## Chapter 4 Network Analysis

4.0 Introduction

▲ The KVL and KCL in circuit theory are no longer valid. ▲Use the Maxwell's equations to solve all "microwave circuits:?

Field Problem E, H (Vectors)



### Example:

For a complicates coplanar stripline (CPS) low-pass filter, can we design it using circuit point of view?



Agilent 8510 with head modules

## 4.1 Impedance, Equivalent Voltage and Current

Equivalent voltage and current are valid only for TEM lines



For non-TEM lines, equivalent voltage and current are: (a)Not unique but useful for an engineering approach, (b)For certain mode only, (c)  $VI^*$  = power flow of the mode, and

(d)The characteristic impedance is defined as  $Z_0 = V/I$ 

## The Concept of Impedances

An important link between EM field and circuit theory (1) $\eta = \sqrt{\mu/\varepsilon}$  = Intrinsic impedance of a medium. also the wave impedance of a plane wave.

- $(2)Z_{w}$  = Wave impedance (TE, TM, or TEM wave) a characteristic for a particular mode, e.g.the  $TE_{10}$  mode.
- $(3)Z_0$  = Characteristic impedance, the ratio of V and I for a traveling wave on a transmission li

ne 
$$(I_0^-=0, V_0^-=0, V=V_0^+, I=I_0^+)$$



### Example: Application of Waveguide Impedance

Find the reflection of a  $TE_{10}$  wave incident on the interface from z < 0. Let a = 22.86 mm, b = 10.16 mm, and f = 10 GHz.  $z < 0, \varepsilon_{r1} = 1; z > 0, \varepsilon_{r2} = 2.54$ 

$$Sol: \beta_{air} = \sqrt{k_0^2 - \left(\frac{\pi}{a}\right)^2} = 158.0m^{-1}$$

$$\beta_{die} = \sqrt{\varepsilon_r k_0^2 - \left(\frac{\pi}{a}\right)^2} = 304.1m^{-1}$$

$$Z_{Wair} = \frac{k_0 \eta_0}{\beta_{air}} = \frac{209.4 \times 377}{158} = 500\Omega$$

$$Z_{Wdie} = \frac{k_d \eta_d}{\beta_{die}} = \frac{209.4 \times 377}{304.1} = 259.6\Omega$$

$$\Gamma = \frac{Z_{Wdie} - Z_{Wair}}{Z_{Wdie} + Z_{Wair}} = \frac{259.6 - 500}{259.6 + 500} = -0.316$$



Note: The calculation of  $\Gamma$  is valid only when the two waveguides have identical cross section dimensions.

## **One-Port** Network



▲ Input impedance

$$Z_{in} = \frac{V}{I} = \frac{VI^{*}}{|I|^{2}} = \frac{2P}{|I|^{2}} = \frac{P_{\ell} + 2j\omega(W_{m} - W_{e})}{\frac{1}{2}|I|^{2}}$$
$$= R + jX, \quad X = \frac{4\omega(W_{m} - W_{e})}{|I|^{2}} \begin{cases} X > 0, \ W_{m} > W_{e}, \text{ inductive load} \\ X < 0, \ W_{m} < W_{e}, \text{ capactive load} \end{cases}$$

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$$\begin{split} & \psi(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(\omega) e^{i\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \psi(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} V^{*}(\omega) e^{-i\omega t} d\omega \\ & \psi(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} V^{*}(\omega) e^{-i\omega t} d\omega \\ & \text{Since } v(t) \text{ is real } (v(t) = v^{*}(t)), V^{*}(\omega) = V(-\omega) \end{split}$$

$$\begin{split} & \text{Re}[V(\omega)] = \text{Re}[V(-\omega)], \text{ Im}[V(\omega)] = -\text{Im}[V(-\omega)] \\ \text{Re}[I(\omega)] = \text{Re}[I(-\omega)], \text{ Im}[I(\omega)] = -\text{Im}[I(-\omega)] \\ \text{Re}[J(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \psi(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \psi(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(-\omega) e^{-i\omega t} d\omega \\ & \text{Re}[I(\omega)] = \frac{$$

Even and Odd Properties of Z ( $\omega$ ) and  $\Gamma(\omega)$ 

### *N*-Port Microwave Network



- (1) Each ports has  $V_k^+, V_k^-, I_k^+, I_k^-, k = 1, 2, ..., N$ , defined at  $t_k$  plane.
- (2) There could be many modes near each transmission line-network junction.
- (3) Each  $t_k$  plane is chosen at a plane with a suitable distance away from the junction ensuring no higher order mode at  $t_k$ .

