(\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2} \right)$$

$$\Rightarrow \psi(\mathbf{x}) = \sqrt{\frac{8}{abc}} \sin \frac{l\pi x}{a} \sin \frac{m\pi y}{b} \sin \frac{n\pi z}{c}$$

Sub.
$$\psi(\mathbf{x})$$
 into (3.160): $G(\mathbf{x}, \mathbf{x}') = 4\pi \sum_{j} \frac{\psi_{j}^{*}(\mathbf{x}')\psi_{j}(\mathbf{x})}{\lambda_{j} - \lambda}$, we obtain

$$G(\mathbf{x}, \mathbf{x}') = \frac{32}{\pi abc} \sum_{l,m,n=1}^{\infty} \frac{\sin \frac{l\pi x'}{a} \sin \frac{l\pi x}{a} \sin \frac{m\pi y'}{b} \sin \frac{m\pi y}{b} \sin \frac{n\pi z'}{c} \sin \frac{n\pi z}{c}}{\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2}}$$

(b) The potential is given by

$$\begin{split} &\Phi(\mathbf{x}) = \frac{1}{4\pi\varepsilon_0} \int_{\mathcal{V}} \rho(\mathbf{x}') G(\mathbf{x}, \mathbf{x}') d^3 x' - \frac{1}{4\pi} \oint_{\mathcal{S}} \underbrace{\Phi(\mathbf{x}')}_{=0} \frac{\partial}{\partial n'} G(\mathbf{x}, \mathbf{x}') da' \\ &\lambda = \lambda_0 \sin(\frac{z\pi}{c}) \Rightarrow \rho(\mathbf{x}') = \lambda_0 \sin(\frac{z'\pi}{c}) \delta(x' - \frac{a}{2}) \delta(y' - \frac{b}{2}) \\ &G(\mathbf{x}, \mathbf{x}') = \frac{32}{\pi abc} \sum_{l,m,n=1}^{\infty} \frac{\sin\frac{l\pi x'}{a} \sin\frac{l\pi x}{a} \sin\frac{m\pi y'}{b} \sin\frac{m\pi y}{b} \sin\frac{n\pi z'}{c} \sin\frac{n\pi z}{c}}{\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2}} \end{split}$$

$$\pi abc_{l,m,n=1} \qquad \frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2}$$

$$\Phi(\mathbf{x}) = \frac{1}{a^2} + \frac{32}{b^2} \lambda \sum_{l=1}^{\infty} \frac{\sin\frac{l\pi}{2}\sin\frac{m\pi}{2}\sin\frac{l\pi x}{a}\sin\frac{m\pi y}{b}\sin\frac{n\pi z}{c}}{\sin\frac{n\pi z}{c}} \int_{-\infty}^{c} \sin\frac{n\pi z}{a}\sin\frac{n\pi z}{c} ds$$

$$\Phi(\mathbf{x}) = \frac{1}{4\pi\varepsilon_0} \frac{32}{\pi abc} \lambda_0 \sum_{l,m,n=1}^{\infty} \frac{\sin\frac{l\pi}{2}\sin\frac{m\pi}{2}\sin\frac{l\pi x}{a}\sin\frac{m\pi y}{b}\sin\frac{n\pi z}{c}}{\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2}} \int_0^c \sin\frac{n\pi z^2}{c}\sin\frac{n\pi z}{c}\sin\frac{n\pi z}{c}dz'$$

$$\Phi(\mathbf{x}) = \frac{1}{4\pi\varepsilon_0} \frac{32}{\pi abc} \lambda_0 \frac{c}{2} \sum_{l,m=1}^{\infty} \frac{\sin\frac{l\pi}{2}\sin\frac{m\pi}{2}\sin\frac{l\pi x}{a}\sin\frac{m\pi y}{b}\sin\frac{\pi z}{c}}{\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{1}{c^2}}$$