## 4.2 Z-Matrix of an N-port Network



 $I_{k} = I_{k}^{+} + I_{k}^{-}$   $V_{k} = \text{ total voltage at port } k.$   $I_{k} = \text{ total current at port } k.$   $\begin{bmatrix} V_{1} \\ V_{2} \\ \vdots \\ V_{N} \end{bmatrix} = \begin{bmatrix} Z_{11} & Z_{12} & \cdots & Z_{1N} \\ Z_{21} & & \vdots \\ \vdots & & \vdots \\ Z_{N1} & \cdots & \cdots & Z_{NN} \end{bmatrix} \begin{bmatrix} I_{1} \\ I_{2} \\ \vdots \\ I_{N} \end{bmatrix},$ 

[Z] = Impedance matrix

 $V_k = V_k^+ + V_k^-$ , k = 1, 2, ..., N at  $t_k$  plane

[V] = [Z][I]For an arbitrary *N*-port linear network: [Z] =Imped

 $Z_{ii} = \frac{V_i}{V_i}$ 

= transfer impedance between ports *i* and *j* when all other ports are open circuited.



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▲ In other words, a two-port specified by [Y] can be represented by



If there is no *active device or anisotropic materials*, like BJT, FET, ferrite and plasma, inside the network, one can use the reciprocity theorem to prove that

 $V_{N}$ 

 $Z_{ij} = Z_{ji}$  and  $Y_{ij} = Y_{ji}$ , i.e. [Z] and [Y] are symmetric matrices.

Lossless Network  

$$\begin{aligned}
& P_{uv} = \frac{1}{2} [V]^{r} [I]^{*} = \frac{1}{2} ([Z]](I)^{r} [I]^{*} = \frac{1}{2} \sum_{m=1}^{N} \sum_{m=1}^{N} I_{m} Z_{m} I_{n}^{*} \\
& For N = 2, P_{uv} = (Z_{11}, I_{1}^{*} + Z_{12}, I_{1}^{*} + Z_{21}, I_{2}^{*} + Z_{21}, I_{2}^{*$$

### **Conversion Between Network Parameters**



## Determine [S] from [Z] and [Y] Matrices



## Example 4.5 Properties of *S*-Parameters

Determine if the network is reciprocal or lossless. What is the reflection coefficient at port 1 if port 2 is shor-circuited?

$$\begin{bmatrix} S \end{bmatrix} = \begin{bmatrix} 0.1 & j0.8 \\ j0.8 & 0.1 \end{bmatrix}$$

Sol: (1)  $[S]^{T} = [S]$  symmetric matrix  $\Leftrightarrow$  reciprocal network, (2)  $[S]^{T} [S]^{*} \neq [U] \Leftrightarrow$  It is not a lossless network. (3) If port 2 is short-circuited,  $V_{2}^{-} = S_{21}V_{1}^{+} + S_{22}V_{2}^{+} = -V_{2}^{+} \Rightarrow V_{2}^{+} = -S_{21}/(1+S_{22})V_{1}^{+}$   $V_{1}^{-} = S_{11}V_{1}^{+} + S_{12}V_{2}^{+} = [S_{11} - S_{12}S_{21}/(1+S_{22})]V_{1}^{+}$  $\Gamma = \frac{V_{1}^{-}}{V_{1}^{+}} = S_{11} - \frac{S_{12}S_{21}}{1+S_{22}} = 0.1 - \frac{-0.64}{1+0.1} \approx 0.682e^{j0^{\circ}}$  Network Parameters for Reciprocal & Lossless Networks (1)[Y], [Z] are symmetric matrices for reciprocal networks and a purely imaginary for lossless networks. (2)[S] is a symmetric matrix for reciprocal networks and a unitary matrix for lossless networks. Multiple what is a unitary matrix? If  $[S]^T[S]^* = [U]$ , where  $[S] = \begin{bmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{bmatrix}$ , i.e.  $[S]^T[S]^* = \begin{bmatrix} S_{11} & S_{21} \\ S_{12} & S_{22} \end{bmatrix} \begin{bmatrix} S_{11}^* & S_{12}^* \\ S_{21}^* & S_{22}^* \end{bmatrix} = \begin{bmatrix} |S_{11}|^2 + |S_{21}|^2 & S_{11}S_{12}^* + S_{21}S_{22}^* \\ S_{11}^*S_{12} + S_{21}^*S_{22} & |S_{12}|^2 + |S_{22}|^2 \end{bmatrix}$   $= \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ A $|S_{11}|^2 + |S_{21}|^2 = |S_{12}|^2 + |S_{22}|^2 = 1$  indicates that the energy is conserved. AS<sub>11</sub><sup>\*</sup>S<sub>12</sub> + S<sub>22</sub>S<sub>21</sub><sup>\*</sup> = 0 means the different columns of [S] are orthogonal. 18

### Shift in Reference Planes









### Example

The characteristic admittance of the two  $\lambda/4$ -section is  $2Y_0$ . Find the *S*-parameters of the two-port with reference impedance  $Z_0 = Y_0^{-1}$ . Based on the *S*-matrix, show that the circuit is lossless.

Find the input reflection coefficients  $\Gamma_{ine}$  and  $\Gamma_{ino}$  when the symmetric plane is placed with open and short circuits, respectively. Using the impedance transformation property, we have  $Z_{ine} = 0$  and  $Z_{ino} = jZ_0 / 2$ , and  $\Gamma_{ine} = -1$  and  $\Gamma_{ino} = (j-2)/(j+2)$ .

$$S_{11} = \frac{1}{2} (\Gamma_{ine} + \Gamma_{ino}) = \frac{1}{2} \left( -1 + \frac{j-2}{j+2} \right) = -\frac{2}{j+2} = S_{22}$$
$$S_{21} = \frac{1}{2} (\Gamma_{ine} - \Gamma_{ino}) = \frac{1}{2} \left( -1 - \frac{j-2}{j+2} \right) = -\frac{j}{j+2} = S_{12}$$

One can validate that the circuit is lossless since  $[S][S]^{T^*}$  = identity matrix.

### Y-Matrix of Symmetric Networks



### Example

Let  $Z_A = 2Z_0$ . Find its Z matrix and  $\Gamma$  at port1 (reference impedance =  $Z_0$ ) if port2 is shortcircuited.



The Z-martix can be obtained via the ABCD matrices in p. 24:

$$\begin{bmatrix} Z_{11} & Z_{12} \\ Z_{21} & Z_{22} \end{bmatrix} = \frac{Z_0}{1+j4} \begin{bmatrix} j & 4 \\ 4 & 2 \end{bmatrix}$$

Applying the even-odd analysis,

$$\begin{bmatrix} V_{1} \\ V_{2} \end{bmatrix} = \begin{bmatrix} V_{eve} + V_{odd} \\ V_{eve} - V_{odd} \end{bmatrix} = \begin{bmatrix} Z_{in}^{eve} & Z_{in}^{odd} \\ Z_{in}^{eve} & -Z_{in}^{odd} \end{bmatrix} \begin{bmatrix} I_{eve} \\ I_{odd} \end{bmatrix} \quad Z_{in}^{eve} = Z_{0} \frac{2Z_{A} + jZ_{o} \tan 45^{\circ}}{Z_{0} + j2Z_{A} \tan 45^{\circ}} = Z_{0} \frac{4 + j}{1 + j4}$$
$$= \begin{bmatrix} Z_{in}^{eve} & Z_{in}^{odd} \\ Z_{in}^{eve} & -Z_{in}^{odd} \end{bmatrix} \begin{bmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & -\frac{1}{2} \end{bmatrix} \begin{bmatrix} I_{1} \\ I_{2} \end{bmatrix} \qquad Z_{in}^{odd} = jZ_{o} \tan 45^{\circ} = jZ_{o}$$
$$Z_{11} = Z_{22} = \frac{Z_{in}^{eve} + Z_{in}^{odd}}{2} = Z_{0} \frac{j}{1 + j4}$$
$$= \frac{1}{2} \begin{bmatrix} Z_{in}^{eve} + Z_{in}^{odd} & Z_{in}^{eve} - Z_{in}^{odd} \\ Z_{in}^{eve} - Z_{in}^{odd} & Z_{in}^{eve} + Z_{in}^{odd} \end{bmatrix} \begin{bmatrix} I_{1} \\ I_{2} \end{bmatrix} \qquad Z_{12} = Z_{21} = \frac{Z_{in}^{eve} - Z_{in}^{odd}}{2} = Z_{0} \frac{4}{1 + j4}$$

### Equivalent Circuit for Two-Port Networks



Measurement plane for port 1 Reference plane for device port 2 Measurement plane for device port 1 Reference plane for device port 2 Measurement plane for port 2 Reference plane for device port 2 Measurement plane for port 2 Measurement plane for

 $\Delta$  <u>Error Box</u> : Losses and phase delays caused by the effects of connectors, cables, and/or transitions.

 $\Delta$  A calibration procedure is used to characterize the *S*-parameters of the error box before measurement of DUT.

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## Theory of TRL Network Analyzer Calibration

### Why calibration is necessary in a microwave measurement?



(a) The S-parameters include magnitude and phase measurement.

(b) Different circuits may have different reference planes, transitions, connectors, housing structures, and/or frequency bands.

Thru---Remove the DUT and Let the S-parameters be [T]



 $S_{21} = S_{12}$  (Reciprocal). The error box are symmetric and identical.

$$u = S_{21}a_1 + S_{22}v = S_{21}a_1 + S_{22}(S_{22}v + S_{21}a_2) = S_{21}\frac{a_1 + S_{22}a_2}{1 - S_{22}^2}$$

$$v = S_{22}u + S_{21}a_2 = S_{21}S_{22}\frac{a_1 + S_{22}a_2}{1 - S_{22}^2} + S_{21}a_2 = S_{21}\frac{S_{21}a_1 + a_2}{1 - S_{22}^2}$$

$$b_1 = S_{11}a_1 + S_{12}v = S_{11}a_1 + S_{12}(S_{21}\frac{S_{21}a_1 + a_2}{1 - S_{22}^2}) \equiv T_{11}a_1 + T_{12}a_2$$

$$T_{11} = S_{11} + \frac{S_{12}S_{21}S_{22}}{1 - S_{22}^2} = S_{11} + \frac{S_{21}^2S_{22}}{1 - S_{22}^2} = T_{22}, T_{12} = \frac{S_{12}S_{21}}{1 - S_{22}^2} = T_{21}$$

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## TRL Calibration (cont'd) $T_{11} = S_{11} + \frac{S_{22}S_{12}^2}{1 - S_{12}^2} = L_{11} = S_{11} + \frac{S_{12}^2S_{22}e^{-2\gamma\ell}}{1 - S_{12}^2 - 2\gamma\ell}$

$$T_{12} = \frac{S_{12}^{2}}{1 - S_{22}^{2}}$$

$$R_{11} = S_{11} + \frac{S_{12}^{2}\Gamma_{L}}{1 - S_{22}^{2}\Gamma_{L}}$$

$$L_{12} = \frac{S_{12}^{2}e^{-\gamma\ell}}{1 - S_{22}^{2}e^{-2\gamma\ell}}$$
Solve  $\underbrace{S_{11}, S_{12}, S_{22}, \Gamma_{L}, and e^{-\gamma\ell}}{5 \text{ unknowns}}$ 

After some simple algebraic manipulations, one can obtain

$$e^{-\gamma\ell} = \frac{L_{12}^{2} + T_{12}^{2} - (T_{11} - L_{11})^{2} \pm \sqrt{\left[L_{12}^{2} + T_{12}^{2} - (T_{11} - L_{11})^{2}\right]^{2} - 4L_{12}^{2}T_{12}^{2}}}{2L_{12}^{2}T_{12}^{2}}$$

$$S_{22} = \frac{T_{11} - L_{11}}{T_{12} - L_{12}e^{-\gamma\ell}}, S_{11} = T_{11} - S_{22}T_{12}, S_{12}^{2} = T_{12}\left(1 - S_{22}^{2}\right)$$

$$\Gamma_{L} = \frac{R_{11} - S_{11}}{S_{12}^{2} + S_{22}(R_{11} - S_{11})}$$

Line---Remove the DUT and Insert a Section 
$$Z_0$$
 Line  

$$\begin{bmatrix}
I_1 \\
I_2 \\
I_3 \\
I_4 \\
I_5 \\$$

### TRL Calibration (cont'd)

- (1) The S-parameters for the error boxes are known.
- (2) The S-parameters for DUT can then be obtained as



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### Optional

### Example (Midterm 2004, part 1)

Find the *Y*-parameters of the network and the input admittances seen at port 1 when port 2 is terminated in  $Z_2 = 0.5Z_0$ ,  $Z_2 = \infty$ . The reference impedance is  $Z_0$ .

$$\overbrace{\ell = 0 \quad \text{Port 1, } Z_o}^{\begin{array}{c} \begin{array}{c} \begin{array}{c} \partial_{\tau} Z_o \\ \partial_{\tau} Z_o \\ \partial_{\tau} Z_o \end{array}} \xrightarrow{\begin{array}{c} \partial_{\tau} Z_o \\ \partial_{\tau} Z_o \\ \partial_{\tau} Z_o \\ \partial_{\tau} Z_o \end{array}} \xrightarrow{\begin{array}{c} \partial_{\tau} Z_o \\ \partial_{\tau}$$

If can be known that  $Y_{11} = Y_{22}, Y_{12} = Y_{21}$ . When  $V_2 = 0, V^- = -V^+$ 

$$\frac{V^{*}e^{j2\theta}}{V^{*}e^{j2\theta}} \qquad V^{*}_{V^{*}} \qquad (1)Y_{11} = \frac{I_{1}}{V_{1}}\Big|_{V_{2}=0} = jY_{0} \tan \theta - j\frac{Y_{0}}{2}\cot 2\theta = Y_{22}$$

$$(2)V_{11} = V^{*}e^{j2\theta} + V^{-}e^{-j2\theta} = 2jV^{*} \sin 2\theta$$

$$I_{2} = -\left(\frac{V^{*}}{2Z_{0}} - \frac{V^{-}}{2Z_{0}}\right) = -\frac{V^{*}}{Z_{0}}$$

$$Y_{12} = Y_{21} = \frac{I_{2}}{V_{1}}\Big|_{V_{2}=0} = j\frac{Y_{0}}{2}\csc 2\theta$$

$$Y_{12} \text{ cannot see the open stub?} \qquad 4$$

## Part IL Excitation of a TE<sub>mn</sub> Waveguide Mode (Pozar 4.7&4.8)

- This study presents a methodology of exciting a pure circular  $TE_{mn}$  mode using cascaded Y-type power dividers.
- •The dividers partition the input signal into several parts which are then coupled to a circular waveguide through apertures. The coupling apertures induce magnetic dipoles. With proper arrangement of the magnetic dipoles on the circumference of the circular waveguide, they then jointly excite the desired mode.
- •The coupling strength and the mode purity are calculated using the reciprocity theorem and the magnetic current sources.
- •Three mode converters,  $TE_{21}$ ,  $TE_{01}$ , and  $TE_{41}$ , were designed, built, and tested at *W*-band. Back-to-back transmission measurements exhibit excellent agreement to the results of computer simulations when the conductor loss is taken into consideration. The measured transmissions are high and the bandwidths are broad.
- These Y-type converters are structurally simple but the machining errors are critical. The factors affecting the performance will be discussed in detail.

T. H. Chang\*, C. S. Li, C. N. Wu, and C. F. Yu, "Exciting circular TEmn modes at low terahertz region", Appl. Phys. Lett. 93, 111503 (2008).



### Example (Midterm 2004 part 2)

Find the *Y*-parameters of the network and the input admittances seen at port 1 when port 2 is terminated in  $Z_2 = 0.5Z_0$ ,  $Z_2 = \infty$ . The reference impedance is  $Z_0$ .



Part II

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### Applications

- Exciting a specific waveguide mode plays a key role in many applications, such as the gyrotron traveling-wave tube (gyro-TWT) and the gyrotron backward-wave oscillator (gyro-BWO).
- In the gyro-TWT, the mode converter launches a wave of a specific mode into the interaction structure to interact with the electron beam; while, in the gyro-BWO, the mode converter extracts the wave power at the upstream end.
- In addition to the gyrotron applications, the mode converters can be used in microwave/plasma systems, radar/antennas systems, and rotary joints.

### Techniques to Excite a Specific Waveguide Mode

- By coupling method, they can be classified into two types: serpentine/corrugated structure and sidewall coupling structure.
- Taking gyrotron as an example, the general requirements for the mode converter/launcher include high converting efficiency, high mode purity, broad bandwidth, high-power capability, and short converting length.
- Y-type converters are superior over other converters in the gyrotron applications.

Part II

### Current Sheets and Arbitrary Current Source



### Part II

# Properties and Characteristics of the Waveguide Modes

Table 1 summarizes the desired modes and their corresponding coupling structures, radii, and parasitic modes.

Desired mode	$TE_{21}$	$TE_{01}$	$TE_{41}$
Coupling structure	Dual-feed	Quad-feed	Quad-feed
Waveguide radius	1.74 mm	2.15 mm	3.00 mm
Parasitic modes	TE <sub>11,A</sub> , TE <sub>11,B</sub>	TE <sub>11,A</sub> , TE <sub>11,B</sub>	TE <sub>11,A</sub> , TE <sub>11,B</sub>
	$TM_{01}$	$TE_{21,A}, TE_{21,B}$	$TE_{21,A}$ , $TE_{21,B}$
		TM <sub>01</sub> ,	$TE_{01}$
		$TM_{11,A}, TM_{11,B}$	$TE_{31,A}$ , $TE_{31,B}$
		1788 - 1787 Maria	$TE_{12,A}, TE_{12,B}$
			TM <sub>01</sub> ,
			$TM_{11,A}, TM_{11,B}$
	No.		TM <sub>21,A</sub> , TM <sub>21,B</sub>

### Part II

### Difficulties to Excite a Higher-Order Mode

Cutoff frequency vs waveguide radius. For the *W*-band operation, the desired waveguide radii are 1.74, 2.15, and 3.00 mm, for  $TE_{21}$ ,  $TE_{01}$ , and  $TE_{41}$ , respectively.



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### Mode Synthesizing

Excitation of a given mode can be achieved using properly arranged electric or magnetic current sources. These current sources come from the probe feeds, the loop feeds, or the coupling apertures. The sidewall coupling apertures, capable of high power operation, thus are the best choice for present study. Mode excitation from an arbitrary electric and magnetic current source can be found in Pozar's Chap.4. Here we present a complete approach to exciting a cylindrical  $TE_{mn}$  mode using multiple magnetic current sources.

- Induce current sources
- Synthesize the desired mode
- Analyze mode purity

### Part II

Electric/Magnetic Dipoles and Currents



TABLE 4.3 Electric and Magnetic Polarizations

$$\begin{aligned} \overline{P}_e &= \alpha_e \varepsilon_0 E_n \hat{n} \delta(\overline{x} - \overline{x}_0) \\ \overline{P}_m &= -\alpha_m \overline{H}_t \delta(\overline{x} - \overline{x}_0) \end{aligned}$$

Aperture Shape	$\alpha_e$	$\alpha_m$
Round hole	$\frac{2r_0^3}{3}$	$\frac{4r_0^3}{3}$
Rectangular slot $(\bar{H} \text{ across slot})$	$\frac{\pi \ell d^2}{16}$	$\frac{\pi \ell d}{16}$

#### Part II

### Induce Current Sources

The idea of the sidewall coupling structure is to equip the waveguide with some dipoles.

Figure shows the geometry of the sidewall coupling structure under study.



The coupling aperture induce equivalent electric and magnetic dipole moments. The electric dipole moment is proportional to the normal electric field while the magnetic dipole moment is proportional to the tangential magnetic field

$$\overline{P}_e = \alpha_e \varepsilon_0 E_n \hat{n} \delta(\overline{x} - \overline{x}_0)$$

$$\overline{P}_m = -\alpha_m \overline{H}_t \delta(\overline{x} - \overline{x}_0)$$

where  $\bar{x}$  and  $\bar{x}_0$  are positions of the observer and the aperture;  $\alpha_e$  and  $\alpha_m$  are constants that depend on the size and shape of the aperture, and  $\varepsilon_0$  is the permittivity.

### Part II

### **Equivalent Polarization Currents**

The electric and magnetic dipole moments are

$$\overline{P}_e = \alpha_e \varepsilon_0 E_n \hat{n} \delta(\overline{x} - \overline{x}_0)$$
  
$$\overline{P}_m = -\alpha_m \overline{H}_t \delta(\overline{x} - \overline{x}_0)$$

The equivalent electric and magnetic polarization currents are

$$\begin{split} \overline{J} &= \sum j\omega \overline{P}_e = 0\\ \overline{M} &= \sum j\omega \mu_0 \overline{P}_m = -j\omega \mu_0 \alpha_m H_z \hat{\mathbf{z}} \sum_i \delta(\overline{x} - \overline{x}_i) \end{split}$$

where the wave is assumed to be sinusoidal time-dependent  $(e^{j\alpha t})$  $\overline{x}_i$  indicates the position of the *i*th coupling hole and  $H_z$  is the amplitude of the tangential magnetic field.

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Part II

## Synthesize the desired mode: the reciprocity theorem (I), Pozar 1.9

The current sources just induced might excite desired mode as well as parasitic modes. In this section we provide a method to show how the desired mode is synthesized and how the parasitic modes are suppressed.

A useful theorem in electromagnetism is to be mentioned---the reciprocity theorem.

Two sets of current sources:  $(\overline{J}_1, \overline{M}_1)$  and  $(\overline{J}_2, \overline{M}_2)$ Their corresponding fields:  $(\overline{E}_1, \overline{H}_1)$  and  $(\overline{E}_2, \overline{H}_2)$ 

Assume  $\overline{M}_2$  to be the only nonzero current source. Then, the reciprocity theorem reads:

$$\oint_{S} (\overline{E}_{1} \times \overline{H}_{1} - \overline{E}_{2} \times \overline{H}_{2}) \cdot ds = \int_{V} (\overline{M}_{2} \cdot \overline{H}_{1}) dv$$

### Part II

Synthesize the desired mode: *How many dipoles are needed?* Side view Top view



Side view of the magnetic field and top view of the electric field for (a)  $TE_{21}$  mode, (b)  $TE_{01}$ mode, and (c)  $TE_{41}$  mode. The field profiles are calculated using HFSS. Two couplings can excite  $TE_{21}$  modes, while four couplings with proper phase control can excite either  $TE_{01}$  or  $TE_{41}$  mode. Part II

### Synthesize the desired mode: the reciprocity theorem (II)

Since  $\overline{J}_1 = \overline{M}_1 = 0$ , the corresponding  $\overline{E}_1$  and  $\overline{H}_1$  are assumed to be certain circular TE<sub>mn</sub> waveguide mode.

 $\overline{E}_2$  and  $\overline{H}_2$  are the electric and magnetic fields due to  $\overline{M}_2$ .

The resulting electric and magnetic fields in the cylindrical waveguide can be expressed as:

$$\overline{E}_{2}^{+} = \sum_{n}^{n} A_{n}^{+} \left(\overline{e}_{n} + \hat{\mathbf{z}}e_{zn}\right) \cdot e^{-j\beta_{n}z}$$

$$\overline{H}_{2}^{+} = \sum_{n}^{n} A_{n}^{+} \left(\overline{h}_{n} + \hat{\mathbf{z}}h_{zn}\right) \cdot e^{-j\beta_{n}z}$$
where  $\left(\overline{e}_{n} + \hat{\mathbf{z}}e_{zn}\right)$  and  $\left(\overline{h}_{n} + \hat{\mathbf{z}}h_{zn}\right)$  are the normalized fields of the *n*th mode and  $A_{n}^{+}$  represents the amplitude of the *n*th mode.

Part II

## Synthesize the desired mode:

### Demonstration

How the magnetic dipoles are formulated? The case for the  $\mathrm{TE}_{01}$  mode is demonstrated.

$$\overline{P}_m = -\alpha_m H_z \hat{\mathbf{z}} \delta(\rho - \rho_0) \delta(z - z_0) \cdot \left| \delta(\phi - 0) + \delta(\phi - \frac{\pi}{2}) + \delta(\phi - \pi) + \delta(\phi - \frac{3\pi}{2}) \right|$$

where  $H_z$  is the tangential *H* field (along *z* direction) and  $\alpha_m$  is a constant depending on the geometry of the hole. Four magnetic dipoles evenly distributed in the circumference are assumed.

With the lengthy calculation, we would find the solution for the amplitude of the *n*th mode:

$$A_{n}^{+} = \frac{1}{P_{n}} \int_{V} (\hat{z}h_{zn}) \cdot j\omega\mu_{0} \overline{P}_{m} e^{j\beta_{n}z} dV = -\frac{1}{P_{n}} \alpha_{m} H_{z0} j\omega\mu_{0} \cdot \left[ h_{zn}(0) + h_{zn}(\frac{\pi}{2}) + h_{zn}(\pi) + h_{zn}(\frac{3\pi}{2}) \right]$$

where  $P_n = 2 \oint_{S_0} ((\overline{e}_n \times \overline{h}_n) \cdot \hat{\mathbf{z}}) ds$  is a normalization constant proportional to the power flow of the nth mode and  $h_{zn}(\phi) = (A \sin m\phi + B \cos m\phi)J_m(p'_{mn})$ .

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### Analyze mode purity

Using the above mentioned approaches, we could synthesize the desired mode. However, some unwanted modes may inevitably be generated and result in a serious mode-competition problem in a gyrotron experiment. Therefore, mode purity is an important issue in the design of a mode converter.

• TE<sub>21</sub> mode  

$$\overline{P}_{m} = -\alpha_{m}H_{z}\delta(\rho - \rho_{0})\delta(z - z_{0})\hat{z} \cdot [\delta(\phi - 0) + \delta(\phi - \pi)]$$

$$A_{\text{TE21,B}}^{+} = -\frac{1}{P_{21}}\alpha_{m}H_{z0}j\omega\mu_{0} \cdot 2BJ_{m}(p'_{mn})$$

$$A_{\text{TE21,A}}^{+} = 0, A_{\text{TE11,A}}^{+} = A_{\text{TE11,B}}^{+} = 0, A_{\text{TM01}}^{+} = 0$$

The arrangement of dual inputs only excites the linear polarization of the  $TE_{21}$  mode, e.g., and luckily, it eliminates non two-fold symmetric modes. So the major parasitic mode  $TE_{11}$  mode cannot be excited.

### Part II

### Analyze mode purity

The field pattern of  $TE_{41}$  mode suggests a quad-feed structure. The magnetic dipoles can be expressed as:

•  $TE_{41}$  mode  $\overline{P}_m = -\alpha_m H_z \delta(\rho - \rho_0) \delta(z - z_0) \hat{\mathbf{z}} \cdot [\delta(\phi - 0) + \delta(\phi - \frac{\pi}{2}) + \delta(\phi - \pi) + \delta(\phi - \frac{3\pi}{2})]$  $A_{\text{TE41,B}}^{+} = -\frac{4B}{P_{41}}\alpha_{m}H_{z0}j\omega\mu_{0}J_{4}(p_{41}')$  $A_{\text{TE01}}^{+} = -\frac{4B}{P_{01}}\alpha_{m}H_{z0}j\omega\mu_{0}J_{0}(p_{01}') \neq 0$  $P_{41}^{total} = \frac{Z_0 k_0 \beta_{41} \pi}{2k_{-41}^k \varepsilon_{04}} (p_{41}'^2 - 4^2) J_4^2 (p_{41}') 16 \cdot (\alpha_m H_{z0} j \omega \mu_0)^2 J_4^2 (p_{41}')$  $P_{01}^{total} = \frac{Z_0 k_0 \beta_{01} \pi}{2k_{c\,01}^4 \varepsilon_{00}} (p_{01}'^2 - 0^2) J_0^2(p_{01}') 16 \cdot (\alpha_m H_{z0} j \omega \mu_0)^2 J_0^2(p_{01}')$ The power ratio is:  $\frac{P_{01}^{total}}{P_{41}^{total}} = \frac{\beta_{41} p'_{01}{}^4 \varepsilon_{04}}{\beta_{01} p'_{41}{}^4 \varepsilon_{00}} \frac{(p'_{01}{}^2 - 0^2)}{(p'_{41}{}^2 - 4^2)} \frac{J_0^4(p'_{01})}{J_4^4(p'_{41})}$ 

Part II

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### Analyze mode purity

The field pattern of  $TE_{01}$  mode is azimuthally symmetric. We can choose dual feeds or quad feeds. Quad-feed structure is sufficient to eliminate these two unwanted modes. The magnetic dipoles can be expressed as:

•  $TE_{01}$  mode

$$\begin{split} \overline{P}_{m} &= -\alpha_{m}H_{z}\delta(\rho - \rho_{0})\delta(z - z_{0})\hat{\mathbf{z}} \cdot [\delta(\phi - 0) + \delta(\phi - \frac{\pi}{2}) + \delta(\phi - \pi) + \delta(\phi - \frac{3\pi}{2})] \\ A_{\text{TE01}}^{+} &= \frac{1}{P_{n}}\int_{V}(\hat{z}h_{zn}) \cdot j\omega\mu_{0}\overline{P}_{m}e^{j\beta_{n}z}dv = -\frac{4B}{P_{n}}\alpha_{m}H_{z0}j\omega\mu_{0}J_{0}(p_{01}') \\ A_{\text{TE11,A}}^{+} &= A_{\text{TE11,B}}^{+} = 0, \ A_{\text{TE21, A}}^{+} = A_{\text{TE21, B}}^{+} = 0 \\ A_{\text{TM01}}^{+} &= 0, \ A_{\text{TM11,A}}^{+} = A_{\text{TM11,B}}^{+} = 0 \end{split}$$

The quad-feed arrangement avoids exciting the parasitic modes. So the mode purity would be very



Calculated field pattern and the mode purity for (a)  $TE_{21}$  mode, (b)  $TE_{01}$ mode, and (c)  $TE_{41}$  mode. The transmissions of the desired modes are shown in blue lines and the reflections are shown in dashed lines. The minor and major parasitic modes are shown in gray and red, respectively.





Two identical converters joined back-to-back for the three modes of interest. Each set consists of three pieces made of oxygenfree high-conductivity copper.

### Part II

### **Experimental Setup**



(a) Photo of the experimental setup for directly measuring the backto-back transmission. The test set is connected to the head modules which are calibrated. The test set is enlarged.

(b) Slotted plate of  $TE_{01}$  converter as an example.

### Part II

## Simulated and Measured Results



Field pattern and transmission for (a)  $TE_{21}$ , (b)  $TE_{01}$ , and (c)  $TE_{41}$ , respectively. The field patterns are HFSS's simulation results. The solid dots represent the measured results and the lines are the simulations. Three different resistivity of the copper are displayed.

### Part II

### Conclusion

Using Y-type power divider to excite pure  $TE_{mn}$  modes was reported. Three mode converters were designed, fabricated, and tested. These converters feature a high back-to-back converting efficiency, high mode purity, broad bandwidth, and compact converting section. Such a converter is suitable for a variety of applications, especially the gyrotrons to generate low-terahertz radiation [34]. At higher frequency, like terahertz region, the micro-fabrication technique is need, which is currently under investigation. The authors would like to thank the technical support of Mr. C. Lee of Ansoft, Taiwan Branch.

T. H. Chang, C. H. Li, C. N. Wu, and C. F. Yu, "Generating pure circular TE<sub>mn</sub> modes using Y-type power dividers", IEEE Trans. Microwave Theory Tech. 58, 1543 (2010).

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The End of Chap. 4